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Functional derivative

In the <u>calculus of variations</u>, a field of <u>mathematical analysis</u>, the **functional derivative** (or **variational derivative**)^[1] relates a change in a Functional (a functional in this sense is a function that acts on functions) to a change in a <u>function</u> on which the functional depends.

In the calculus of variations, functionals are usually expressed in terms of an <u>integral</u> of functions, their <u>arguments</u>, and their <u>derivatives</u>. In an integral *L* of a functional, if a function *f* is varied by adding to it another function δf that is arbitrarily small, and the resulting integrand is expanded in powers of δf , the coefficient of δf in the first order term is called the functional derivative.

For example, consider the functional

$$J[f] = \int_a^b L(x, f(x), f'(x)) dx ,$$

where $f'(x) \equiv df/dx$. If f is varied by adding to it a function δf , and the resulting integrand $L(x, f + \delta f, f' + \delta f')$ is expanded in powers of δf , then the change in the value of J to first order in δf can be expressed as follows:^{[1][Note 1]}

$$\delta J = \int_{a}^{b} \left(\frac{\partial L}{\partial f} \delta f(x) + \frac{\partial L}{\partial f'} \frac{d}{dx} \delta f(x) \right) \, dx = \int_{a}^{b} \left(\frac{\partial L}{\partial f} - \frac{d}{dx} \frac{\partial L}{\partial f'} \right) \delta f(x) \, dx + \frac{\partial L}{\partial f'} (b) \delta f(b) - \frac{\partial L}{\partial f'} (a) \delta f(a)$$

where the variation in the derivative, $\delta f'$ was rewritten as the derivative of the variation (δf) ', and <u>integration by parts</u> was used.

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Definition

In this section, the functional derivative is defined. Then the functional differential is defined in terms of the functional derivative.

Functional derivative

Given a manifold M representing (continuous/smooth) functions ρ (with certain boundary conditions etc.), and a functional F defined as

 $F: M \to \mathbb{R} \quad \text{or} \quad F: M \to \mathbb{C} \,,$

the **functional derivative** of $F[\rho]$, denoted $\delta F/\delta \rho$, is defined through^[2]

$$\int rac{\delta F}{\delta
ho}(x) \phi(x) \; dx = \lim_{arepsilon o 0} rac{F[
ho+arepsilon \phi] - F[
ho]}{arepsilon} \ = \left[rac{d}{darepsilon} F[
ho+arepsilon \phi]
ight]_{arepsilon=0},$$

where ϕ is an arbitrary function. The quantity $\varepsilon \phi$ is called the variation of ρ .

In other words,

$$\phi\mapsto\left[rac{d}{darepsilon}F[
ho+arepsilon\phi]
ight]_{arepsilon=0}$$

is a linear functional, so one may apply the <u>Riesz–Markov–Kakutani representation theorem</u> to represent this functional as integration against some measure. Then $\delta F/\delta \rho$ is defined to be the <u>Radon–Nikodym derivative</u> of this measure.

One thinks of the function $\delta F / \delta \rho$ as the gradient of *F* at the point ρ and

$$\int \frac{\delta F}{\delta \rho}(x)\phi(x)\ dx$$

as the directional derivative at point ρ in the direction of ϕ . Then analogous to vector calculus, the inner product with the gradient gives the directional derivative.

Functional differential

The differential (or variation or first variation) of the functional $F[\rho]$ is $\frac{[3] [Note 2]}{2}$

$$\delta F[
ho;\phi] = \int {\delta F\over \delta
ho}(x) \ \phi(x) \ dx \ dx$$

Heuristically, ϕ is the change in ρ , so we 'formally' have $\phi = \delta \rho$, and then this is similar in form to the <u>total differential</u> of a function $F(\rho_1, \rho_2, \ldots, \rho_n)$,

$$dF = \sum_{i=1}^n rac{\partial F}{\partial
ho_i} \ d
ho_i \ ,$$

where $\rho_1, \rho_2, \ldots, \rho_n$ are independent variables. Comparing the last two equations, the functional derivative $\delta F/\delta \rho(x)$ has a role similar to that of the partial derivative $\partial F/\partial \rho_i$, where the variable of integration x is like a continuous version of the summation index i.^[4]

Properties

Like the derivative of a function, the functional derivative satisfies the following properties, where $F[\rho]$ and $G[\rho]$ are functionals: [Note 3]

Linearity:^[5]

$$rac{\delta(\lambda F+\mu G)[
ho]}{\delta
ho(x)}=\lambdarac{\delta F[
ho]}{\delta
ho(x)}+\murac{\delta G[
ho]}{\delta
ho(x)},$$

where λ , μ are constants.

