

# Precision Measurement of the Proton Elastic Form Factor Ratio at Low $Q^2$

Shalev Gilad, MIT

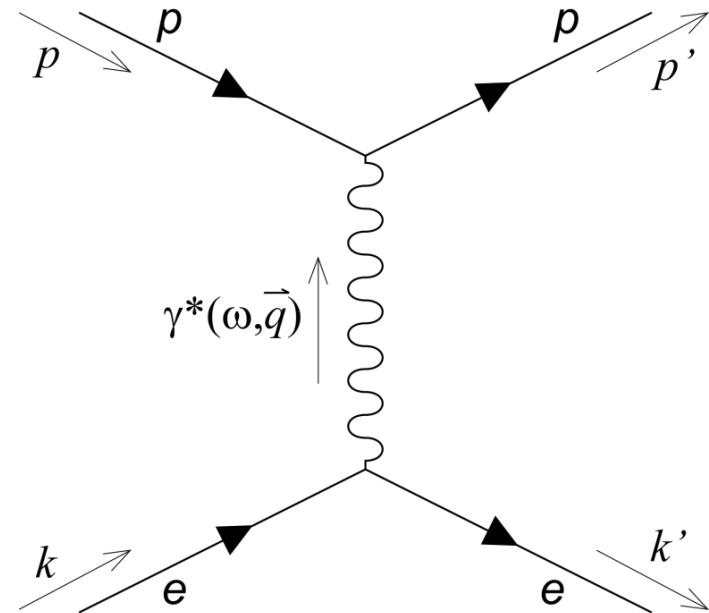
- Nucleons Form Factors
- Recoil polarimetry
- World data
- New results
- Impact
- Outlook and summary

# Dirac and Pauli Form Factors

- Pioneered by Hofstadter *et. al* at Stanford in 1950s, first proton form factor measurement reported in 1955
- As theory for Strong force, QCD has been tested well in the asymptotic region, understanding hadron structure in confinement region still challenging
- Dirac and Pauli form factors:  $F_1$  ,  $F_2$

$$J_{hadronic}^{\mu} = e\bar{u}(p')[\gamma^{\mu}F_1(Q^2) + \frac{i\sigma^{\mu\nu}q_{\nu}}{2M}F_2(Q^2)]u(p)$$

$$Q^2 = -q^2$$



single photon exchange  
(Born approximation)

$$\frac{d\sigma}{d\Omega} = \sigma_{Mott} \frac{1}{1 + \tau} \{ F_1^2(Q^2) + \tau [ F_2^2(Q^2) + 2(F_1(Q^2) + F_2(Q^2))^2 \tan^2 \frac{\theta_c}{2} ] \}$$

# Sachs Form Factors

- Linear combination of  $F_1$  and  $F_2$
- Fourier transform of the charge (magnetization) densities in the Breit frame at the non-relativistic limit

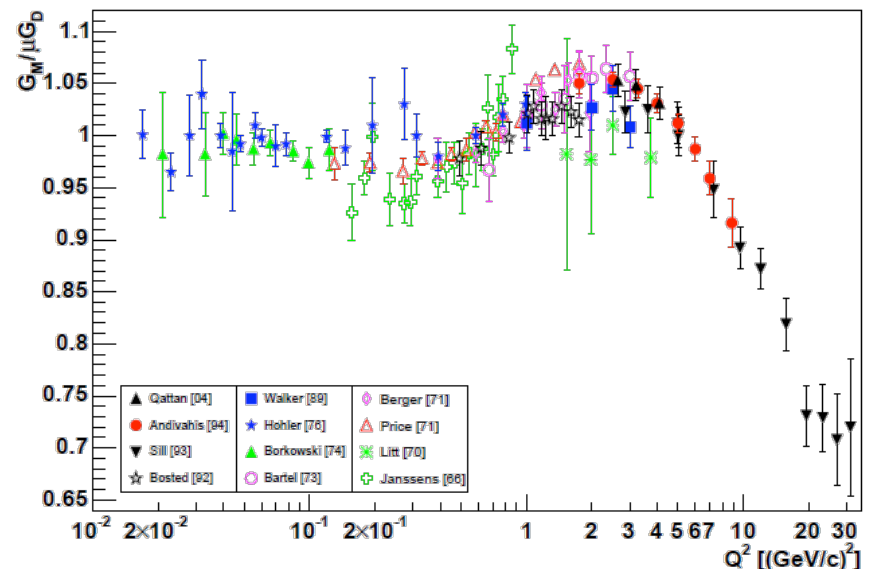
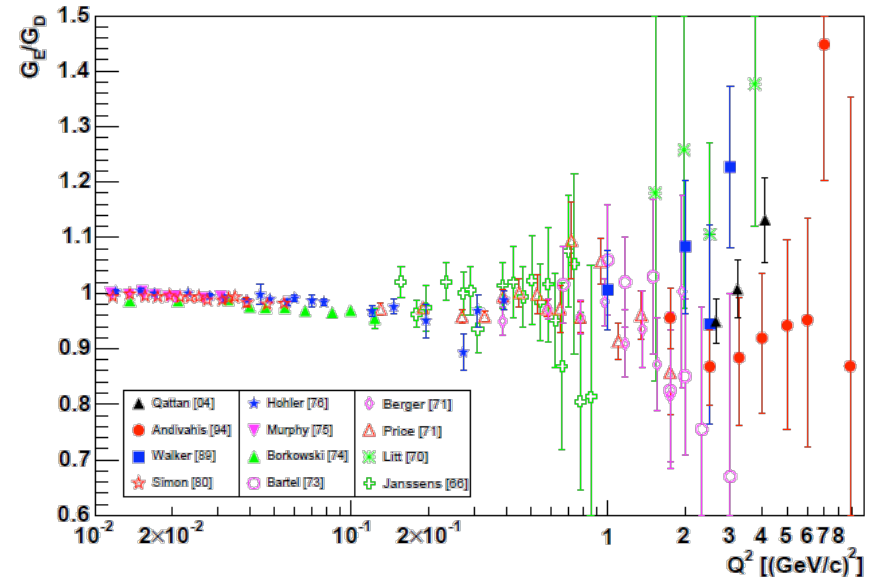
Electric:  $G_E \equiv F_1 - \tau F_2$

Magnetic:  $G_M \equiv F_1 + F_2$

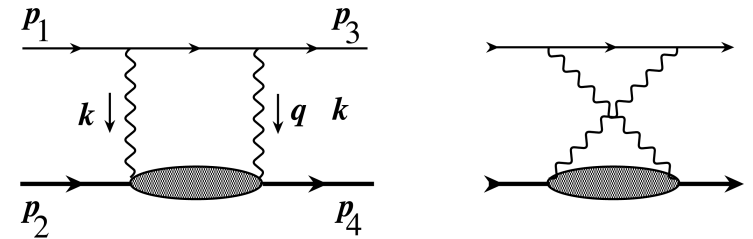
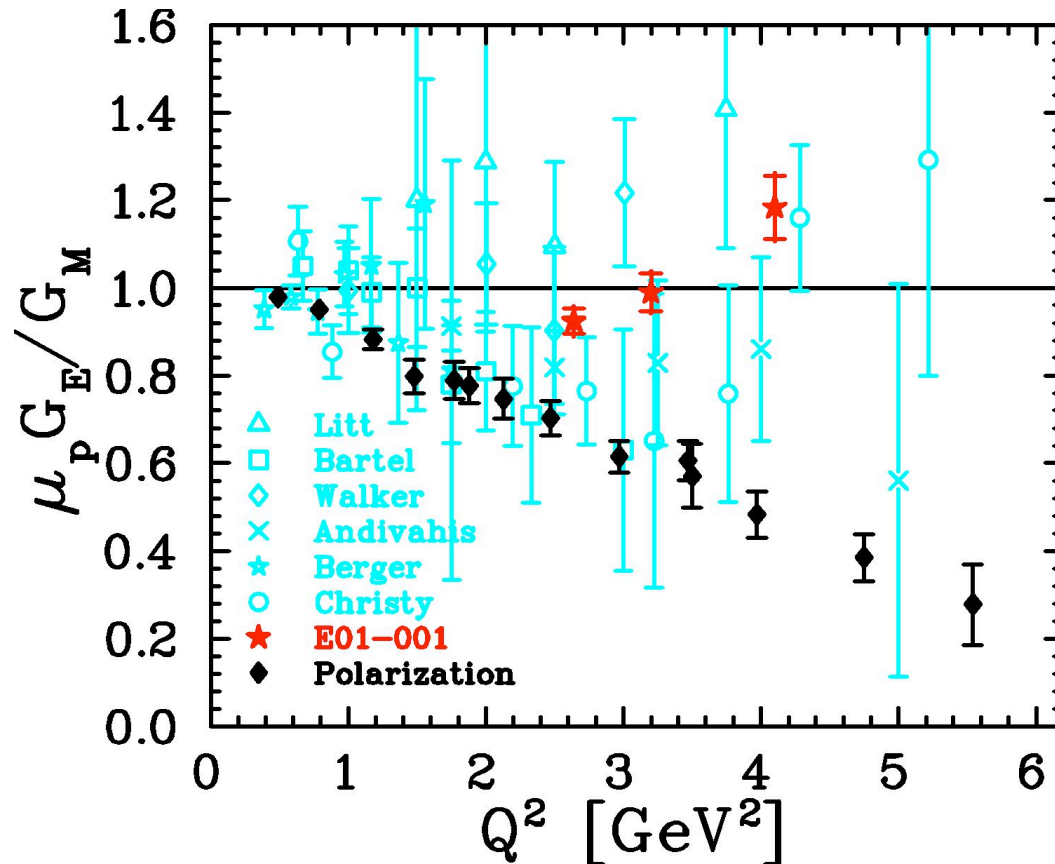
$$\frac{d\sigma}{d\Omega} = \sigma_{\text{Mott}} \frac{1}{1+\tau} [G_E^2 + \frac{\tau}{\epsilon} G_M^2]$$

- Rosenbluth separation  
Hard to determine  $G_E$  at large  $Q^2$ ,  
 $G_M$  at low  $Q^2$  (except at  $\theta_e \approx 90^\circ$ )
- Early experiments found  $\sim$   
dipole form ( $Q^2 < 2 \text{ GeV}^2$ ),  
naively corresponds to an  
exponential shape in space

$$G_D(Q^2) = \left(1 + \frac{Q^2}{0.71 \text{ GeV}^2}\right)^{-2} \quad \mu_P \frac{G_E}{G_M} = 1$$



# Rosenbluth vs. Polarimetry



*P.A.M. Guichon and M. Vanderhaeghen,  
PRL 91, 142303 (2003)*

*M.K. Jones, et al., PRL 84, 1398 (2000)*

*O. Gayou, et al., PRL 88, 092301 (2003)*

*I.A. Qattan, et al., PRL 94, 142301 (2005)*

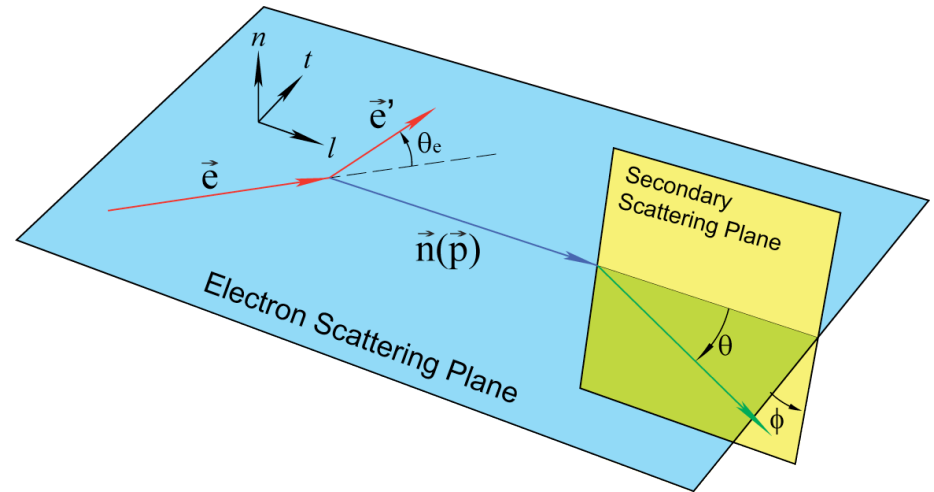
# Recoil Polarimetry

- Direct measurement of form factor ratios by measuring the ratio of the transferred polarization  $P_t$  and  $P_l$

$$l_0 P_t = -2\sqrt{\tau(1+\tau)} G_E G_M \tan \frac{\theta_e}{2}$$

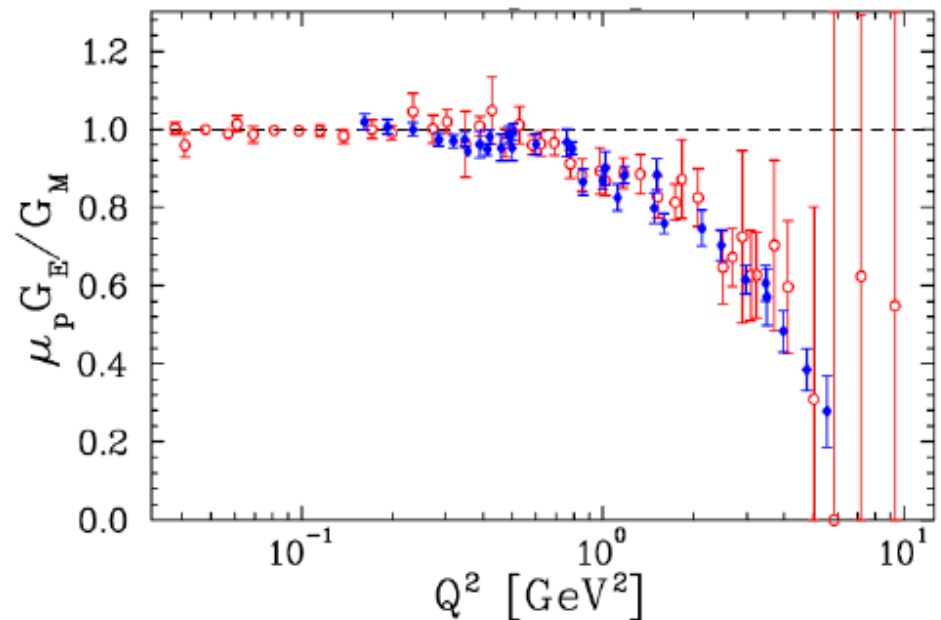
$$l_0 P_l = \frac{E_e + E_{e'}}{M} \sqrt{\tau(1+\tau)} G_M^2 \tan^2 \frac{\theta_e}{2}$$

$$\frac{G_E}{G_M} = -\frac{P_t}{P_l} \frac{(E_e + E_{e'})}{2M} \tan \frac{\theta_e}{2}$$



