

# Quasifree ( $e, e'p$ ) Reactions on Nuclei with Neutron Excess

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## MOTIVATION

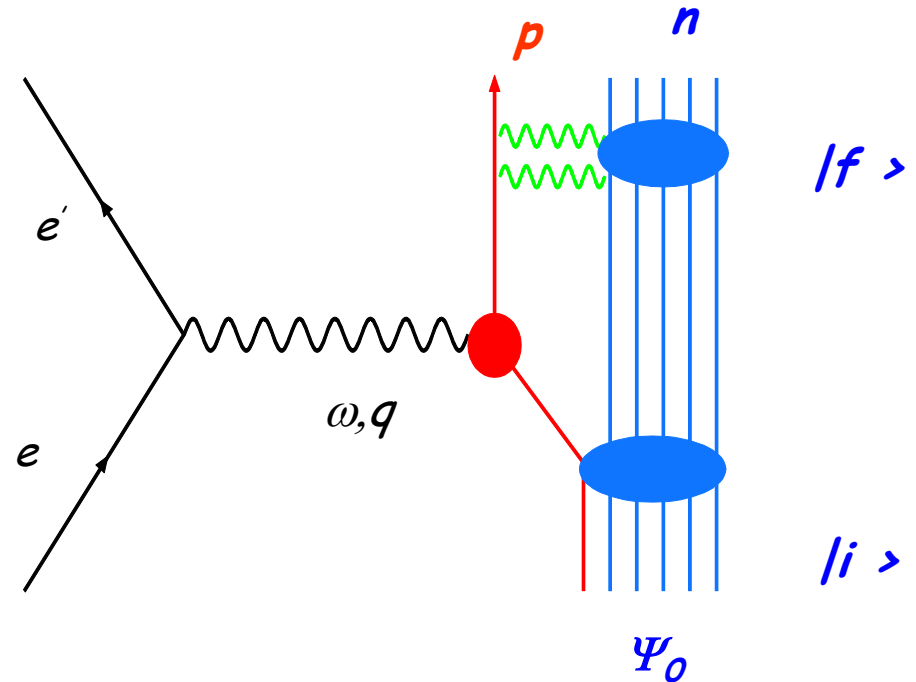
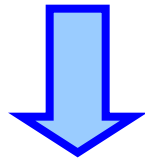
- understanding the evolution of nuclear properties as a function of  $N/Z$
- nuclear reactions main source of information on nuclear properties
- direct reactions give insight into the s.p. properties
- advantages of the elm probe:  $(e,e'p)$  preferential tool to study proton-hole states, bound protons, validity and limits of IPSM
- large amount of  $(e,e'p)$  data, accurate information on s.p. properties of stable nuclei
- advent of RIB facilities will provide data on unstable nuclei
- electron RIB colliders that use storage rings under construction (GSI, RIKEN) will offer unprecedented opportunities to study exotic nuclei with electron scattering (ELISE at FAIR, SCRIT at RIKEN)
- exclusive  $(e,e'p)$  knockout experiments (ELISE at FAIR, SCRIT at RIKEN)

# OUTLINE

- DWIA model for (e,e'p)
- NIKHEF data  $^{40}\text{Ca}$   $^{48}\text{Ca}$
- original analysis DWIA
- comparison of different models DWIA, RDWIA, different s.p. wave functions
- calculations performed for Ca isotopes: 40, 48, 52, 60
- evolution of nuclear properties with models of proven reliability in stable isotopes will test the ability of the established nuclear theory in the domain of exotic nuclei

# Direct knockout DWIA (e,e'p)

- ☼ exclusive reaction:  $n$
- ☼ DKO mechanism: the probe interacts through a one-body current with one nucleon that is then emitted the remaining nucleons are spectators



$$\langle f | J^\mu(\mathbf{q}) | i \rangle \longrightarrow \lambda_n^{1/2} \langle \chi_{\mathbf{p}}^{(-)} | j^\mu(\mathbf{q}) | \phi_n \rangle$$

## Direct knockout DWIA (e,e'p)

$$\lambda_n^{1/2} \langle \chi^{(-)} | j^\mu | \phi_n \rangle$$

- $j^\mu$  one-body nuclear current
- $\chi^{(-)}$  s.p. scattering w.f.  $H^+(\omega+E_m)$
- $\phi_n$  s.p. bound state overlap function  $H(-E_m)$
- $\lambda_n$  spectroscopic factor
- $\chi^{(-)}$  and  $\phi$  consistently derived as eigenfunctions of a Feshbach optical model Hamiltonian

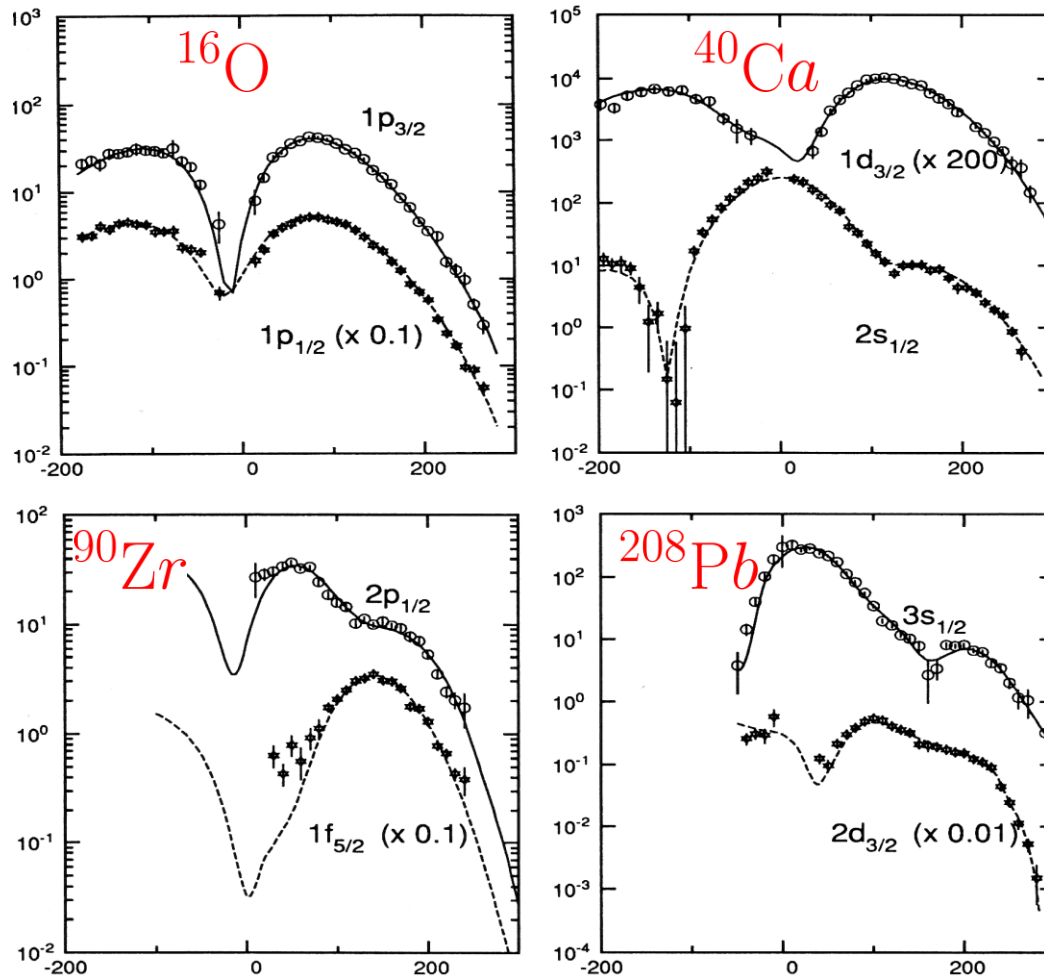
$$\mathcal{H}(E) = P H P + P H Q \frac{1}{E - Q H Q + i\eta} Q H P$$