Product rule:^[6]

$$rac{\delta(FG)[
ho]}{\delta
ho(x)} = rac{\delta F[
ho]}{\delta
ho(x)} G[
ho] + F[
ho] rac{\delta G[
ho]}{\delta
ho(x)}\,,$$

Chain rules:

If *F* is a functional and *G* another functional, then^[7] $\frac{\delta F[G[\rho]]}{\delta \rho(y)} = \int dx \frac{\delta F[G]}{\delta G(x)} \frac{\delta G[\rho](x)}{\delta \rho(y)} \cdot \frac{\delta G[\rho](x)}{\delta \rho(y)} \cdot \frac{\delta G[\rho](x)}{\delta \rho(y)} \cdot \frac{\delta F[g(\rho)]}{\delta \rho(y)} = \frac{\delta F[g(\rho)]}{\delta g[\rho(y)]} \frac{dg(\rho)}{d\rho(y)} \cdot \frac{dg(\rho)}{\delta g[\rho(y)]} \cdot \frac{dg(\rho)}{\delta g[\rho(y)]} \cdot \frac{\delta F[g(\rho)]}{\delta g[\rho(y)]} \frac{dg(\rho)}{\delta g[\rho(y)]} \cdot \frac{\delta F[g(\rho)]}{\delta g[\rho(y)]} = \frac{\delta F[g(\rho)]}{\delta g[\rho(y)]} \frac{dg(\rho)}{\delta g[\rho(y)]} \cdot \frac{\delta F[g(\rho)]}{\delta g[\rho(y)]} \cdot \frac{$

Determining functional derivatives

A formula to determine functional derivatives for a common class of functionals can be written as the integral of a function and its derivatives. This is a generalization of the <u>Euler–Lagrange equation</u>: indeed, the functional derivative was introduced in <u>physics</u> within the derivation of the <u>Lagrange</u> equation of the second kind from the <u>principle of least action</u> in <u>Lagrangian mechanics</u> (18th century). The first three examples below are taken from <u>density functional theory</u> (20th century), the fourth from <u>statistical mechanics</u> (19th century).

Formula

Given a functional

$$F[
ho] = \int f(oldsymbol{r},
ho(oldsymbol{r}),
abla
ho(oldsymbol{r}))\,doldsymbol{r},$$

and a function $\phi(\mathbf{r})$ that vanishes on the boundary of the region of integration, from a previous section Definition,

$$\begin{split} \int \frac{\delta F}{\delta \rho(\boldsymbol{r})} \,\phi(\boldsymbol{r}) \,d\boldsymbol{r} &= \left[\frac{d}{d\varepsilon} \int f(\boldsymbol{r},\rho+\varepsilon\phi,\nabla\rho+\varepsilon\nabla\phi) \,d\boldsymbol{r} \right]_{\varepsilon=0} \\ &= \int \left(\frac{\partial f}{\partial \rho} \,\phi + \frac{\partial f}{\partial \nabla \rho} \cdot \nabla\phi \right) d\boldsymbol{r} \\ &= \int \left[\frac{\partial f}{\partial \rho} \,\phi + \nabla \cdot \left(\frac{\partial f}{\partial \nabla \rho} \,\phi \right) - \left(\nabla \cdot \frac{\partial f}{\partial \nabla \rho} \right) \phi \right] d\boldsymbol{r} \\ &= \int \left[\frac{\partial f}{\partial \rho} \,\phi - \left(\nabla \cdot \frac{\partial f}{\partial \nabla \rho} \right) \phi \right] d\boldsymbol{r} \\ &= \int \left(\frac{\partial f}{\partial \rho} - \nabla \cdot \frac{\partial f}{\partial \nabla \rho} \right) \phi(\boldsymbol{r}) \,d\boldsymbol{r} \,. \end{split}$$

The second line is obtained using the total derivative, where $\partial f / \partial \nabla \rho$ is a derivative of a scalar with respect to a vector. [Note 4]