## Advantages:

- Only one measurement is needed for each  $Q^2$
- Much better precision than a cross section measurement
- Complementary to cross section measurements
- Discrepancy between Rosenbluth and polarized measurement, mostly explained by  $2-\gamma$  exchange



(J. Arrington, et al., Phys. Rev. C **76** 035205 (2007))

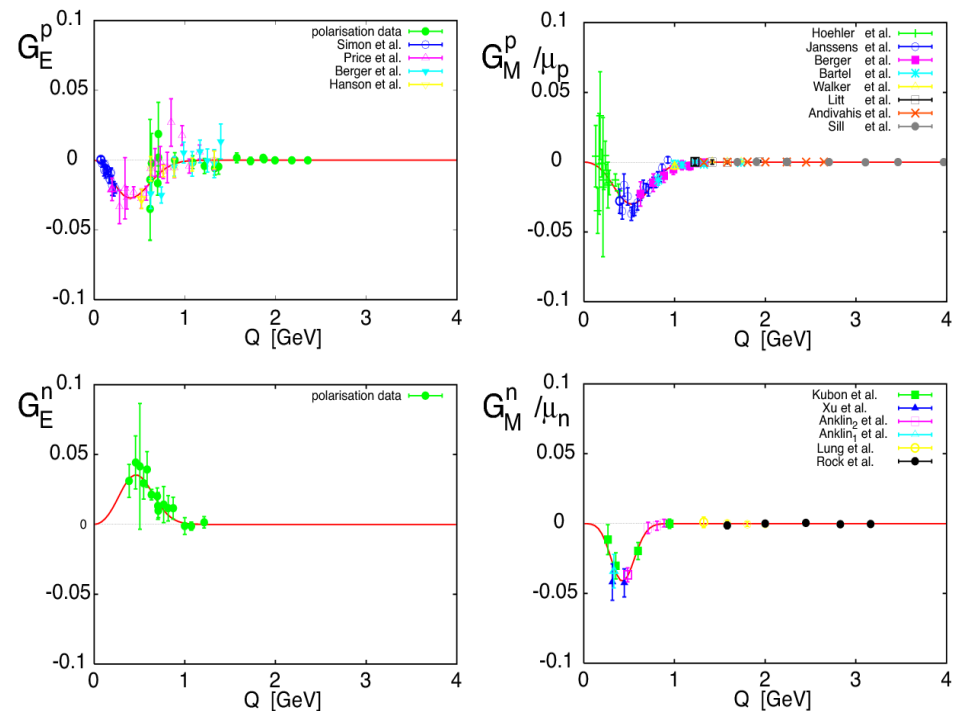
# Form Factors at Low $Q^2$

- Small  $Q^2 \rightarrow$  larger length scale, closely related to proton size

- Improved EMFFs:
  - Strange form factors through PV
  - Proton Zemach radius and hydrogen hyperfine splitting
  - Isoscalar and isovector form factors for Lattice QCD
  - Proton RMS radius

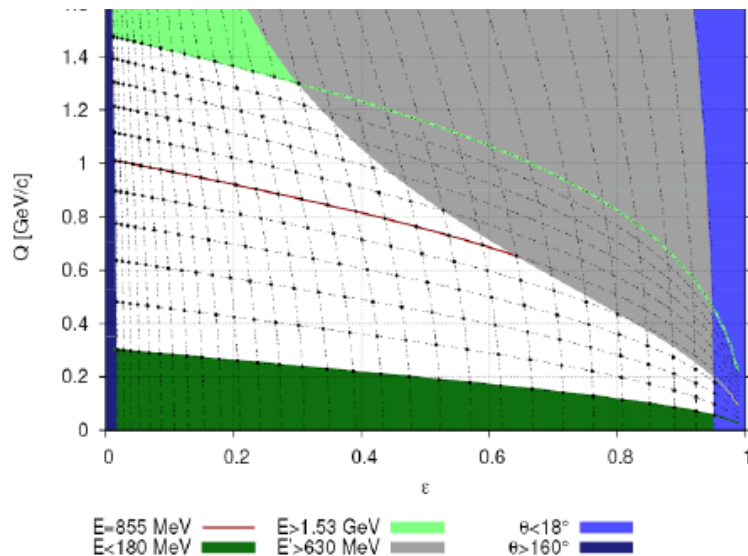
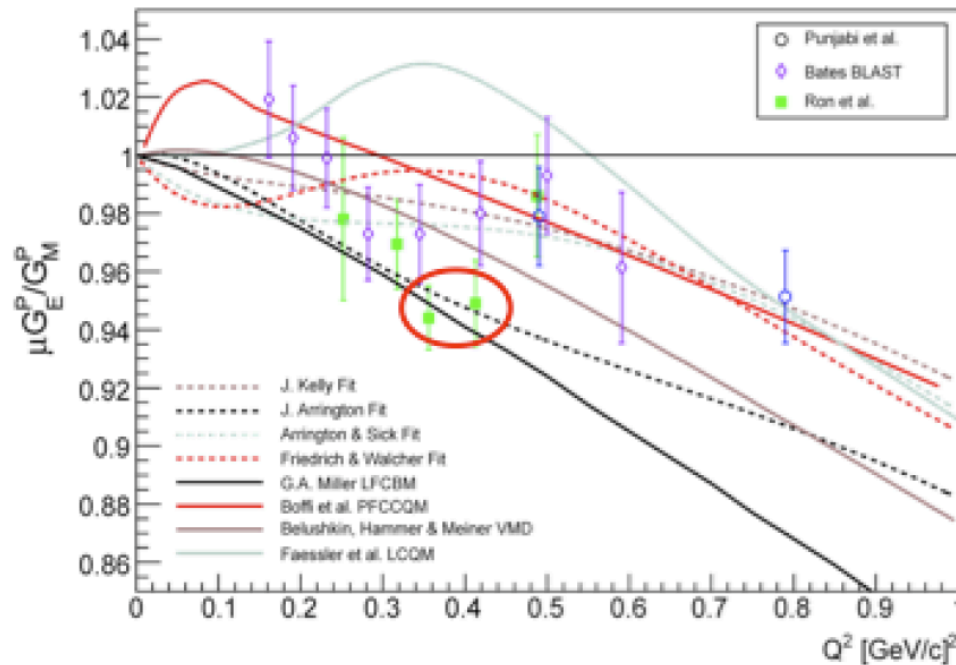
$$\langle r_{E,M}^2 \rangle = \frac{-6}{G_{E,M}(0)} \left[ \frac{d}{dQ^2} G_{E,M}(Q^2) \right]_{Q^2=0}$$

- 2003 – Fit by Friedrich & Walcher
  - Smooth dipole form + “bump & dip”
  - All four FFs exhibit similar structure at small momentum transfer ( $Q^2 \sim 0.25 \text{ GeV}^2$ )
  - Interpretation: effect of pion cloud



J. Friedrich and Th. Walcher, *Eur. Phys. J. A* **17**, 607 (2003)

# World Data at Low $Q^2$



- Bates **BLAST** result consistent with 1

Crawford et al., *Phys. Rev. Lett* 98 052301 (2007)

- Substantial deviation from unity observed in **LEDEX** – **smaller  $G_E$**   
G. Ron et al., *Phys. Rev. Lett* 99 202002 (2007)

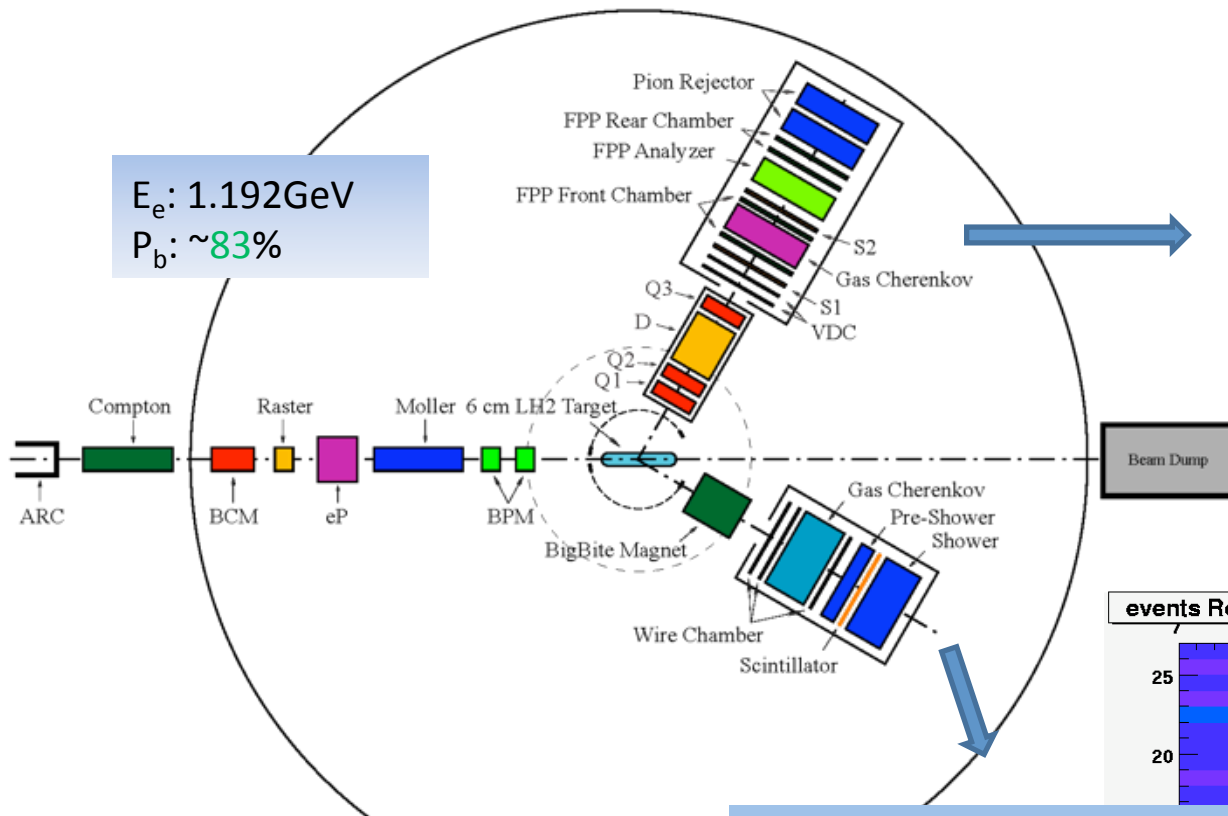
- Both data inconsistent with F&W fit

- Complementary to high precision XS measurement at **Mainz** ( $Q^2 \sim 0.003 - 1 \text{ GeV}^2$ )