# DWIA calculations

- ☀ phenomenological ingredients usually adopted
- ☀  $\chi^{(-)}$  phenomenological optical potential
- ☀  $\phi_n$  phenomenological s.p. wave functions WS, HF (some calculations including correlations are available)
- ☀  $\lambda_n$  extracted in comparison with data: reduction factor applied to the calculated c.s. to reproduce the magnitude of the experimental c.s.
- ☀ DWIA RDWIA calculations with Coulomb distortion  
excellent description of (e,e'p) data

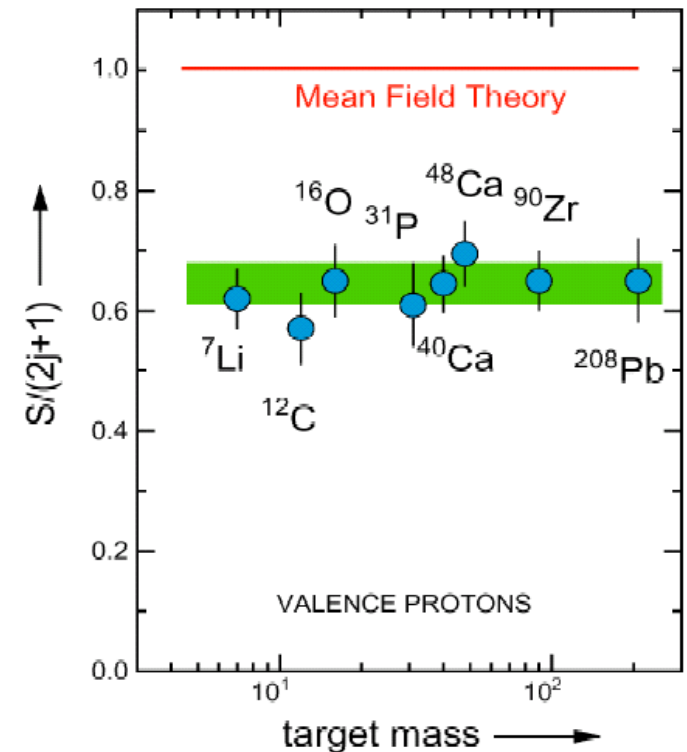
Experimental data:  $E_m$  and  $p_m$  distributions

# Experimental data: $p_m$ distributions



reduction factors applied:  
spectroscopic factors

0.6 - 0.7



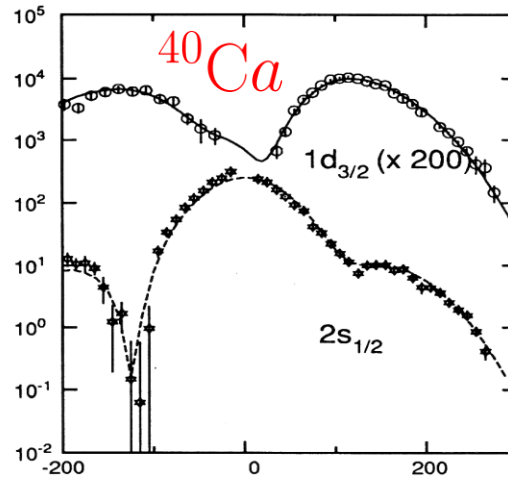
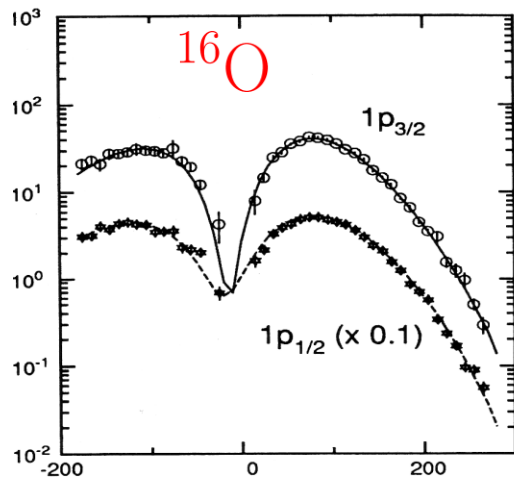
NIKHEF data & CDWIA calculations

## NIKHEF data: $^{40}\text{Ca}(e,e'p)$ , $^{48}\text{Ca}(e,e'p)$

- NIKHEF data  $^{40}\text{Ca}$   $^{48}\text{Ca}$
- original analysis: DWIA with phenomenological WS bound state w.f., depth of the WS well adjusted to give the experimentally observed separation energy, rms radius determined to fit the experimental momentum distribution

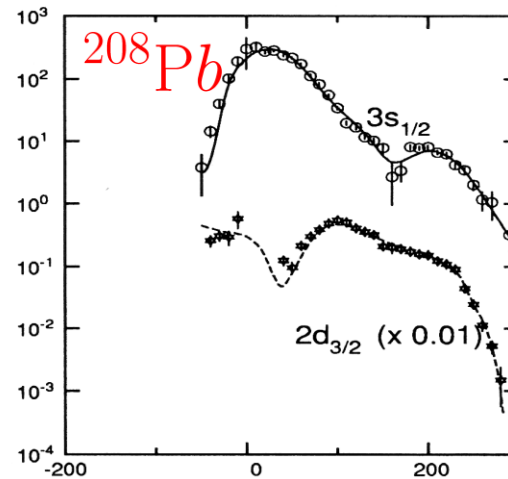
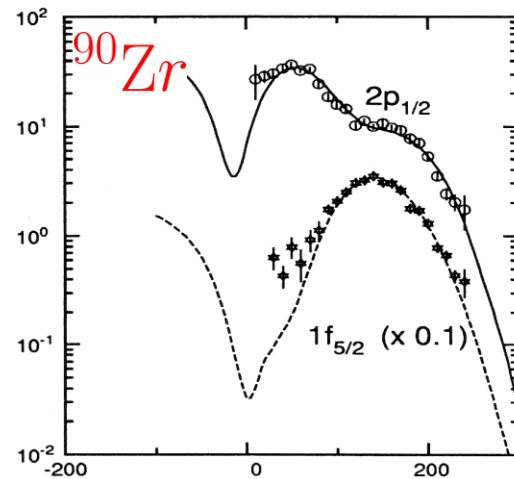