The third line was obtained by use of a product rule for divergence. The fourth line was obtained using the divergence theorem and the condition that $\phi=0$ on the boundary of the region of integration. Since ϕ is also an arbitrary function, applying the fundamental lemma of calculus of variations to the last line, the functional derivative is

$$rac{\delta F}{\delta
ho(m{r})} = rac{\partial f}{\partial
ho} -
abla \cdot rac{\partial f}{\partial
abla
ho}$$

where $\rho = \rho(\mathbf{r})$ and $f = f(\mathbf{r}, \rho, \nabla \rho)$. This formula is for the case of the functional form given by $F[\rho]$ at the beginning of this section. For other functional forms, the definition of the functional derivative can be used as the starting point for its determination. (See the example Coulomb potential energy functional.)

The above equation for the functional derivative can be generalized to the case that includes higher dimensions and higher order derivatives. The functional would be,

$$F[
ho(oldsymbol{r})] = \int f(oldsymbol{r},
ho(oldsymbol{r}),
abla
ho(oldsymbol{r}),
abla^{(2)}
ho(oldsymbol{r}),\dots,
abla^{(N)}
ho(oldsymbol{r}))\,doldsymbol{r},$$

where the vector $\mathbf{r} \in \mathbb{R}^n$, and $\nabla^{(i)}$ is a tensor whose n^i components are partial derivative operators of order i,

An analogous application of the definition of the functional derivative yields

$$egin{aligned} rac{\delta F[
ho]}{\delta
ho} &= rac{\partial f}{\partial
ho} -
abla \cdot rac{\partial f}{\partial (
abla
ho)} +
abla^{(2)} \cdot rac{\partial f}{\partial \left(
abla^{(2)}
ho
ight)} + \cdots + (-1)^N
abla^{(N)} \cdot rac{\partial f}{\partial \left(
abla^{(N)}
ho
ight)} \ &= rac{\partial f}{\partial
ho} + \sum_{i=1}^N (-1)^i
abla^{(i)} \cdot rac{\partial f}{\partial \left(
abla^{(i)}
ho
ho
ight)} \,. \end{aligned}$$

In the last two equations, the n^i components of the tensor $\frac{\partial f}{\partial (\nabla^{(i)} \rho)}$ are partial derivatives of f with respect to partial derivatives

of ρ,

$$\left[rac{\partial f}{\partial\left(
abla^{(i)}
ho
ight)}
ight]_{lpha_1lpha_2\cdotslpha_i} = rac{\partial f}{\partial
ho_{lpha_1lpha_2\cdotslpha_i}} \qquad ext{ where } \ \
ho_{lpha_1lpha_2\cdotslpha_i} \equiv rac{\partial^{\,i}
ho}{\partial r_{lpha_1}\ \partial r_{lpha_2}\cdots\partial r_{lpha_i}} \ ,$$

and the tensor scalar product is,

$$\nabla^{(i)} \cdot \frac{\partial f}{\partial \left(\nabla^{(i)} \rho\right)} = \sum_{\alpha_1, \alpha_2, \cdots, \alpha_i = 1}^n \frac{\partial^i}{\partial r_{\alpha_1} \partial r_{\alpha_2} \cdots \partial r_{\alpha_i}} \frac{\partial f}{\partial \rho_{\alpha_1 \alpha_2 \cdots \alpha_i}} \cdot \frac{[\text{Note 6}]}{\partial \rho_{\alpha_1 \alpha_2 \cdots \alpha_i}}$$

Examples

Thomas-Fermi kinetic energy functional

The <u>Thomas–Fermi model</u> of 1927 used a kinetic energy functional for a noninteracting uniform <u>electron gas</u> in a first attempt of density-functional theory of electronic structure:

$$T_{
m TF}[
ho] = C_{
m F} \int
ho^{5/3}({f r}) \, d{f r} \, .$$

Since the integrand of $T_{\text{TF}}[\rho]$ does not involve derivatives of $\rho(\mathbf{r})$, the functional derivative of $T_{\text{TF}}[\rho]$ is,^[9]

$$egin{aligned} rac{\delta T_{ ext{TF}}}{\delta
ho(m{r})} &= C_{ ext{F}} rac{\partial
ho^{5/3}(\mathbf{r})}{\partial
ho(\mathbf{r})} \ &= rac{5}{3} C_{ ext{F}}
ho^{2/3}(\mathbf{r}) \,. \end{aligned}$$