- New dedicated experiment **JLAB E08-007**

# E08-007: Low $Q^2$ GEp/GMp

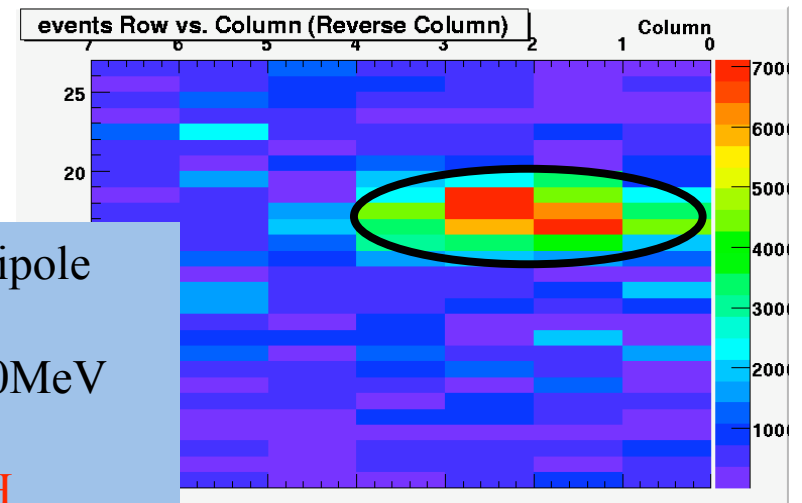
LHRS



$E_e$ : 1.192 GeV  
 $P_b$ : ~83%

- $\Delta p/p_0$ :  $\pm 4.5\%$  ,
- out-of-plane:  $\pm 60$  mrad
- in-plane:  $\pm 30$  mrad
- $\Delta\Omega$ : 6.7 msr
- QQDQ
- Dipole bending angle  $45^\circ$
- **VDC+FPP**
- $P_p$  : 0.55 ~ 0.93 GeV/c

BigBite

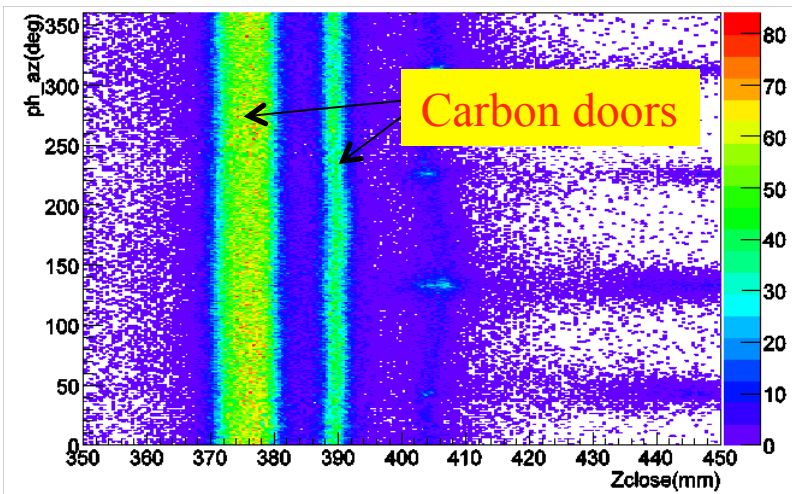
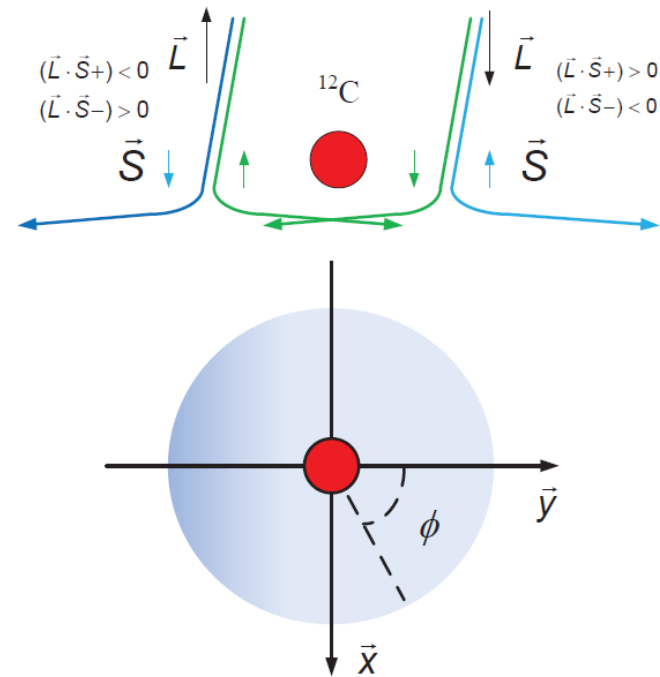
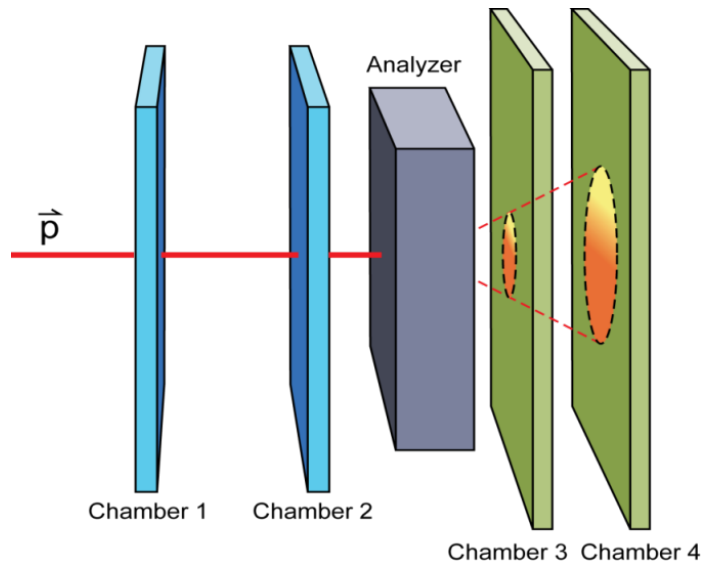


- A high precision ( $<1\%$ ) survey of the proton FF ratio
- 8  $Q^2$  data points: 0.3 ~ 0.7 (GeV/c) $^2$

- Non-focusing Dipole
- Big acceptance.
  - $\Delta p$ : 200-900 MeV
  - $\Delta\Omega$ : 96 msr
- PS + Scint. + **SH**



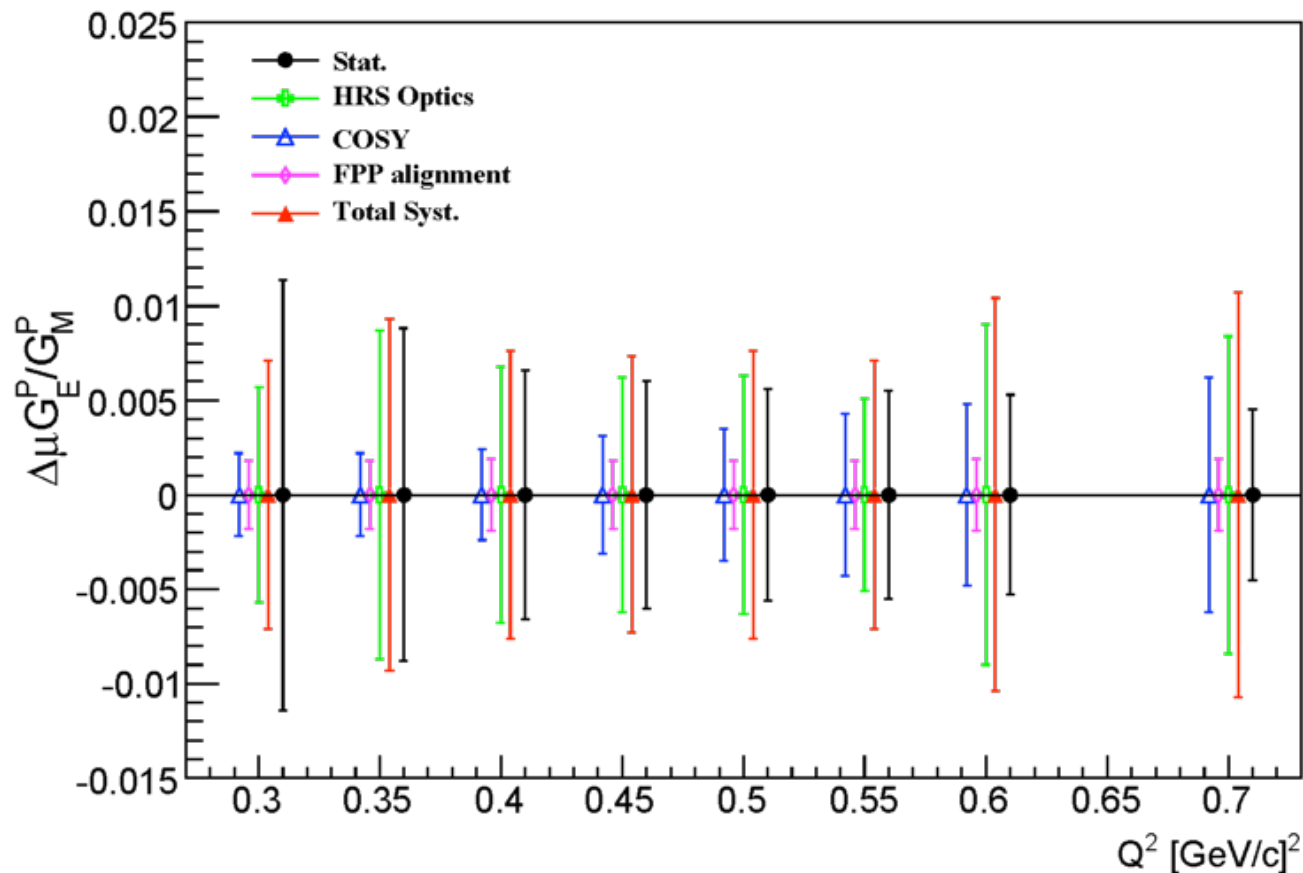
# Focal Plane Polarimeter (FPP)



- Left-right asymmetry gives the vertical component while the up-down asymmetry gives the horizontal component
- Need well determined scattering azimuthal angle,  $\phi_{fpp}$ , chamber alignment checked with straight through data

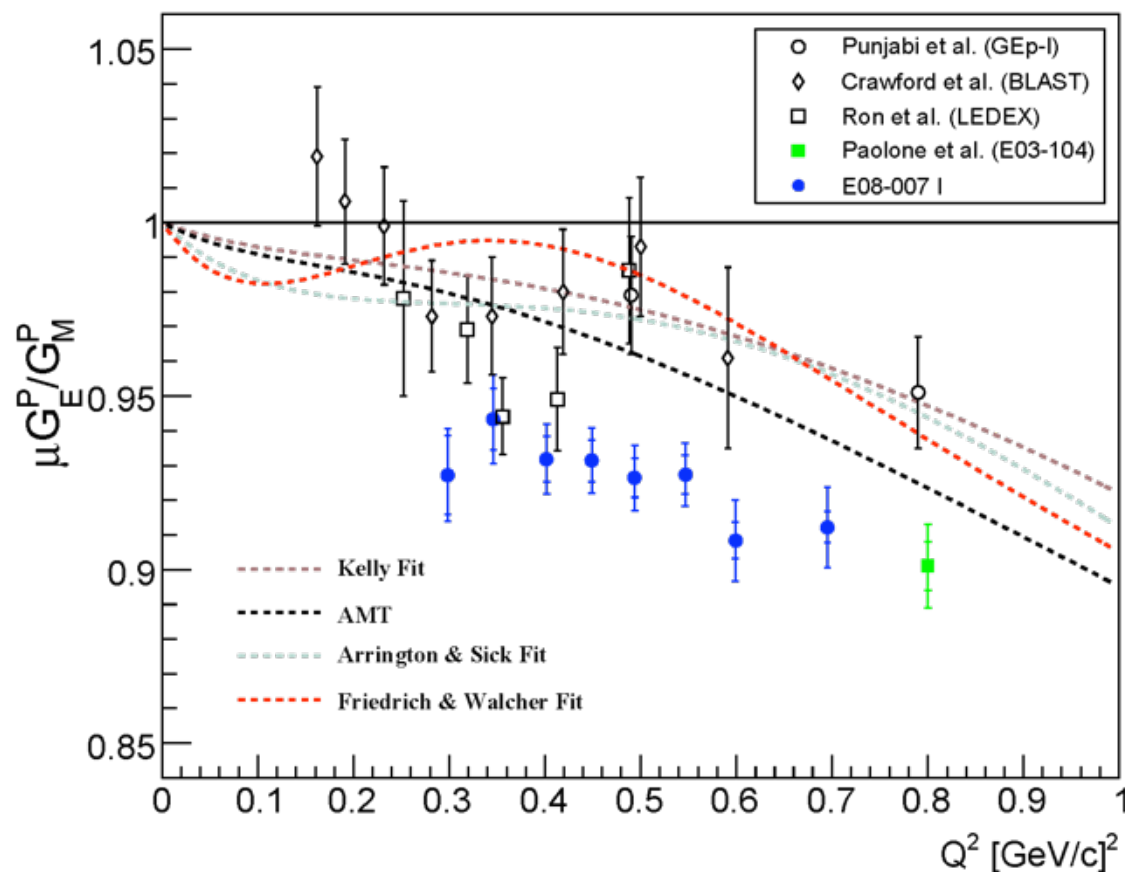
# Systematic Budget

- Spin transport: OPTICS and COSY---major uncertainty (0.7 ~ 1.2 %)
- Others negligible: FPP alignment, Al end cap contamination, VDC reconstruction, spectrometer settings, beam energy, charge asymmetry, pion contamination, etc.