# Experimental data: $p_m$ distributions



$1d_{3/2}$

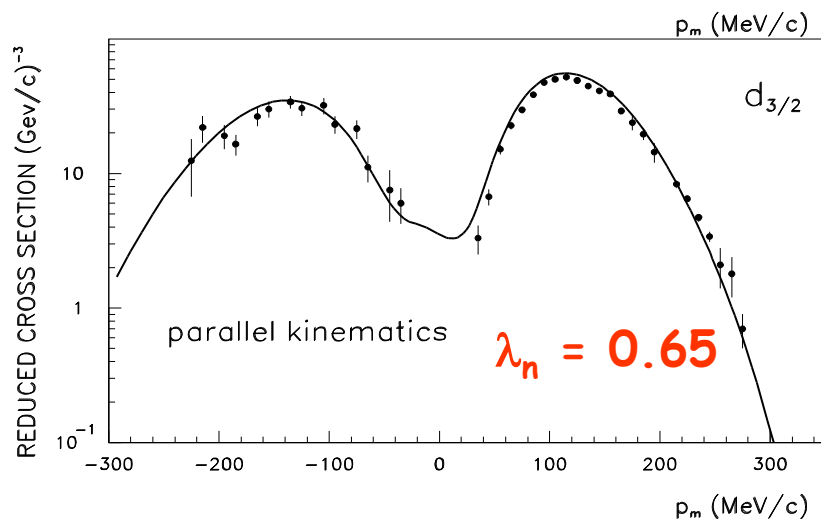
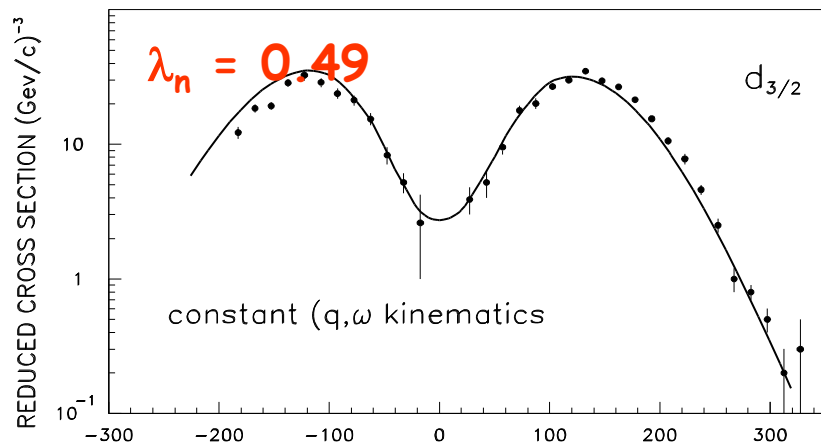
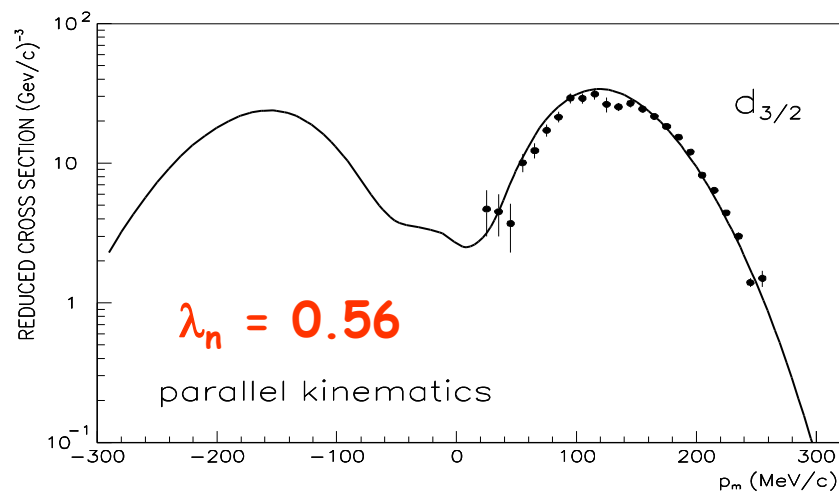
$2s_{1/2}$



NIKHEF data + CDWIA calculations

$^{40}\text{Ca}(e, e'p)$  $^{48}\text{Ca}(e, e'p)$ 

DWIA WS wave function

 $1d_{3/2}$ 

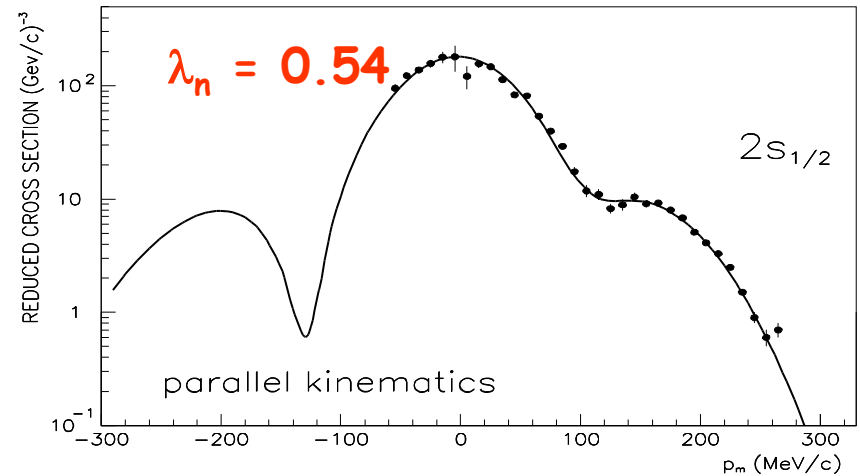
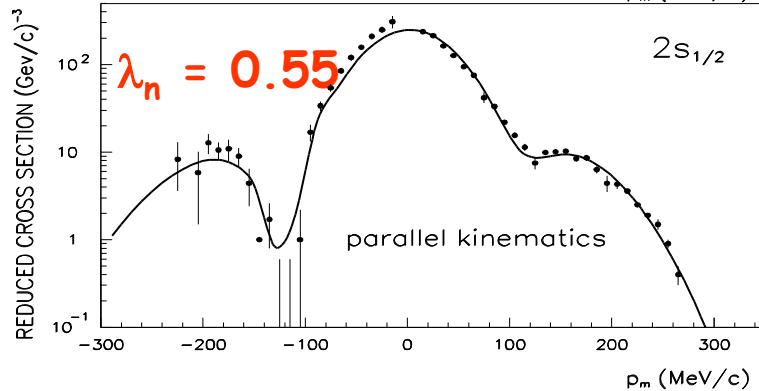
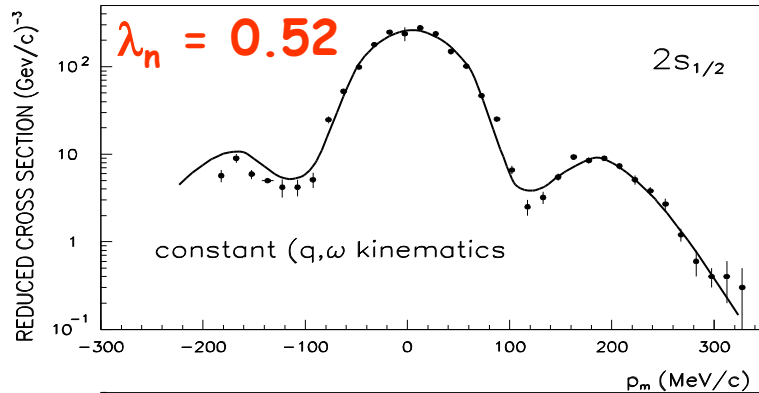
$(\omega, q)$  const:  $E_0 = 483.2$  MeV  $\theta = 61.52$  deg.  $q = 450$  MeV/c  $T_p = 100$  MeV