Coulomb potential energy functional

For the electron-nucleus potential, Thomas and Fermi employed the Coulomb potential energy functional

$$V[
ho] = \int rac{
ho(m{r})}{|m{r}|} \, dm{r}.$$

Applying the definition of functional derivative,

$$\int rac{\delta V}{\delta
ho(m{r})} \, \phi(m{r}) \, dm{r} = \left[rac{d}{darepsilon} \int rac{
ho(m{r}) + arepsilon \phi(m{r})}{|m{r}|} \, dm{r}
ight]_{arepsilon=0}
onumber \ = \int rac{1}{|m{r}|} \, \phi(m{r}) \, dm{r} \, .$$

So,

$$rac{\delta V}{\delta
ho(m{r})} = rac{1}{|m{r}|} \ .$$

For the classical part of the electron-electron interaction, Thomas and Fermi employed the Coulomb potential energy functional

$$J[
ho] = rac{1}{2} \iint rac{
ho(\mathbf{r})
ho(\mathbf{r}')}{|\mathbf{r}-\mathbf{r}'|} \, d\mathbf{r} d\mathbf{r}' \, .$$

From the definition of the functional derivative,

$$egin{aligned} &\int rac{\delta J}{\delta
ho(m{r})} \phi(m{r}) dm{r} = \left[rac{d}{d\epsilon} J[
ho+\epsilon \phi]
ight]_{\epsilon=0} \ &= \left[rac{d}{d\epsilon} \left(rac{1}{2} \iint rac{[
ho(m{r})+\epsilon \phi(m{r})] \left[
ho(m{r}')+\epsilon \phi(m{r}')
ight]}{|m{r}-m{r}'|} \, dm{r} dm{r}'
ight)
ight]_{\epsilon=0} \ &= rac{1}{2} \iint rac{
ho(m{r}') \phi(m{r})}{|m{r}-m{r}'|} \, dm{r} dm{r}' + rac{1}{2} \iint rac{
ho(m{r}) \phi(m{r}')}{|m{r}-m{r}'|} \, dm{r} dm{r}' \end{aligned}$$

The first and second terms on the right hand side of the last equation are equal, since r and r' in the second term can be interchanged without changing the value of the integral. Therefore,

$$\int rac{\delta J}{\delta
ho(m{r})} \phi(m{r}) dm{r} = \int \left(\int rac{
ho(m{r'})}{|m{r}-m{r'}|} dm{r'}
ight) \phi(m{r}) dm{r}$$

and the functional derivative of the electron-electron coulomb potential energy functional $J[\rho]$ is, $\underline{[10]}$

$$rac{\delta J}{\delta
ho(m{r})} = \int rac{
ho(m{r}')}{|m{r}-m{r}'|} dm{r}'\,.$$

The second functional derivative is

$$rac{\delta^2 J[
ho]}{\delta
ho({f r'})\delta
ho({f r})} = rac{\partial}{\partial
ho({f r'})} \left(rac{
ho({f r'})}{|{f r}-{f r'}|}
ight) = rac{1}{|{f r}-{f r'}|}.$$

Weizsäcker kinetic energy functional

In 1935 <u>von Weizsäcker</u> proposed to add a gradient correction to the Thomas-Fermi kinetic energy functional to make it suit better a molecular electron cloud:

$$T_{\mathrm{W}}[
ho] = rac{1}{8} \int rac{
abla
ho(\mathbf{r}) \cdot
abla
ho(\mathbf{r})}{
ho(\mathbf{r})} d\mathbf{r} = \int t_{\mathrm{W}} \; d\mathbf{r} \, ,$$

where

$$t_{
m W} \equiv rac{1}{8} rac{
abla
ho \cdot
abla
ho}{
ho} \qquad ext{and} \ \
ho =
ho(m{r}) \ .$$

Using a previously derived formula for the functional derivative,

$$egin{aligned} rac{\delta T_{\mathrm{W}}}{\delta
ho(m{r})} &= rac{\partial t_{\mathrm{W}}}{\partial
ho} -
abla \cdot rac{\partial t_{\mathrm{W}}}{\partial
abla
ho} \ &= -rac{1}{8} rac{
abla
ho \cdot
abla
ho}{
ho^2} - \left(rac{1}{4} rac{
abla^2
ho}{
ho} - rac{1}{4} rac{
abla
ho \cdot
abla
ho}{
ho^2}
ight) \qquad ext{where} \ \
abla^2 &=
abla \cdot
abla \ , \end{aligned}$$

and the result is, [11]

$$rac{\delta T_{
m W}}{\delta
ho(m{r})} = -rac{1}{8}rac{
abla
ho \cdot
abla
ho}{
ho^2} - rac{1}{4}rac{
abla^2
ho}{
ho} \;.$$