# E08-007 Final Results

X-H. Zhan, MIT Ph.D. thesis (2010), to be submitted

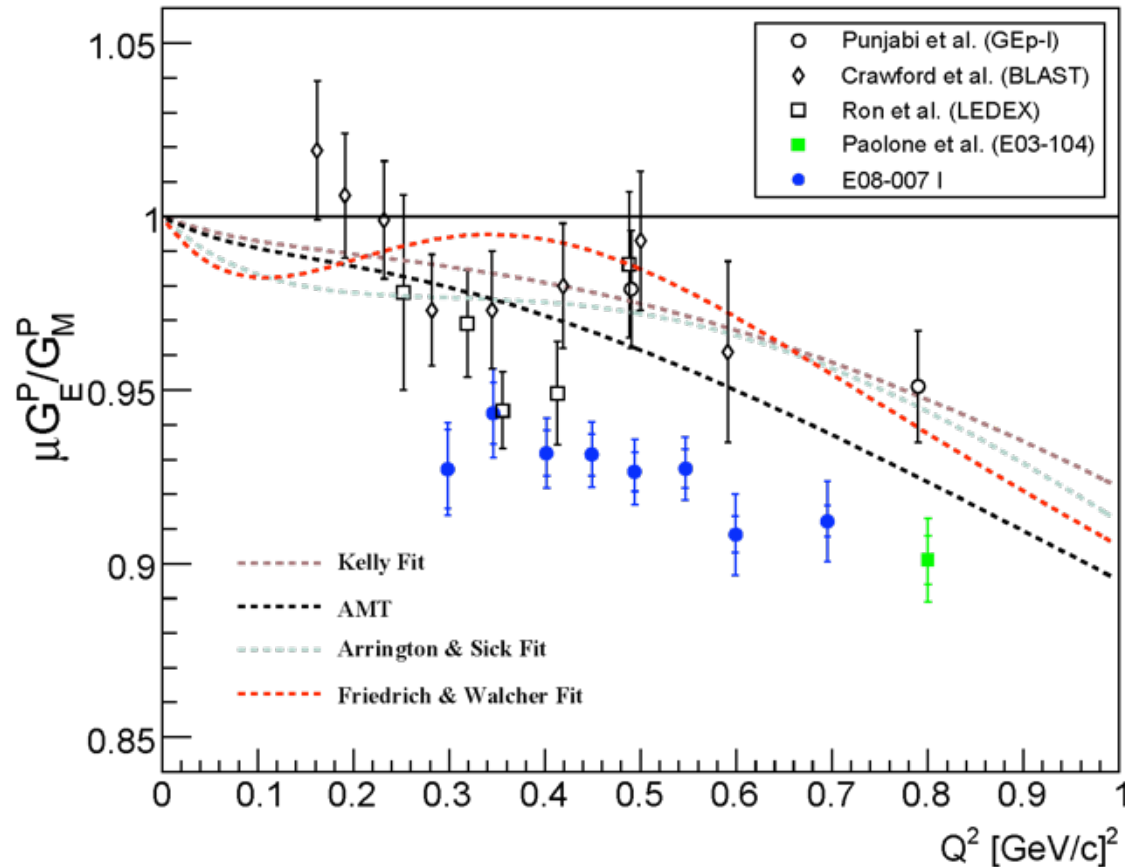


- Agreement with independent analysis of *Paolone et al.* at 0.8  $\text{GeV}^2$
- Slow decrease with  $Q^2$ . A few percent below previous data, fits
- Suggests that  $R = \mu_p G_{Ep}/G_{Mp} < 1$  even at low  $Q^2$
- No obvious indication of "Structure", inconsistent with F&W fit

AMT - W. Melnitchouk J. Arrington and J.A. Tjon, Phys. Rev. C **76**,035205 (2007)  
 J. Arrington and I. Sick. Phys. Rev. C, **76**, 035201 (2007)

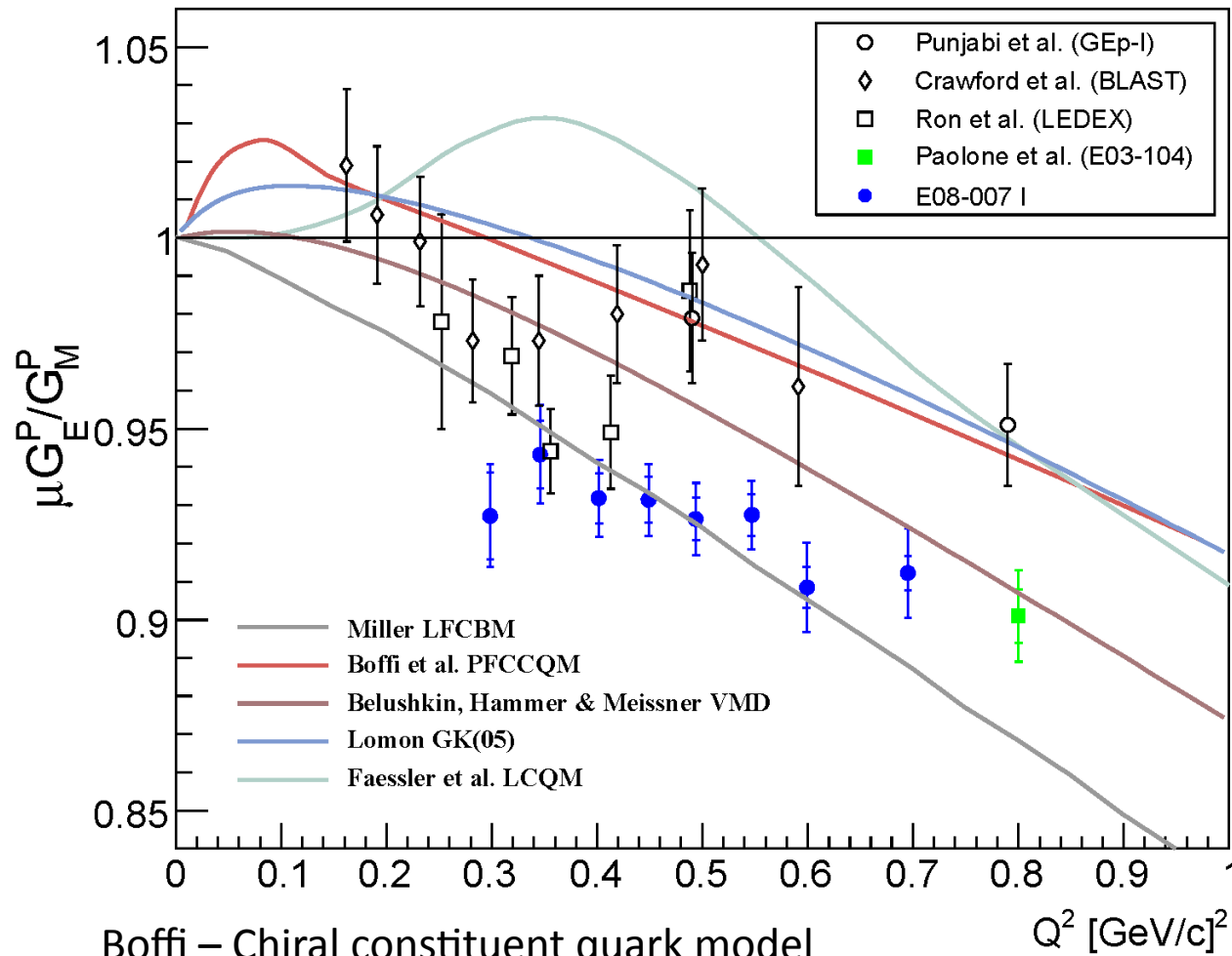
# Comparison to Previous Data

X-H. Zhan, MIT Ph.D. thesis (2010), to be submitted



- Disagreement with  $GE_p$ -I, BLAST, LEDEX
- Agreement with independent analysis of *Paolone et al.* at 0.8  $GeV^2$
- LEDEX re-analysis consistent with new data
- Indications of cut dependence in  $GE_p$ -I point
- No current plans to re-examine BLAST data

# Comparison with Models



◆ No model consistent with data!