parallel kin:  $E_0 = 483.2$  MeV  $T_p = 100$  MeV

NIKHEF data G.J. Kramer Ph. D. Thesis (1990)

$^{40}\text{Ca}(e, e'p)$  $^{48}\text{Ca}(e, e'p)$ 

## DWIA WS wave function

 $2s_{1/2}$ 

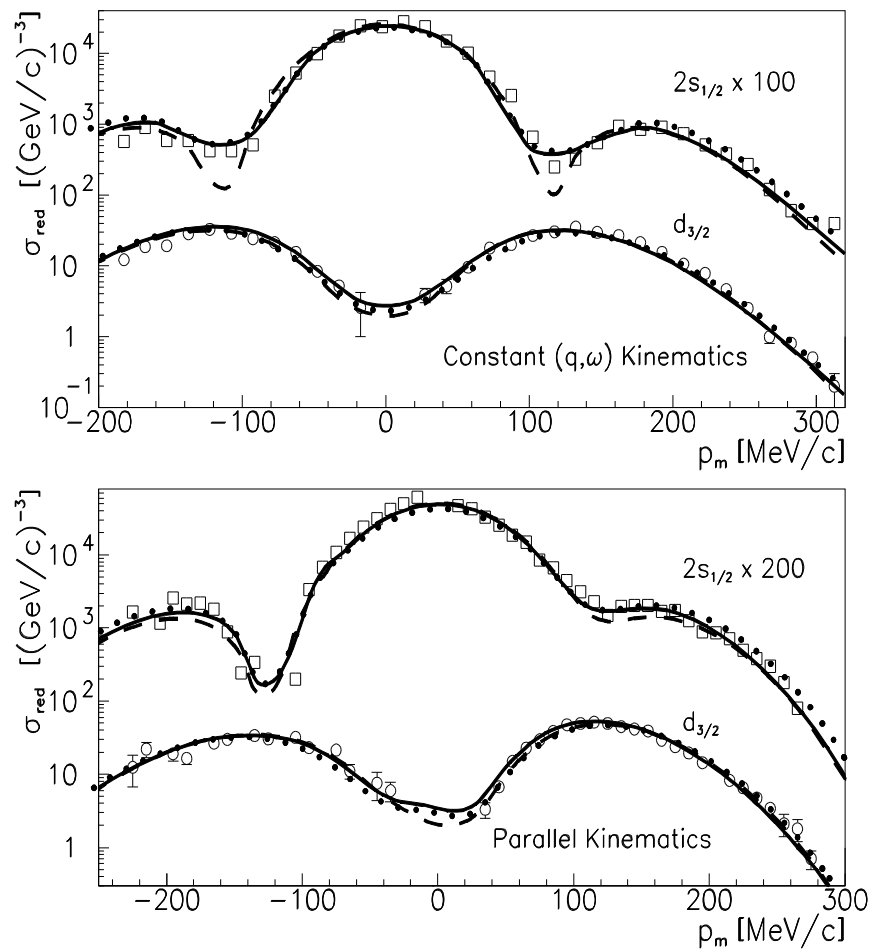
$(\omega, q)$  const:  $E_0 = 483.2$  MeV  $\theta = 61.52$  deg.  $q = 450$  MeV/c  $T_p = 100$  MeV

parallel kin:  $E_0 = 483.2$  MeV  $T_p = 100$  MeV

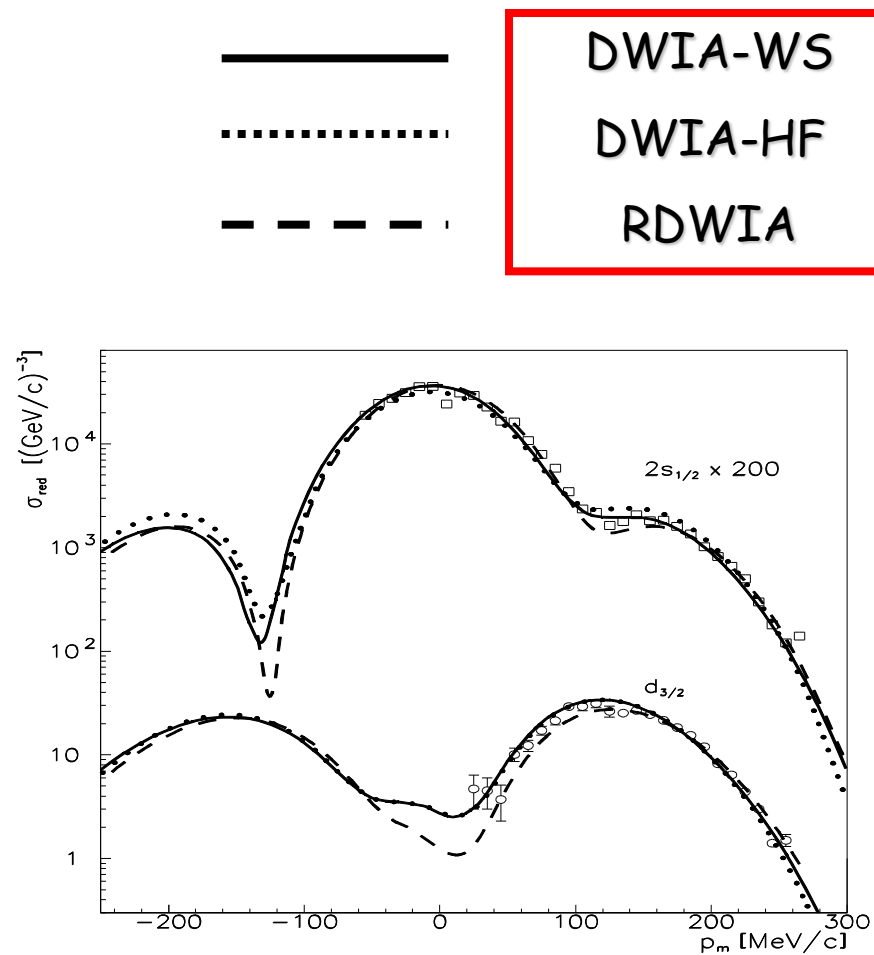
## Comparison of different models

- ☀ DWIA with phenomenological WS wave functions (DWIA-WS)
- ☀ DWIA with HF wave functions from two different parametrizations of the finite-range Gogny interactions D1S and D1M. Results presented for the new D1M force (DWIA-HF)
- ☀ RDWIA relativistic model, ROP for the scattering state, the bound states are obtained in the context of the RMF approach solving the Dirac-Hartree equations. The nucleon interaction is derived from a relativistic Lagrangian containing  $\sigma$ ,  $\omega$ ,  $\rho$  meson fields and also the photon field
- ☀ E- and A-dependent optical potentials contain central, spin-orbit, Coulomb terms and a term dependent on the  $(N-Z)/A$  asymmetry
- ☀ comparison with NIKHEF data on  $^{40}\text{Ca}$   $^{48}\text{Ca}$