Entropy

The entropy of a discrete random variable is a functional of the probability mass function.

$$H[p(x)] = -\sum_x p(x) \log p(x)$$

Thus,

$$egin{aligned} &\sum_x rac{\delta H}{\delta p(x)} \, \phi(x) = \left[rac{d}{d\epsilon} H[p(x) + \epsilon \phi(x)]
ight]_{\epsilon=0} \ &= \left[-rac{d}{d\epsilon} \sum_x \left[p(x) + \epsilon \phi(x)
ight] \, \log[p(x) + \epsilon \phi(x)]
ight]_{\epsilon=0} \ &= -\sum_x \left[1 + \log p(x)
ight] \, \phi(x) \, . \end{aligned}$$

Thus,

$$rac{\delta H}{\delta p(x)} = -1 - \log p(x).$$

Exponential

Let

$$F[arphi(x)]=e^{\int arphi(x)g(x)dx}.$$

Using the delta function as a test function,

$$\begin{split} \frac{\delta F[\varphi(x)]}{\delta \varphi(y)} &= \lim_{\varepsilon \to 0} \frac{F[\varphi(x) + \varepsilon \delta(x-y)] - F[\varphi(x)]}{\varepsilon} \\ &= \lim_{\varepsilon \to 0} \frac{e^{\int (\varphi(x) + \varepsilon \delta(x-y))g(x)dx} - e^{\int \varphi(x)g(x)dx}}{\varepsilon} \\ &= e^{\int \varphi(x)g(x)dx} \lim_{\varepsilon \to 0} \frac{e^{\varepsilon \int \delta(x-y)g(x)dx} - 1}{\varepsilon} \\ &= e^{\int \varphi(x)g(x)dx} \lim_{\varepsilon \to 0} \frac{e^{\varepsilon g(y)} - 1}{\varepsilon} \\ &= e^{\int \varphi(x)g(x)dx} g(y). \end{split}$$

Thus,

$$rac{\delta F[arphi(x)]}{\delta arphi(y)} = g(y) F[arphi(x)].$$

This is particularly useful in calculating the correlation functions from the partition function in quantum field theory.

Functional derivative of a function

A function can be written in the form of an integral like a functional. For example,

$$ho(oldsymbol{r})=F[
ho]=\int
ho(oldsymbol{r}')\delta(oldsymbol{r}-oldsymbol{r}')\,doldsymbol{r}'.$$

Since the integrand does not depend on derivatives of ρ , the functional derivative of $\rho(\mathbf{r})$ is,

$$rac{\delta
ho(m{r})}{\delta
ho(m{r}')}\equiv rac{\delta F}{\delta
ho(m{r}')}=rac{\partial}{\partial
ho(m{r}')}\left[
ho(m{r}')\delta(m{r}-m{r}')
ight] =\delta(m{r}-m{r}').$$

Functional derivative of iterated function

The functional derivative of the iterated function f(f(x)) is given by:

$$rac{\delta f(f(x))}{\delta f(y)} = f'(f(x))\delta(x-y) + \delta(f(x)-y)$$

and

$$rac{\delta f(f(f(x)))}{\delta f(y)}=f'(f(f(x))(f'(f(x))\delta(x-y)+\delta(f(x)-y))+\delta(f(f(x))-y))$$

In general:

$$rac{\delta f^N(x)}{\delta f(y)}=f'(f^{N-1}(x))rac{\delta f^{N-1}(x)}{\delta f(y)}+\delta(f^{N-1}(x)-y)$$

Putting in N=0 gives:

$$rac{\delta f^{-1}(x)}{\delta f(y)} = -rac{\delta (f^{-1}(x)-y)}{f'(f^{-1}(x))}$$

Using the delta function as a test function

In physics, it is common to use the <u>Dirac delta function</u> $\delta(x - y)$ in place of a generic test function $\phi(x)$, for yielding the functional derivative at the point y (this is a point of the whole functional derivative as a <u>partial derivative</u> is a component of the gradient):^[12]

$$rac{\delta F[
ho(x)]}{\delta
ho(y)} = \lim_{arepsilon o 0} rac{F[
ho(x) + arepsilon \delta(x-y)] - F[
ho(x)]}{arepsilon}.$$

This works in cases when $F[\rho(x) + \varepsilon f(x)]$ formally can be expanded as a series (or at least up to first order) in ε . The formula is however not mathematically rigorous, since $F[\rho(x) + \varepsilon \delta(x - y)]$ is usually not even defined.