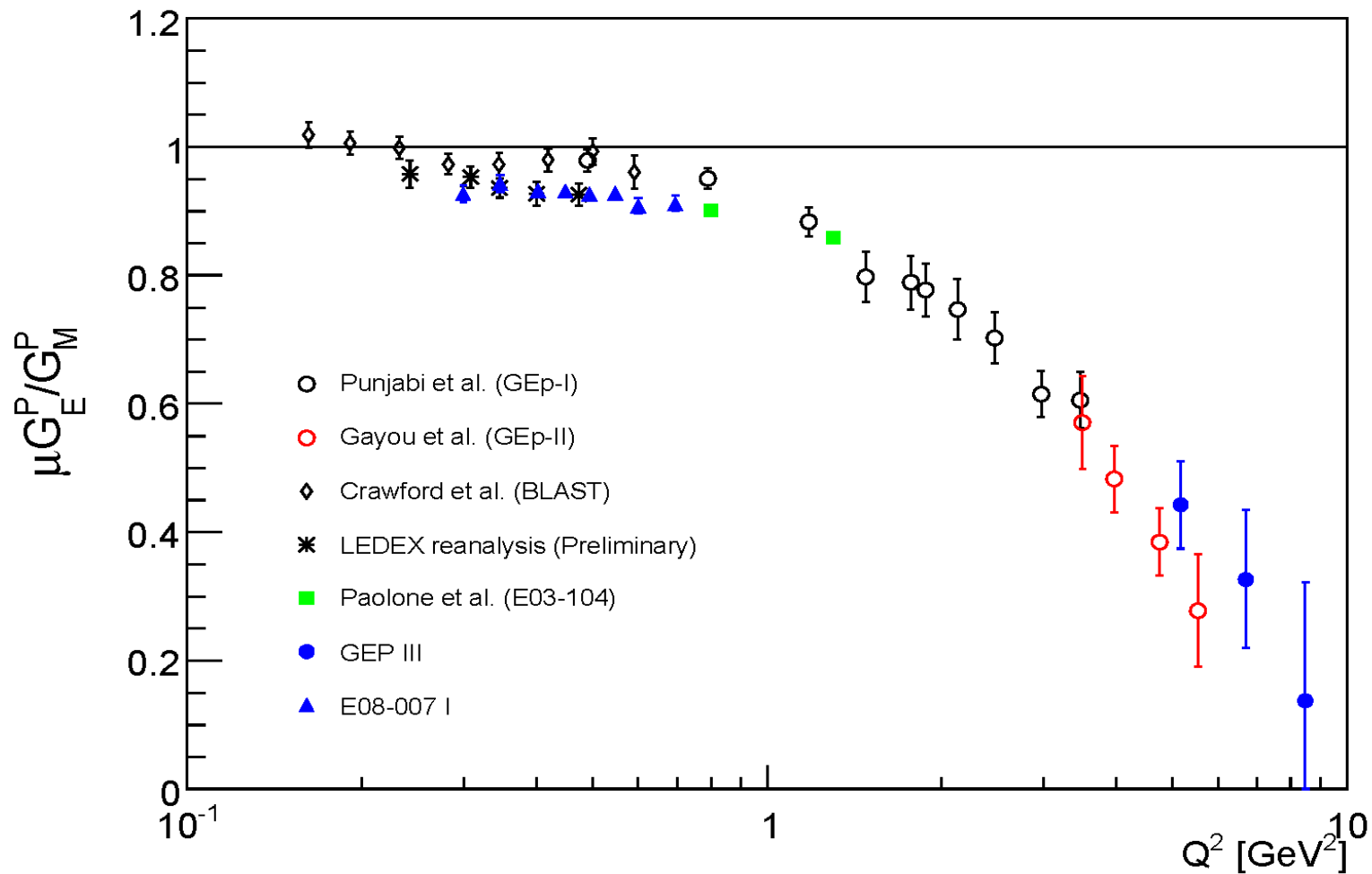
Boffi – Chiral constituent quark model

Faessler – Lorentz covariant chiral quark model

Belushkin, Lomon – VMD

Miller – Light front cloudy bag model

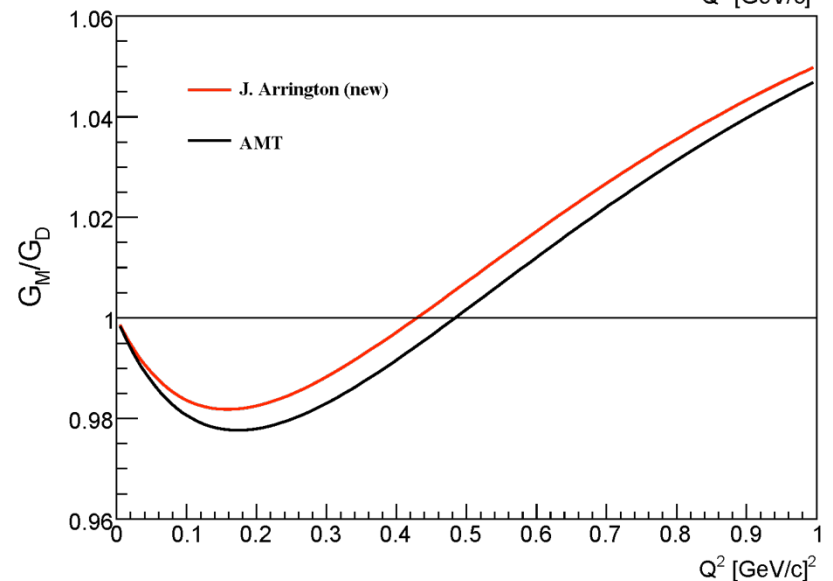
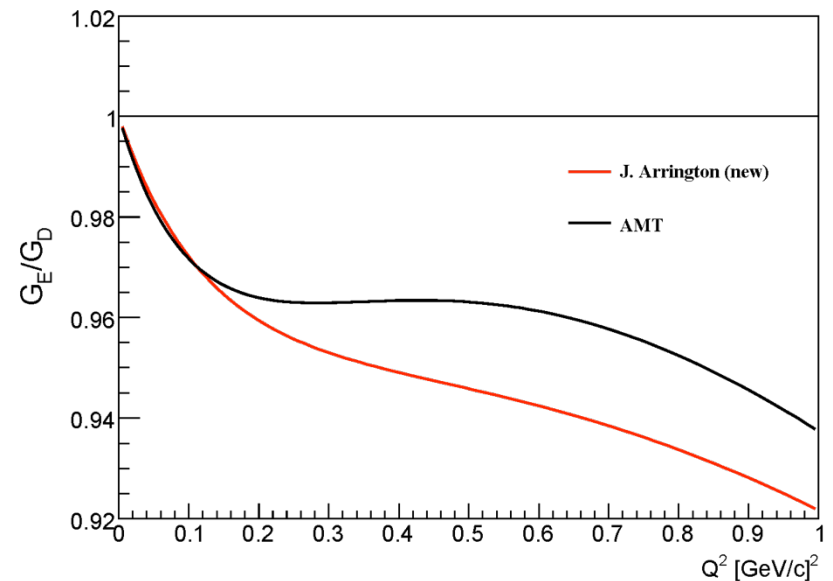
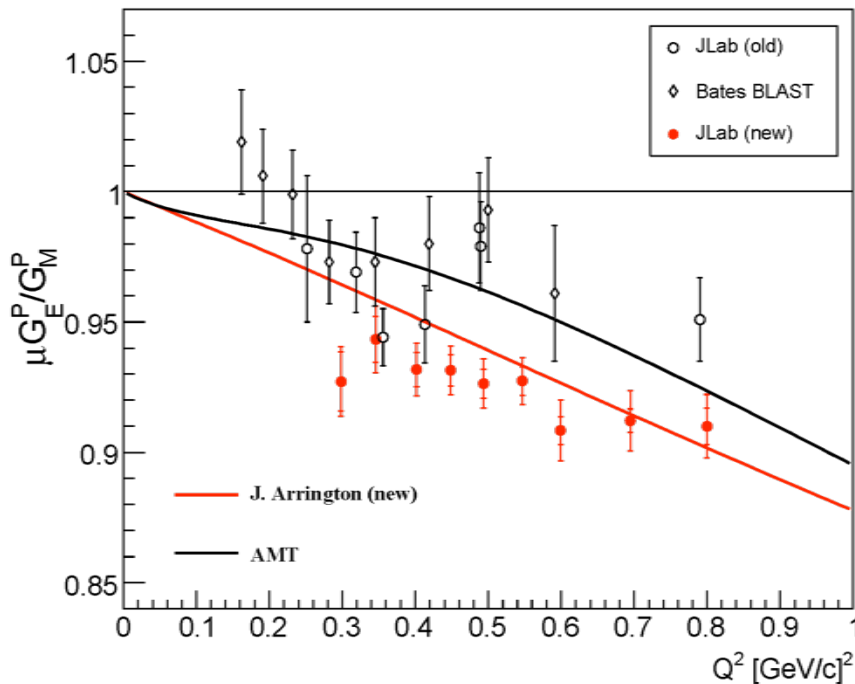
# World Data with Polarization



Note: LEDEX re-analysis  
Examination of Gep-I, Gep-II

# Global Fits

- Combined global fits (**John Arrington**).
- AMT fit (black) : include all previous data with TPE correction
- New fit (**red**) : same procedure, includes new data, remove lowest point of **G<sub>Ep-I</sub>**, (and highest point of **LEDEX**)
- Preliminary fits suggest lower  $G_E$  (~2%), higher  $G_M$  (0.5%)



# Impact I – Strangeness Form Factors by PV

$$A_{PV} = -\frac{G_F Q^2}{4\pi\alpha\sqrt{2}} \left[ (1 - 4\sin^2\theta_W) - \frac{\varepsilon G_{Ep}(G_{En} + G_{Es}) + \tau G_{Mp}(G_{Mn} + G_{Ms})}{\varepsilon(G_{Ep})^2 + \tau(G_{Mp})^2} - \frac{(1 - 4\sin^2\theta_W)\varepsilon' G_{Mp} G_A^Z}{\varepsilon(G_{Ep})^2 + \tau(G_{Mp})^2} \right]$$

- Asymmetry arises from the interference between EM and neutral weak currents

$$\sigma \propto |\mathcal{M}_\gamma + \mathcal{M}_Z|^2$$

$$\mathcal{M}^R = \mathcal{M}_\gamma + \mathcal{M}_Z^R,$$

$$\mathcal{M}^L = \mathcal{M}_\gamma + \mathcal{M}_Z^L.$$

$$A_{PV} = \frac{\sigma_R - \sigma_L}{\sigma_R + \sigma_L} = \frac{|\mathcal{M}^R|^2 - |\mathcal{M}^L|^2}{|\mathcal{M}^R|^2 + |\mathcal{M}^L|^2}$$

- Rely on knowledge of EMFFs
- With New FF parameterization, HAPPEX III results shift  $\sim 0.5\sigma$

$Q^2$	$\Delta A$	$\Delta A/\sigma$	$\Delta A/A$	Exp.
0.38	-0.178	0.42	1.6%	G0 FWD
0.56	-0.347	0.50	1.6%	G0 FWD
1.0	-0.414	0.30	0.8%	G0 FWD
0.50	-0.299	0.50	1.7%	HAPPEX III
0.231	+0.038	0.12	0.2%	G0 BCK
0.65	0.142	0.14	0.3%	G0 BCK

Table: Difference in the extracted asymmetries.



# Impact II – Proton Zemach Radius

## Hyperfine splitting of hydrogen ground state

$$E_{hfs} = (1 + \Delta_{QED} + \Delta_{hvp}^p + \Delta_{\mu vp}^p + \Delta_{weak}^p + \Delta_S) E_F^p$$

$$\Delta_S = \Delta_Z + \Delta_R^p + \Delta_{pol}, \quad \Delta_Z = -2\alpha Z \frac{m_e m_p}{m_e + m_p} r_Z$$

- Leading theoretical uncertainty

$$r_Z = -\frac{4}{\pi} \int_0^\infty \frac{dQ}{Q^2} [G_E(Q^2) G_M(Q^2) / (1 + \kappa_p) - 1]$$

- FFs at Low  $Q^2$  ( $< 1 \text{ GeV}^2$ ) accounts for  $> 70\%$  of  $r_Z$ , and also dominate the uncertainty

Quantity	value (ppm)	uncertainty (ppm)
$(E_{hfs}(e^-p)/E_F^p) - 1$	1 103.48	0.01
$\Delta_{QED}$	1 136.19	0.00
$\Delta_{\mu vp}^p + \Delta_{hvp}^p + \Delta_{weak}^p$	0.14	
$\Delta_Z$ (using [31])	-41.43	0.44
$\Delta_R^p$ (using [31])	5.85	0.07
$\Delta_{pol}$ (this work, using [31])	1.88	0.64
Total	1102.63	0.78
Deficit	0.85	0.78

FFs	$r_z$ (fm)	$\Delta z$	year
Dipole	1.025	-39.29	-
FW	1.049	-40.22	2003
Kelly	1.069	-40.99	2004
AS	1.091	-41.85	2007
AMT	1.080	-41.43	2007
New fit	1.075	-41.21	2009

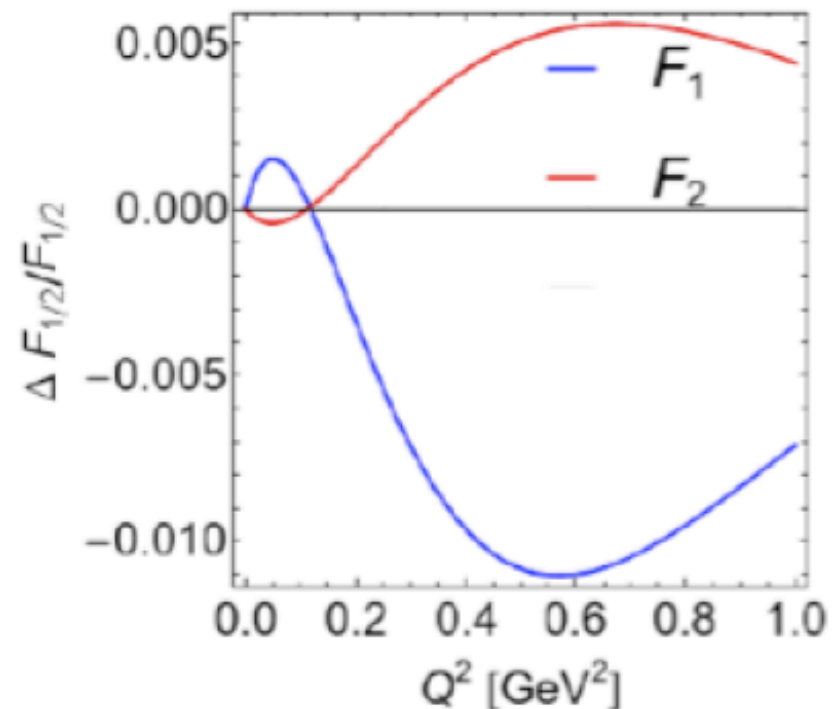
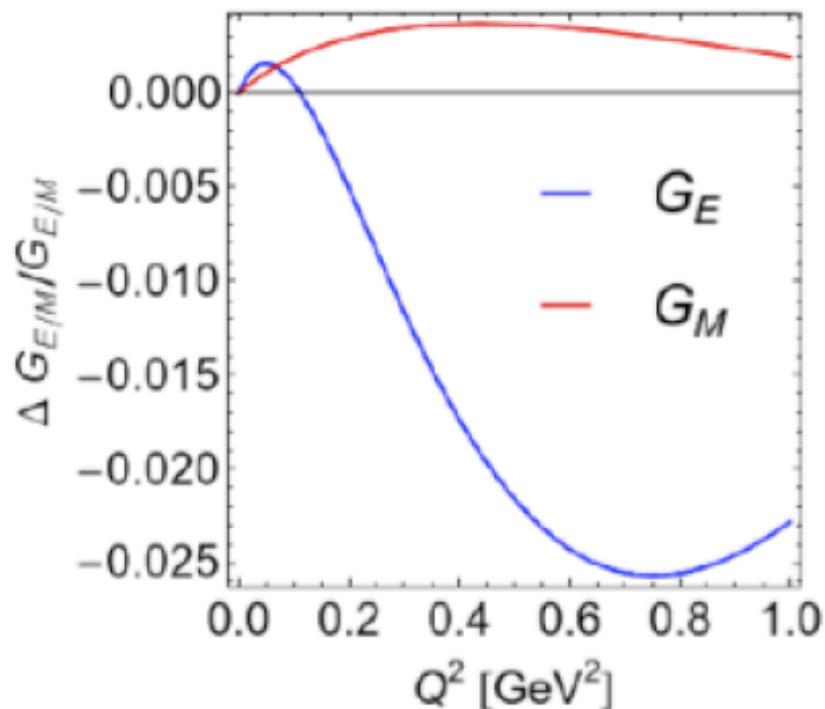
Carlson, Nazaryan, and Griffioen, *arXiv:0805.2603v1* (2009)

# Impact III – Isoscalar & Isovector FFs

- Isoscalar & Isovector FFs (important for Lattice QCD):

$$F_J^S = \frac{1}{2}(F_J^p + F_J^n) \quad F_J^V = \frac{1}{2}(F_J^p - F_J^n)$$

- Change in FFs by using new parameterization vs. old parameterization



# Impact IV - e-d Elastic Cross Section

$$\frac{d\sigma}{d\Omega} = \sigma_{Mott} * \frac{E'}{E} * [A(Q^2) + B(Q^2)\tan^2(\theta/2)]$$

Where

$$A(Q^2) = G_C^2(Q^2) + \frac{8}{9}\eta^2 G_Q^2(Q^2) + \frac{2}{3}\eta G_M^2(Q^2)$$