# $^{40}\text{Ca}(e, e'p)$



# $^{48}\text{Ca}(e, e'p)$



$^{40}\text{Ca}(e,e'p)$

$1d_{3/2}$

$^{48}\text{Ca}(e,e'p)$

$\lambda_n = 0.49$  DWIA-WS

0.51 DWIA-HF

0.49 RDWIA

$(\omega, q)$  const kin

$\lambda_n = 0.65$  DWIA-WS

0.64 DWIA-HF

0.69 RDWIA

parallel kin

$\lambda_n = 0.56$  DWIA-WS

0.55 DWIA-HF

0.52 RDWIA

$^{40}\text{Ca}(e, e'p)$

$2s_{1/2}$

$^{48}\text{Ca}(e, e'p)$

$\lambda_n = 0.55$  DWIA-WS  
0.62 DWIA-HF  
0.51 RDWIA

$(\omega, q)$  const kin

$\lambda_n = 0.56$  DWIA-WS  
0.55 DWIA-HF  
0.52 RDWIA

parallel kin

$\lambda_n = 0.54$  DWIA-WS  
0.58 DWIA-HF  
0.55 RDWIA

$$^{40,48,52,60}\text{Ca}(e,e'p)$$

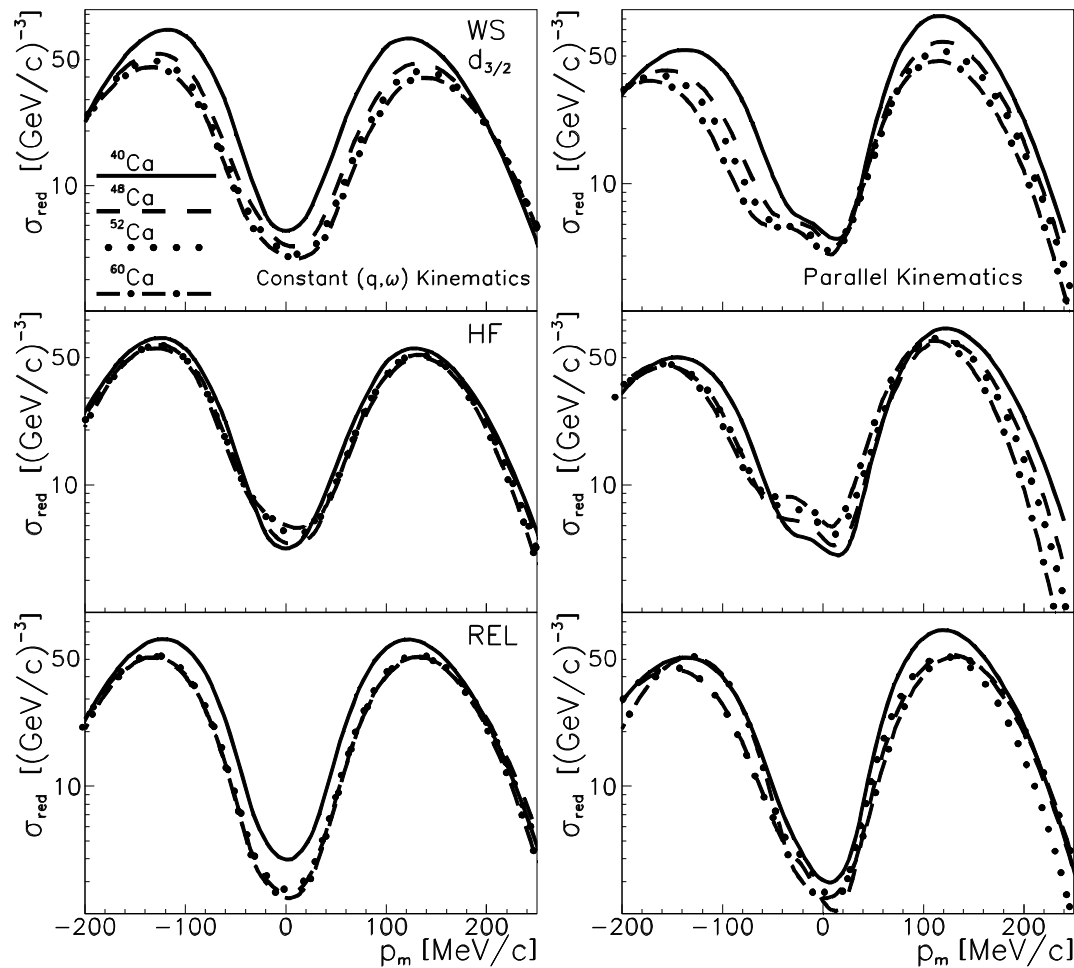
- ☀ DWIA-WS DWIA-HF and RDWIA for Ca isotopes
- ☀ even-even isotopes, spherical nuclei where the s.p. level are fully occupied and pairing effects should be minimized



40,48,52,60Ca(e,e'p)

$1d_{3/2}$

— 40  
 - - - 48  
 ..... 52  
 - . - . 60



DWIA-WS

DWIA-HF

RDWIA

constant ( $q, \omega$ )

parallel

40,48,52,60Ca(e,e'p)