The definition given in a previous section is based on a relationship that holds for all test functions $\phi(x)$, so one might think that it should hold also when $\phi(x)$ is chosen to be a specific function such as the <u>delta function</u>. However, the latter is not a valid test function (it is not even a proper function).

In the definition, the functional derivative describes how the functional $F[\rho(x)]$ changes as a result of a small change in the entire function $\rho(x)$. The particular form of the change in $\rho(x)$ is not specified, but it should stretch over the whole interval on which x is defined. Employing the particular form of the perturbation given by the delta function has the meaning that $\rho(x)$ is varied only in the point y. Except for this point, there is no variation in $\rho(x)$.

Notes

- 1. According to Giaquinta & Hildebrandt (1996), p. 18, this notation is customary in physical literature.
- 2. Called *differential* in (Parr & Yang 1989, p. 246), *variation* or *first variation* in (Courant & Hilbert 1953, p. 186), and *variation* or *differential* in (Gelfand & Fomin 2000, p. 11, § 3.2).

3. Here the notation
$$\frac{\delta F}{\delta \rho}(x) \equiv \frac{\delta F}{\delta \rho(x)}$$
 is introduced.

4. For a three-dimensional cartesian coordinate system,

$$rac{\partial f}{\partial
abla
ho} = rac{\partial f}{\partial
ho_x} \mathbf{\hat{i}} + rac{\partial f}{\partial
ho_y} \mathbf{\hat{j}} + rac{\partial f}{\partial
ho_z} \mathbf{\hat{k}}, \quad ext{ where }
ho_x = rac{\partial
ho}{\partial x}, \
ho_y = rac{\partial
ho}{\partial y}, \
ho_z = rac{\partial
ho}{\partial z} \ ext{and } \mathbf{\hat{i}}, \ \mathbf{\hat{j}}, \ \mathbf{\hat{k}} ext{ are unit vectors along the x, y, z axes.}$$

5. For example, for the case of three dimensions (n = 3) and second order derivatives (i = 2), the tensor $\nabla^{(2)}$ has components,

$$\left[
abla^{(2)}
ight]_{lphaeta} = rac{\partial^{\,2}}{\partial r_{lpha}\,\partial r_{eta}} \qquad ext{ where } \quad lpha,eta=1,2,3 \ .$$

6. For example, for the case n = 3 and i = 2, the tensor scalar product is,

$$abla^{(2)} \cdot rac{\partial f}{\partial \left(
abla^{(2)}
ho
ight)} = \sum_{lpha, eta=1}^3 \; rac{\partial^{\,2}}{\partial r_lpha \; \partial r_eta} \; rac{\partial f}{\partial
ho_{lphaeta}} \qquad ext{where} \; \;
ho_{lphaeta} \equiv rac{\partial^{\,2}
ho}{\partial r_lpha \; \partial r_eta} \; .$$

Footnotes

- 1. (Giaquinta & Hildebrandt 1996, p. 18)
- 2. (Parr & Yang 1989, p. 246, Eq. A.2).
- 3. (Parr & Yang 1989, p. 246, Eq. A.1).
- 4. (Parr & Yang 1989, p. 246).
- 5. (Parr & Yang 1989, p. 247, Eq. A.3).
- 6. (Parr & Yang 1989, p. 247, Eq. A.4).
- 7. (Greiner & Reinhardt 1996, p. 38, Eq. 6).
- 8. (Greiner & Reinhardt 1996, p. 38, Eq. 7).
- 9. (Parr & Yang 1989, p. 247, Eq. A.6).
- 10. (Parr & Yang 1989, p. 248, Eq. A.11).
- 11. (Parr & Yang 1989, p. 247, Eq. A.9).
- 12. Greiner & Reinhardt 1996, p. 37

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External links

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