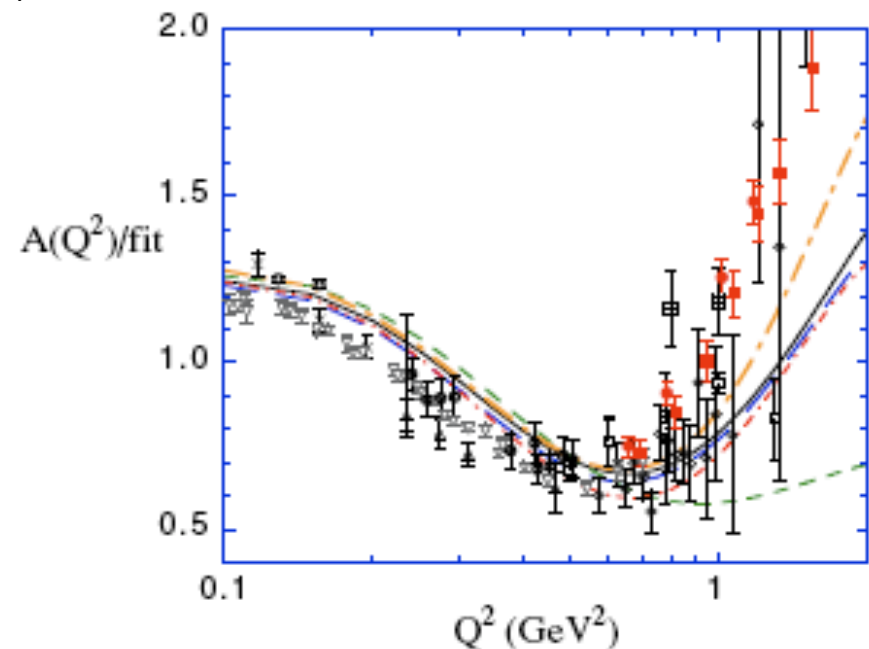
$$B(Q^2) = \frac{4}{3}\eta(1+\eta)G_M^2(Q^2) \quad \text{Small except at } \theta \approx 180^\circ$$

and

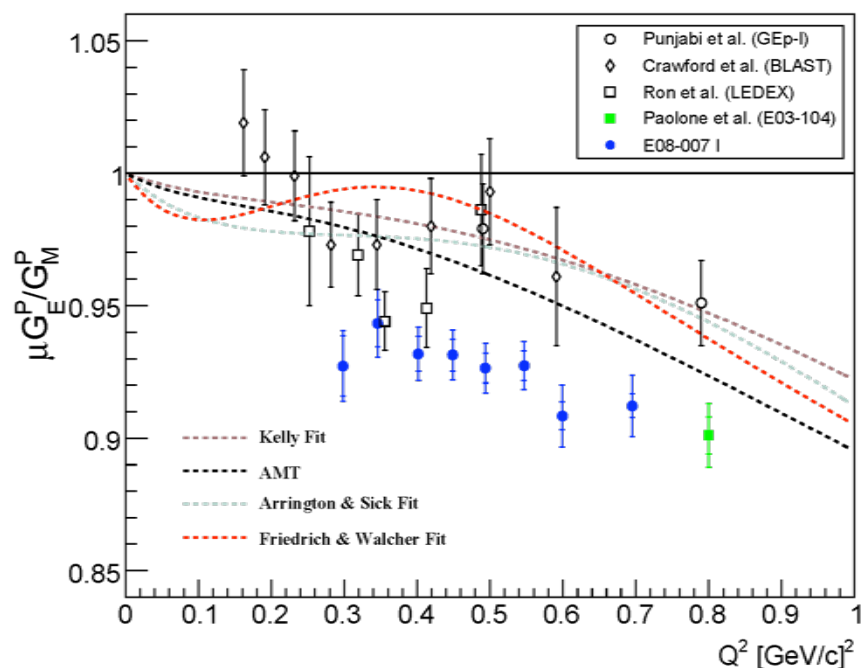
$$\eta = Q^2 / 4m_d^2 \quad \text{Very small at low } Q$$

At low  $Q^2$ , elastic e-d cross section depends mainly on  $GE_p$

New results change model cross sections by a few percent!



# Impact V – Proton RMS Radius



Increased interest in proton radius with new very precise muonic hydrogen data (significantly lower:  $\sim 0.840$  fm)

**New Global fit** with TPE (Arrington):

$$\langle r_p^2 \rangle = 0.873(14) \text{ fm}$$

3% below previous value (Sick):  $0.897(18)$  fm

CODATA(2006) value:  $0.877(7)$  fm

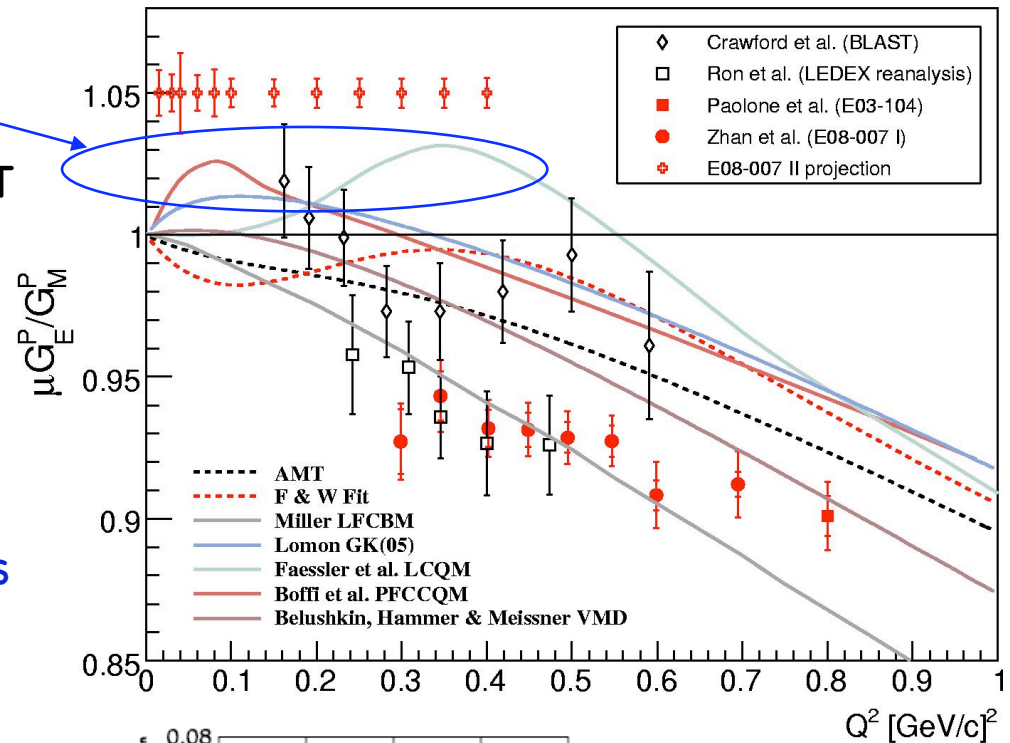
If  $R=1$  as  $Q^2$  approaches zero, yields

**0.015 fm change in charge radius**

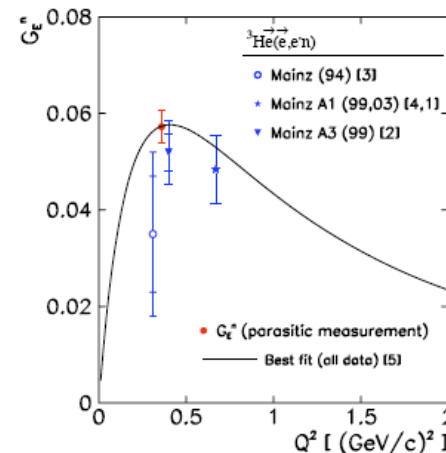
# Near Future Outlook

LEDEX re-analysis almost complete; consistent with X-H. Zhan et al.

- Phase-II (pol. Target - 2012)
  - Extract R down to  $Q^2=0.015$
  - Good overlap with Phase-I, BLAST
  - First precise extraction of magnetic radius
  - Linear approach to  $Q^2=0$  ?
    - ~3% smaller magnetic radius
    - No region where magnetization, charge are simply sum of quarks
    - Disagreement with muonic hydrogen radius?



- New  $G_{En}$  measurement
  - One data point was measured parasitically in Hall A during experiment E05-102.
  - Data in analysis



# Summary

- Nucleon FFs are fundamental quantities describing the nucleon internal structure
- A new high precision measurement was conducted in Jefferson Lab Hall A at low  $Q^2$ , new data strongly deviate from unity, systematically lower than previous world data
- While adding further constraints on theoretical models, the new high precision data also impact determination of other physics quantities: proton Zemach radius, strange form factor through PV, proton RMS radius etc.
- Near-future experiment with polarized target accessing very low  $Q^2$ :
  - ✓ Precisely determine magnetic radius for the first time
  - ✓ Important for precisely determining proton RMS radius, check possible disagreement with determination from muonic hydrogen data
  - ✓ Help to settle disagreement between blast (pol. target) and JLAB (recoil polarization) data

# E08-007 Collaboration

- ◆ Argonne National lab
- ◆ Jefferson Lab
- ◆ Rutgers University
- ◆ St. Mary's University
- ◆ Tel Aviv University
- ◆ UVa
- ◆ CEN Saclay
- ◆ Christopher Newport University
- ◆ College of William & Mary
- ◆ Duke University
- ◆ Florida International University
- ◆ Institut de Physique Nuclaire d'Orsay
- ◆ Kent State University
- ◆ MIT
- ◆ Norfolk State University
- ◆ Nuclear Research Center Negev
- ◆ Old Dominion University
- ◆ Pacific Northwest National Lab
- ◆ Randolph-Macon College
- ◆ Seoul National University
- ◆ Temple University
- ◆ Universite Blaise Pascal
- ◆ University of Glasgow
- ◆ University of Maryland
- ◆ University of New Hampshire
- ◆ University of Regina
- ◆ University of South Carolina

**Thank you!**



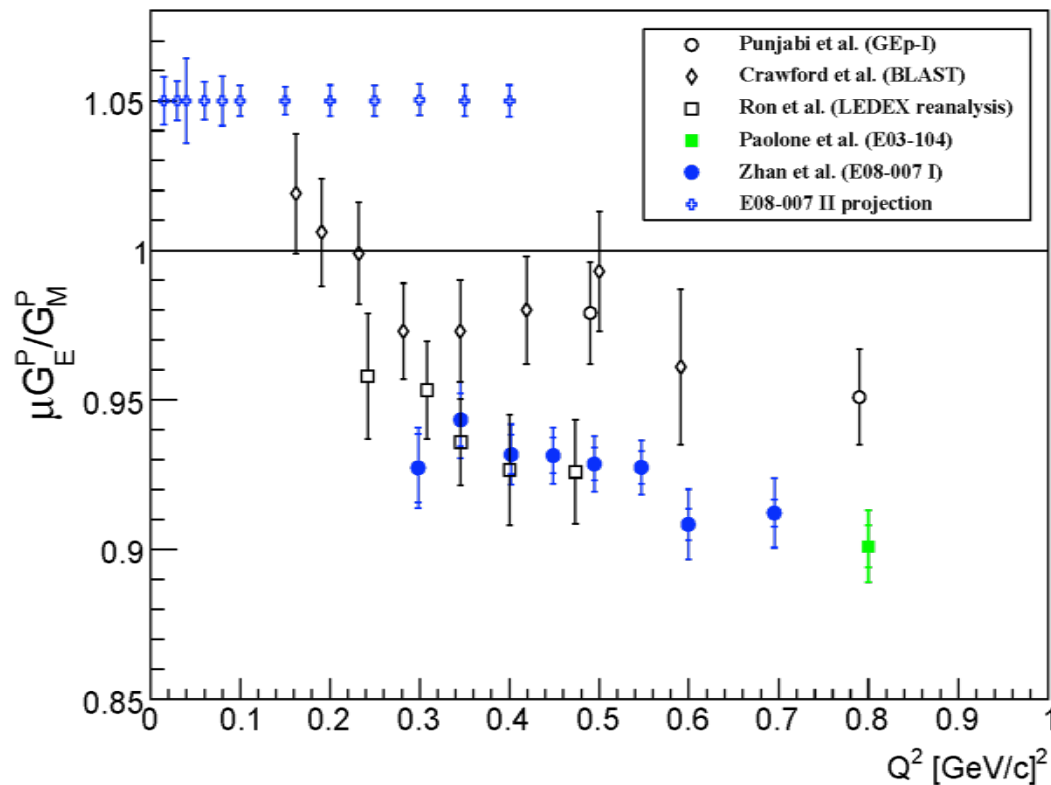
## Acknowledgements

J. Arrington, D. Higinbotham, J. Glister, R. Gilman, S.  
Gilad, E. Piasetzky, M. Paolone, G. Ron, A. Sarty, S.  
Strauch and the entire E08-007 collaboration  
&  
Jefferson Lab Hall A Collaboration

# Back up slides

## Future Outlook

- E08007 analysis finalized.
- Publication in preparation.
- Updated paper for LEDEX (G. Ron *et al.*) in preparation.