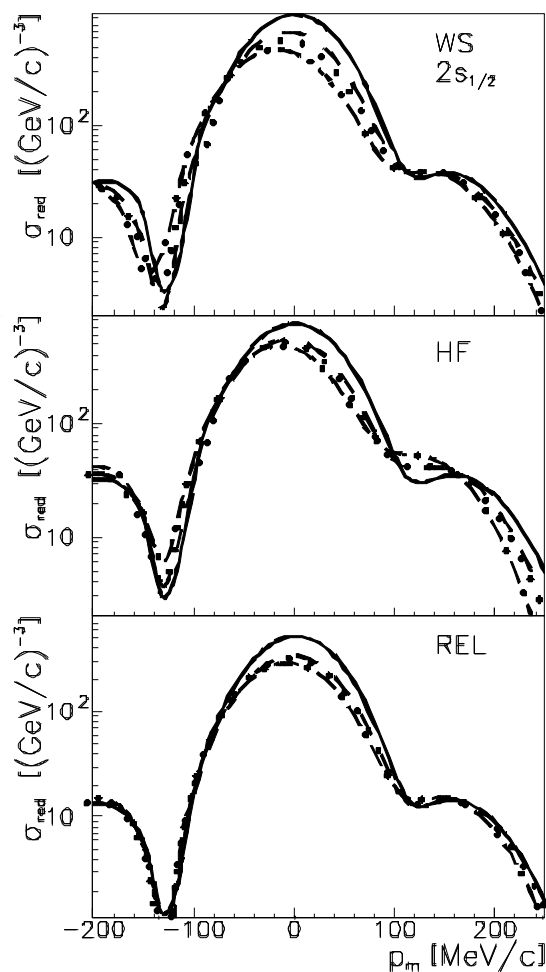
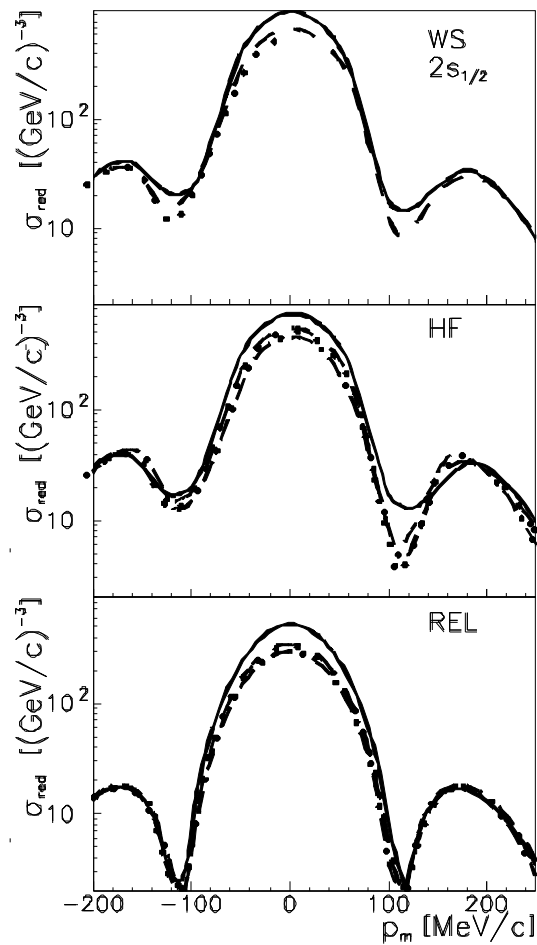
$2s_{1/2}$

— 40  
 - - - 48  
 ..... 52  
 - . - . 60

DWIA-WS

DWIA-HF

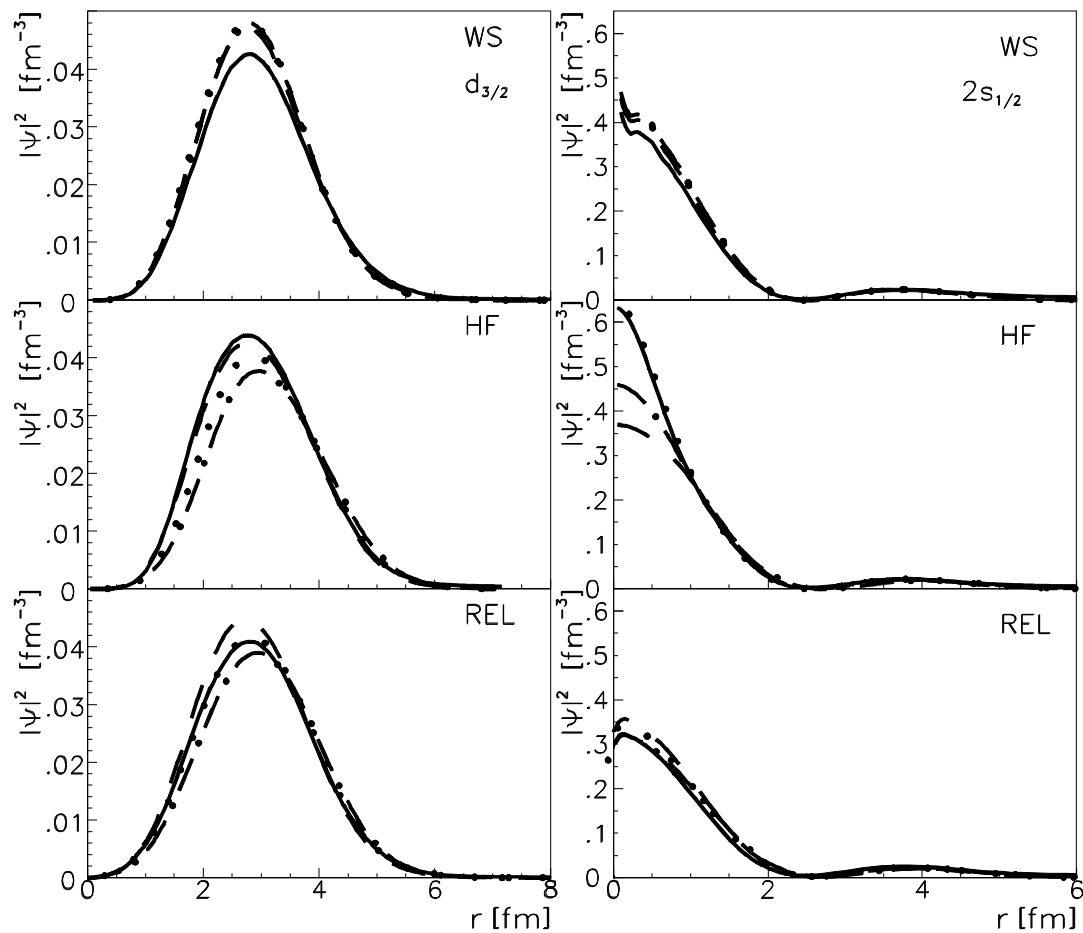
RDWIA



constant  $(q, \omega)$

parallel

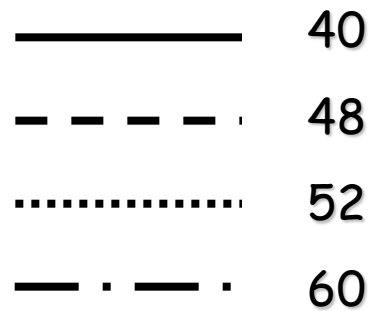
40,48,52,60Ca  $|\phi|^2$



WS

HF

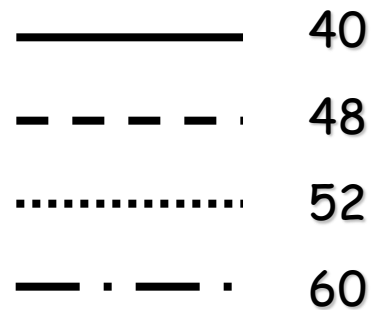
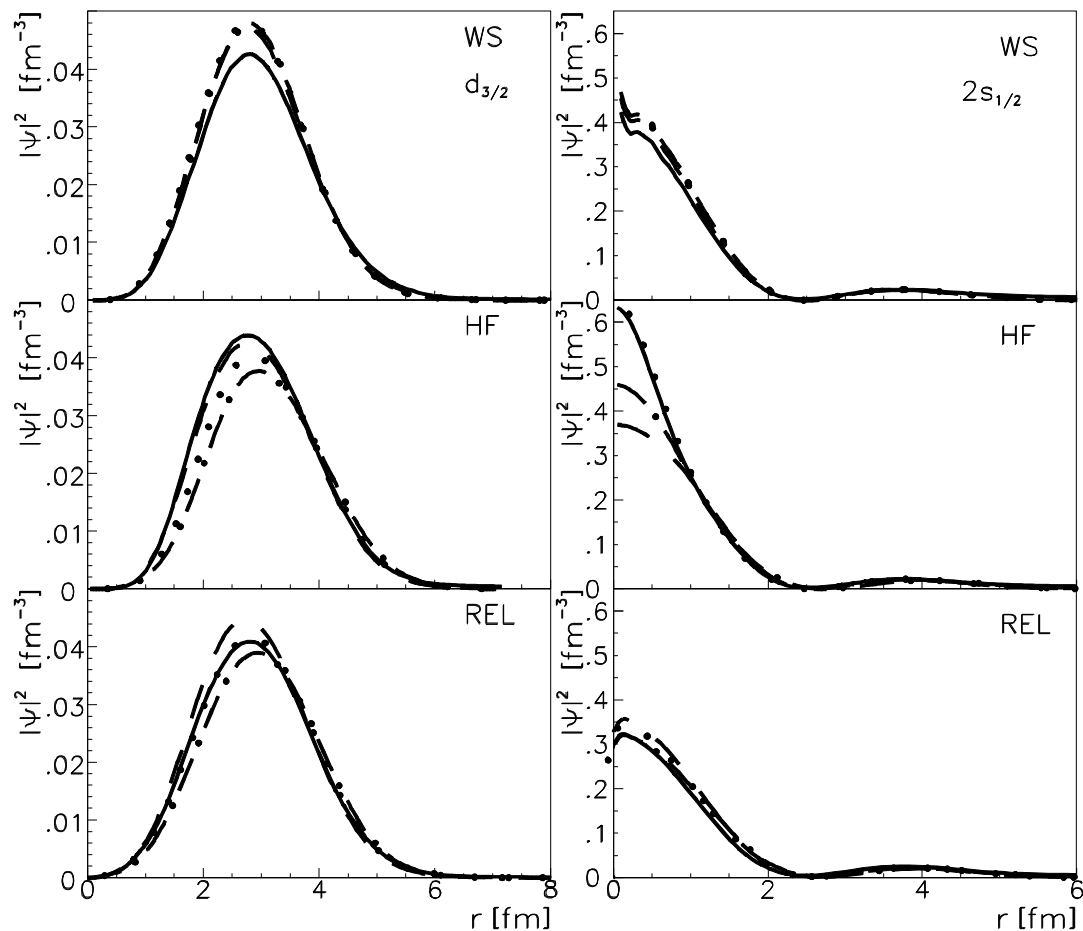
REL



$1d_{3/2}$

$2s_{1/2}$

40,48,52,60Ca  $|\phi|^2$



WS

HF

increasing N/Z different  
behavior for the wave  
functions and the cross  
sections: **FSI**

REL

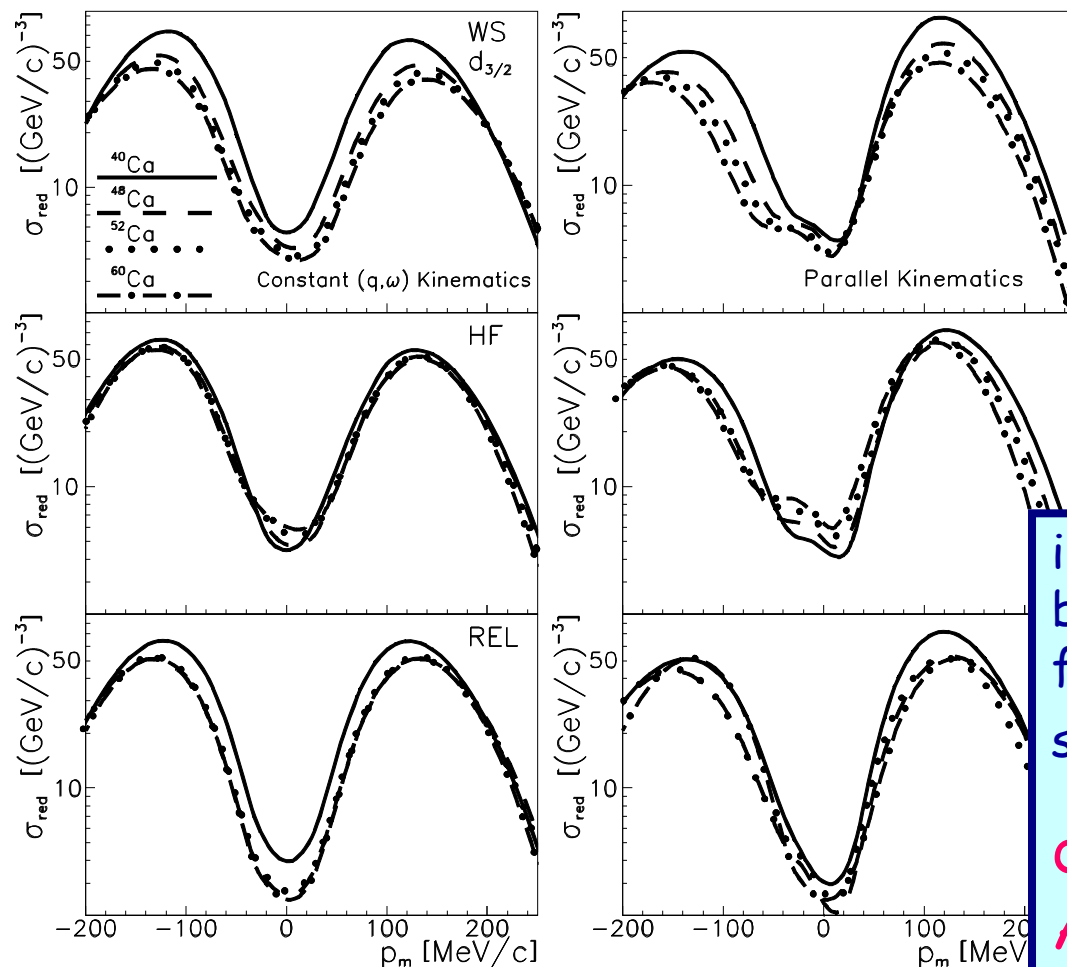
$d_{3/2}$

$2s_{1/2}$

40,48,52,60Ca(e,e'p)

$1d_{3/2}$

— 40  
- - - 48  
... 52  
- . - . 60



DWIA-WS

DWIA-HF

increasing N/Z different  
behavior for the wave  
functions and the cross  
sections: FSI