- Second half of the experiment (DSA) is tentatively scheduled in early 2012

$$A_{\text{phys}} = \frac{v_z \cos \theta' G_M^2 + v_x \sin \theta' \cos \varphi' G_E G_M}{(\varepsilon G_{Ep}^2 + \tau G_{Mp}^2) [\varepsilon (1 + \tau)]}$$

- Opportunity to see the FFR behavior at even lower  $Q^2$  (0.015-0.4  $\text{GeV}^2$ ) region.
- Third independent measurement, direct comparison with **BLAST**, examine any unknown systematic errors for previous measurements.
- Challenges: Solid polarized proton target & effect of target field to septum magnets.

## Elastic Events Selection

- **HRS acceptance cut:**

- out of plane:  $\pm 60$  mr
- in plane:  $\pm 30$  mr
- momentum:  $\pm 0.04$  ( $dp/p_0$ )
- reaction vertex cut

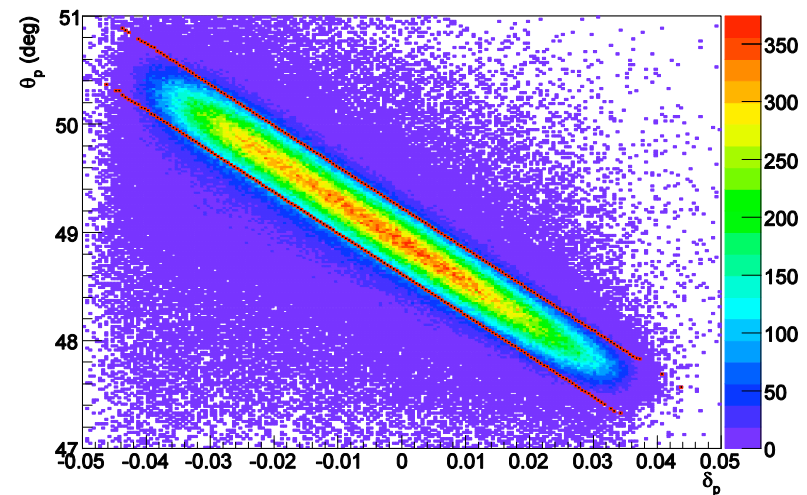
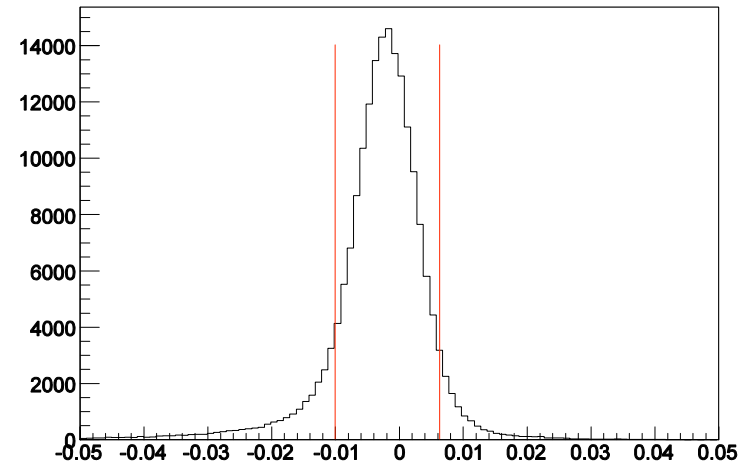
- **FPP cuts:**

- scattering angle  $\theta_{\text{fpp}}$   $5^\circ \sim 25^\circ$
- reaction vertex (carbon door)
- conetest cut

- **Other cuts:**

- Coin. Timing cut
- Coin. event type (trigger)
- single track event
- dpkin (**proton angle vs. momentum**)

proton dpkin



## Focal Plane Asymmetry

- Detection probability at focal plane with azimuthally angle  $\phi_{fpp}$

$$f^{\pm} = \frac{1}{2\pi} \xi [1 \pm A_y(\theta_{fpp})(P_x^{fpp} \sin(\phi_{fpp}) - P_y^{fpp} \cos(\phi_{fpp}))]$$

- Helicity difference:

$$f^{diff} = f^{+} - f^{-} \approx \frac{1}{\pi} [A_y(P_x^{fpp} \sin(\phi_{fpp}) - P_y^{fpp} \cos(\phi_{fpp}))] = C \cos(\phi + \delta)$$

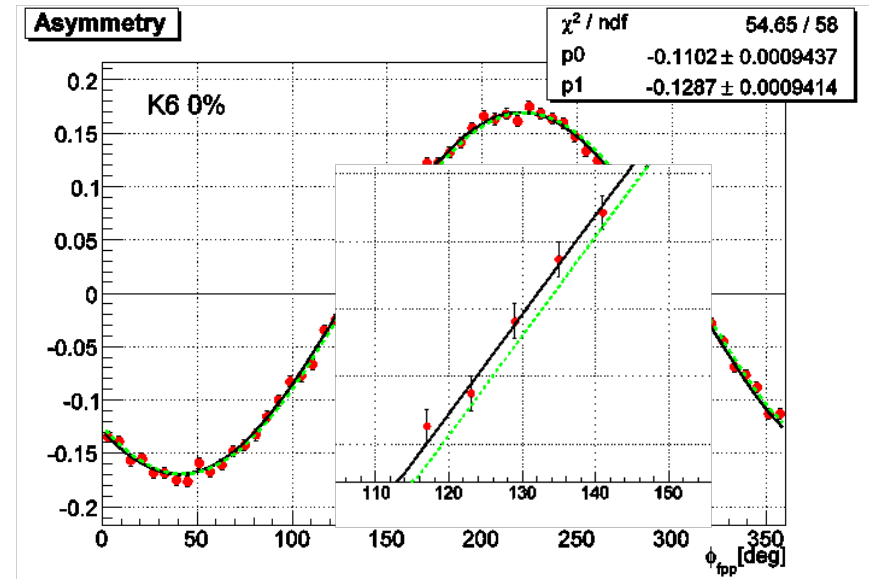
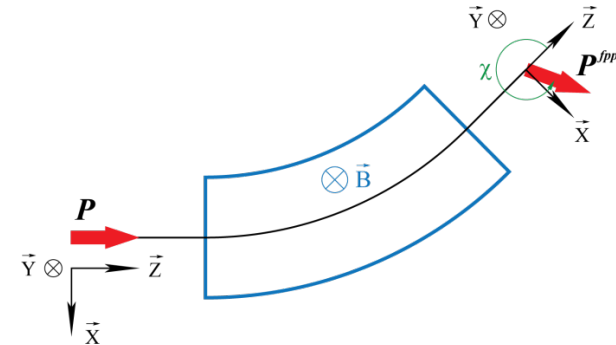
$$C = \frac{1}{\pi} A_y \sqrt{(P_x^{fpp})^2 + (P_y^{fpp})^2}$$

$$\tan \delta = \frac{P_y^{fpp}}{P_x^{fpp}}$$

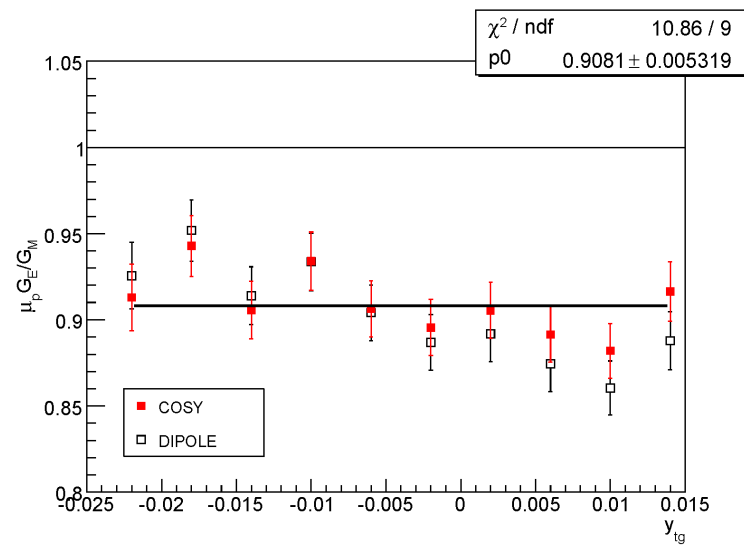
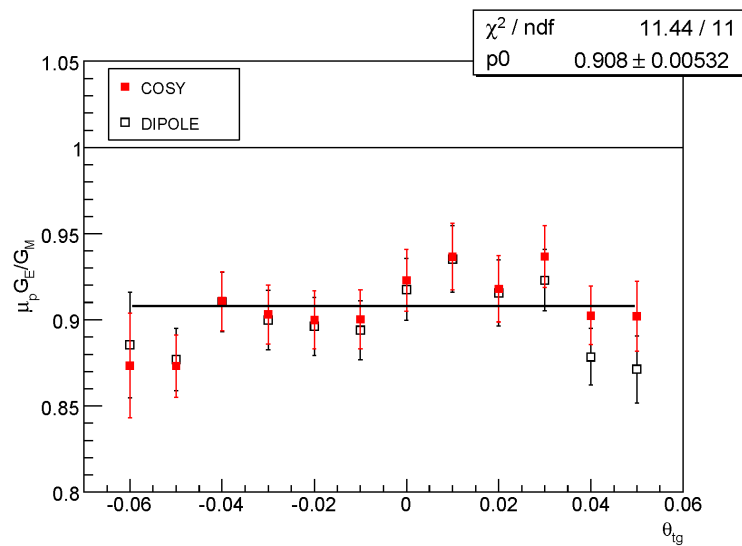
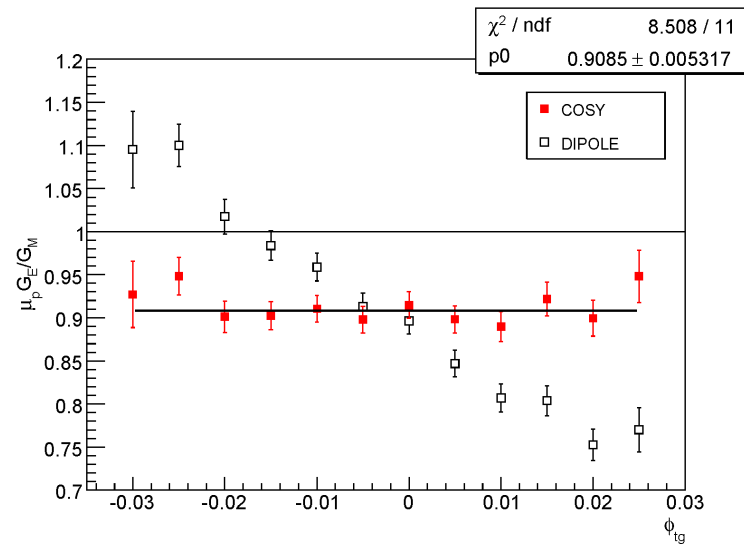
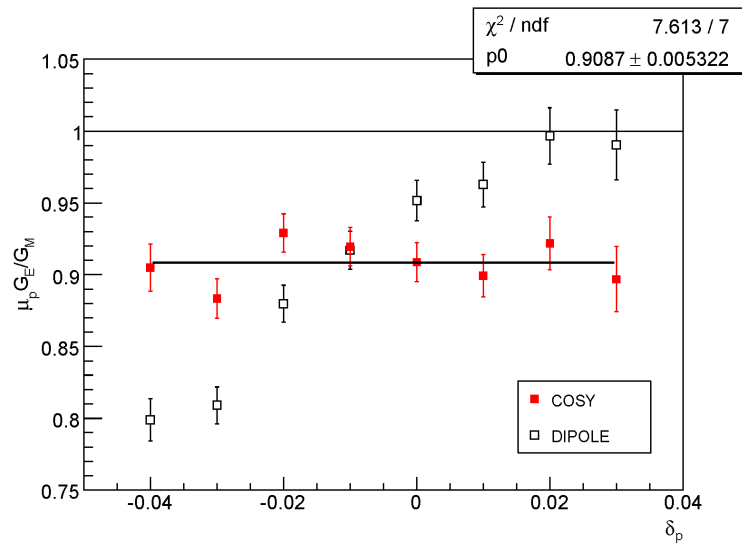
- By dipole approximation:

$$R = \mu_p \frac{G_E}{G_M} \approx \sin \chi \frac{P_x^{fpp}}{P_y^{fpp}} \times K$$

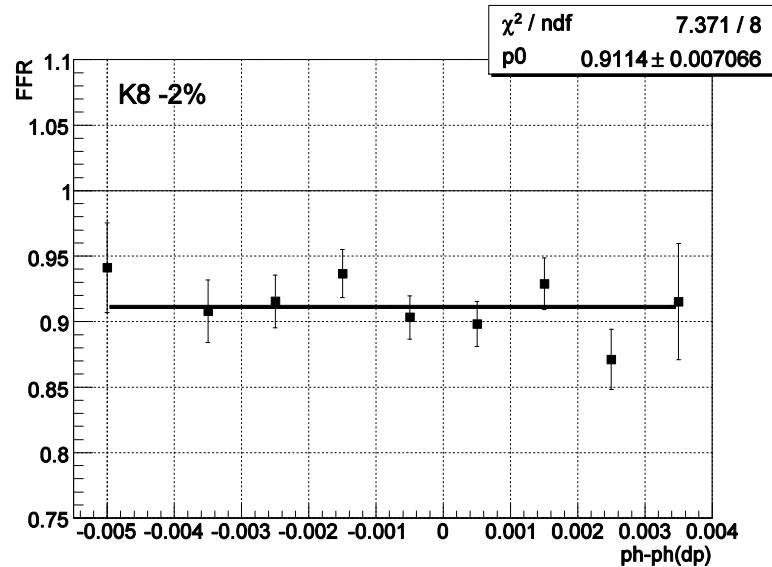
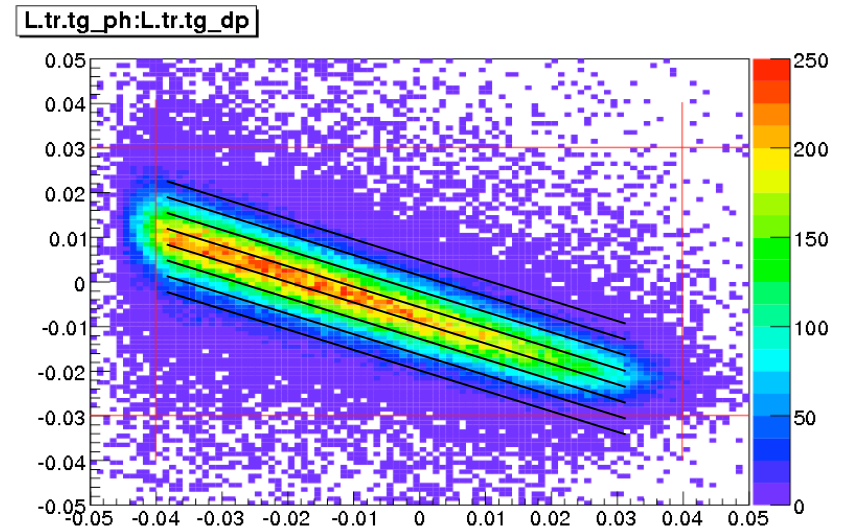
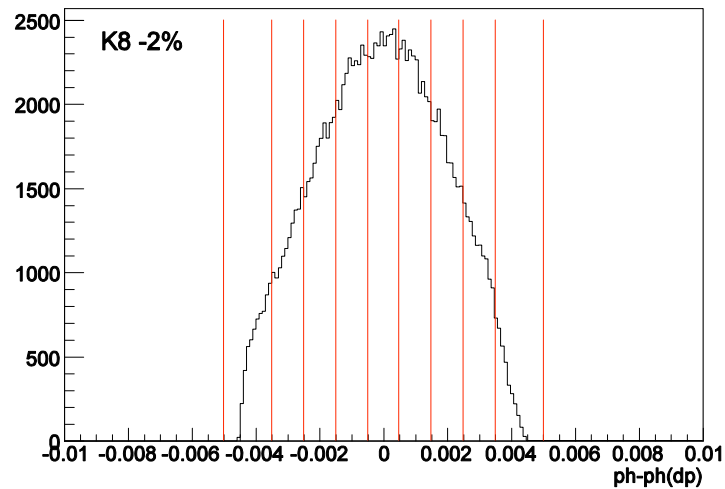
(K: kinematic factor)



# Spin Transport in HRS (COSY)

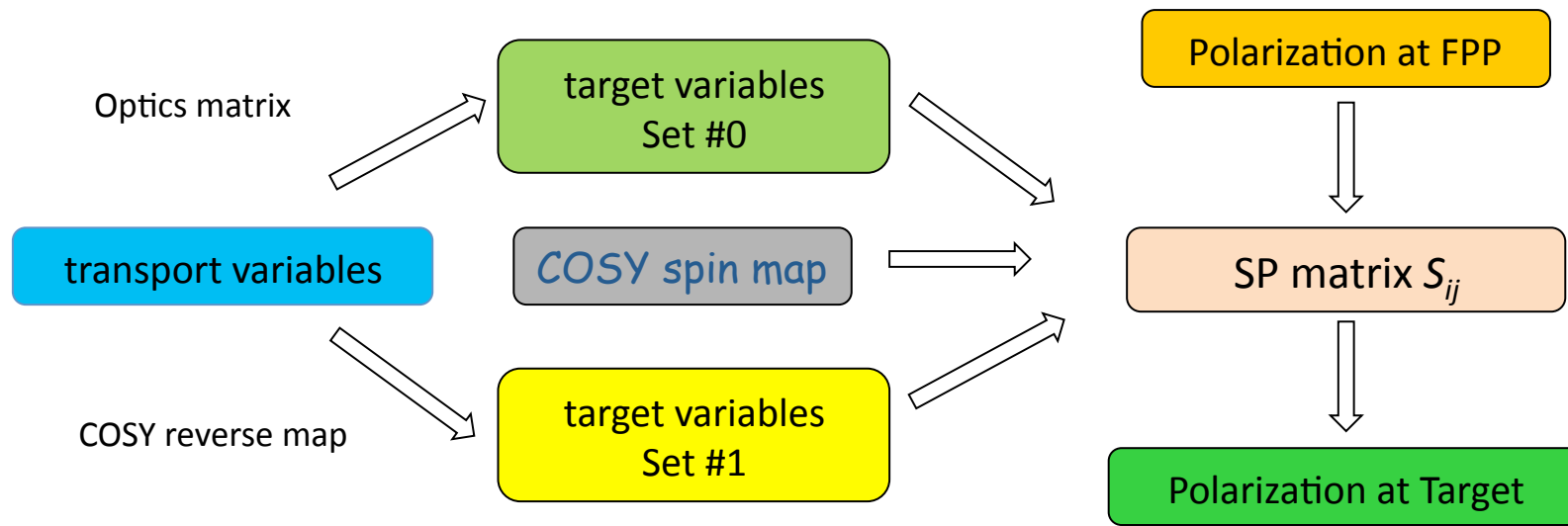


## Spin Transport in HRS



- Binning test for graphical cut.
- A rough check for existence of any possible background under elastic peak.
- No obvious indication of dependence on such variable.

## COSY Spin Precession Matrix



- Different SP matrix were generated by changing the default settings in COSY:

- dipole radius, drift distances, quadrupoles alignment
- central bending angle: **5.5 mrad**
- use COSY transport map to reconstruct target variables

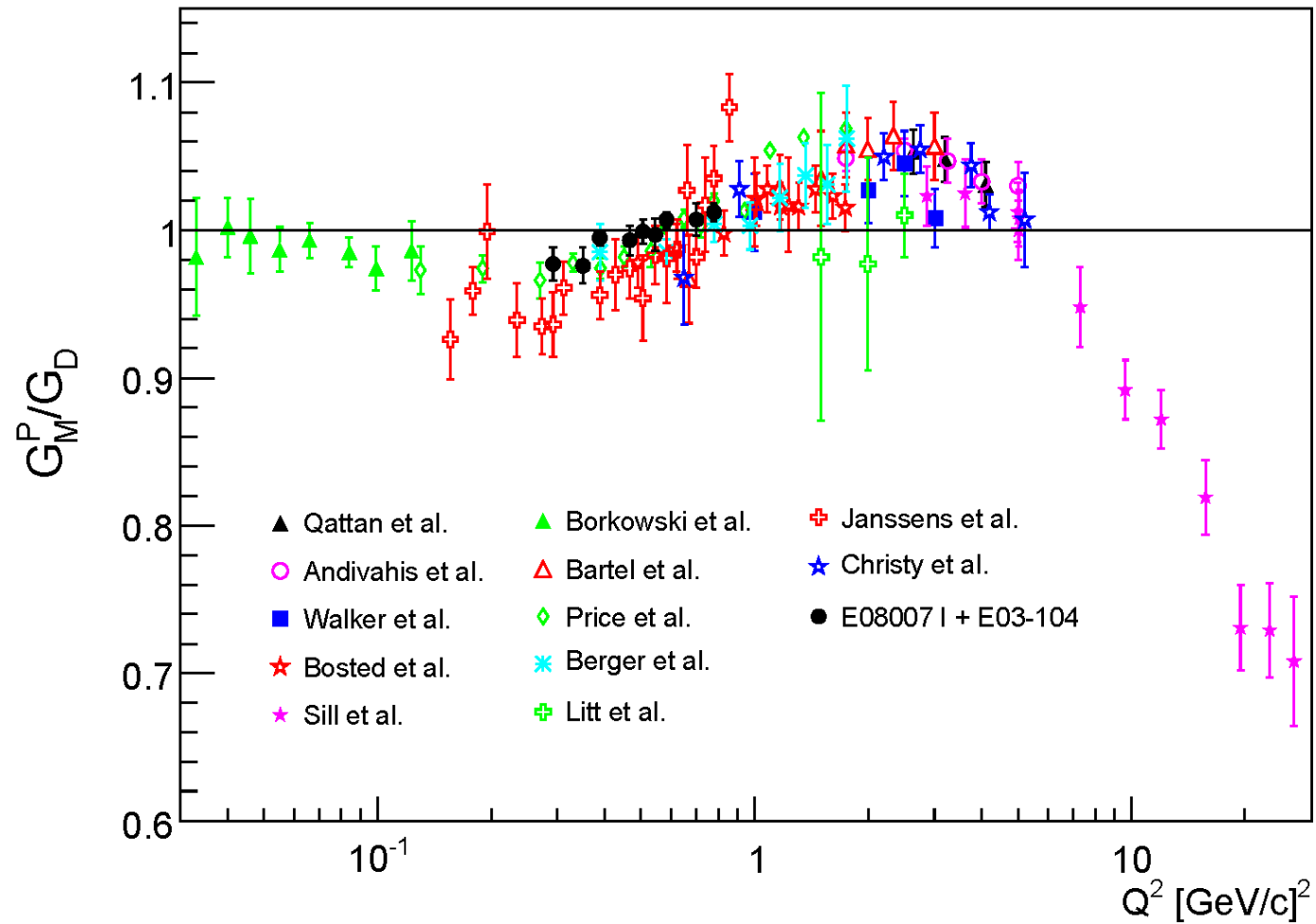
- Uncertainties on target variables (OPTICS):

- dp: **0.001**
- y\_tg: **0.001 m**
- ph\_tg: **0.7~1.2 mrad**
- th\_tg: **1 mrad**



## Individual Form Factors

- With the extract ratio constraint, refit the world reduced cross section data.



## Extraction of Polarization

- Full spin precession by COSY:
  - differential algebra-based.
  - defines the geometry and related setup of magnets.

$$\begin{pmatrix} P_x^{fpp} \\ P_y^{fpp} \end{pmatrix} = \begin{pmatrix} S_{xx} & S_{xy} & S_{xz} \\ S_{yx} & S_{yy} & S_{yz} \end{pmatrix} \begin{pmatrix} P_x^{tg} \\ \eta h P_y^{tg} \\ \eta h P_z^{tg} \end{pmatrix}$$

$$S_{ij} = \sum_{k,l,m,n,p} C_{ij}^{klmnp} x^k \theta^l y^m \phi^n \delta^p$$

focal plane

target frame

- Weighted-sum:

$$f(\phi) = \frac{1}{2\pi} \epsilon (1 + \lambda_x P_x^{tg} + \lambda_y h P_y^{tg} + \lambda_z h P_z^{tg}),$$

$$\lambda_x = A_y (S_{yx} \sin \phi - S_{xx} \cos \phi)$$

$$\lambda_y = \eta A_y (S_{yy} \sin \phi - S_{xy} \cos \phi)$$

$$\lambda_z = \eta A_y (S_{yz} \sin \phi - S_{xz} \cos \phi).$$

- efficiency cancels with different beam helicity

$$\begin{aligned} \int_0^{2\pi} f(\phi) \lambda_y d\phi &= h P_y^{tg} \int_0^{2\pi} f(\phi) \lambda_y^2 d\phi + \\ &\quad h P_z^{tg} \int_0^{2\pi} f(\phi) \lambda_y \lambda_z d\phi + \\ \int_0^{2\pi} f(\phi) \lambda_z d\phi &= h P_y^{tg} \int_0^{2\pi} f(\phi) \lambda_y \lambda_z d\phi + \\ &\quad h P_z^{tg} \int_0^{2\pi} f(\phi) \lambda_z^2 d\phi. \end{aligned}$$



$$\begin{pmatrix} \sum_i \lambda_{y,i} \\ \sum_i \lambda_{z,i} \end{pmatrix} = \begin{pmatrix} \sum_i \lambda_{y,i} \lambda_{y,i} & \sum_i \lambda_{z,i} \lambda_{y,i} \\ \sum_i \lambda_{y,i} \lambda_{z,i} & \sum_i \lambda_{z,i} \lambda_{z,i} \end{pmatrix} \begin{pmatrix} h P_y^{tg} \\ h P_z^{tg} \end{pmatrix}$$

## Impacts III

- Isoscalar & Isovector FFs (important for Lattice QCD):

$$F_i^s = \frac{1}{2}(F_i^p + F_i^n), F_i^v = \frac{1}{2}(F_i^p - F_i^n)$$

- Plots show fractional change in IS and IV FFs by using the new parameterization vs. the old parameterization.