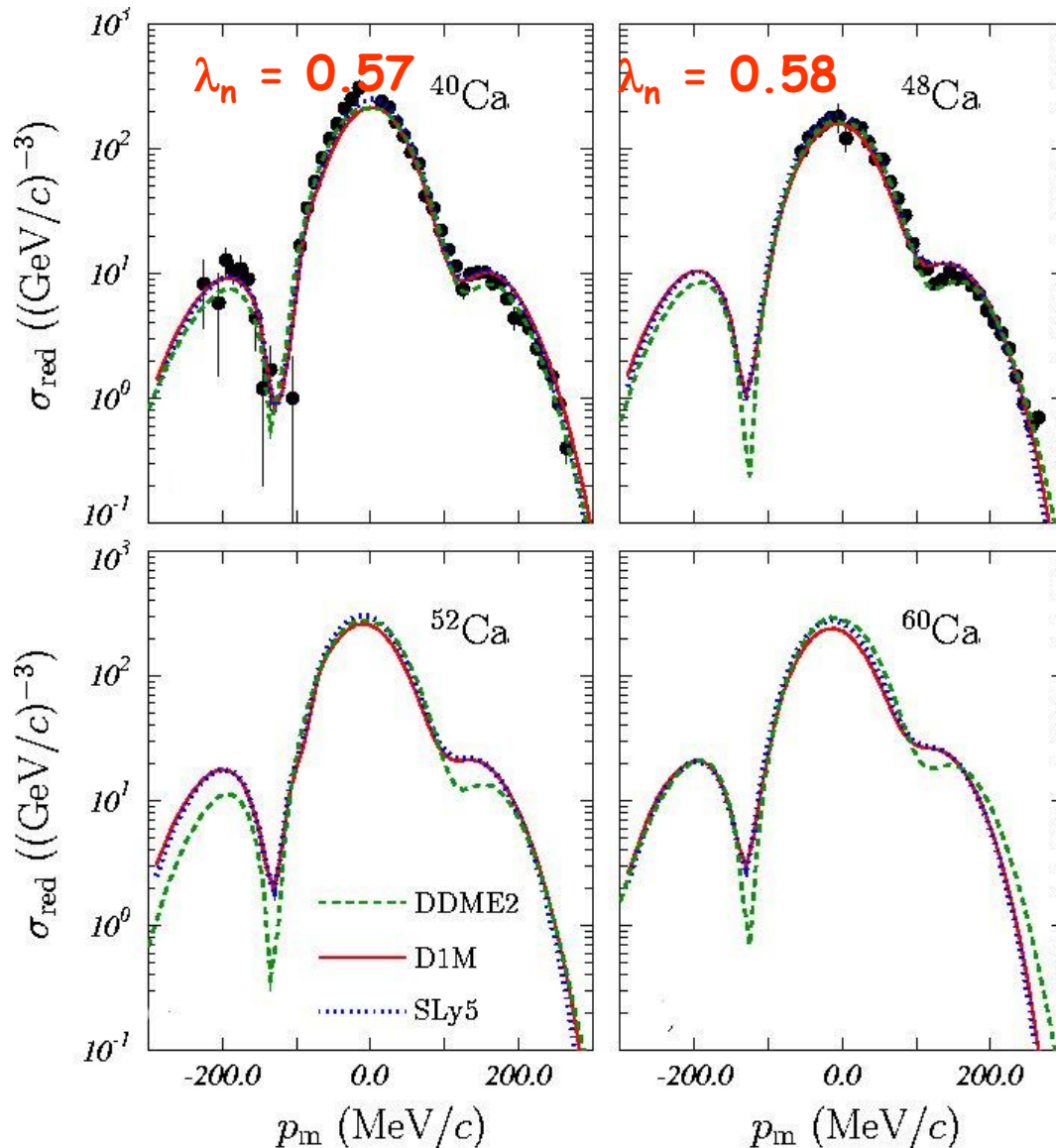
difference due to the  
A- dependence of the  
optical potential

constant (q,  $\omega$ )

parallel

# 40,48,52,60Ca(e,e'p)

$2s_{1/2}$



D1M HF finite-range DWIA

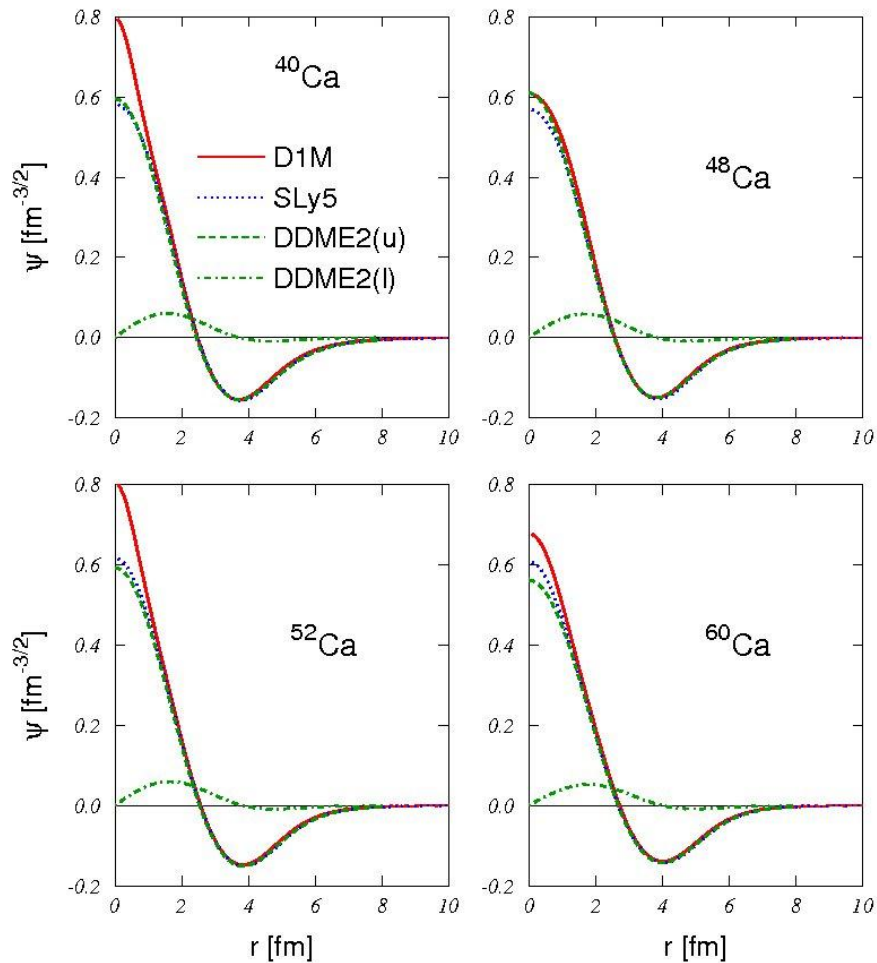
SLy5 HF zero range DWIA

DDME2 relativistic density  
dependent meson-nucleon  
couplings

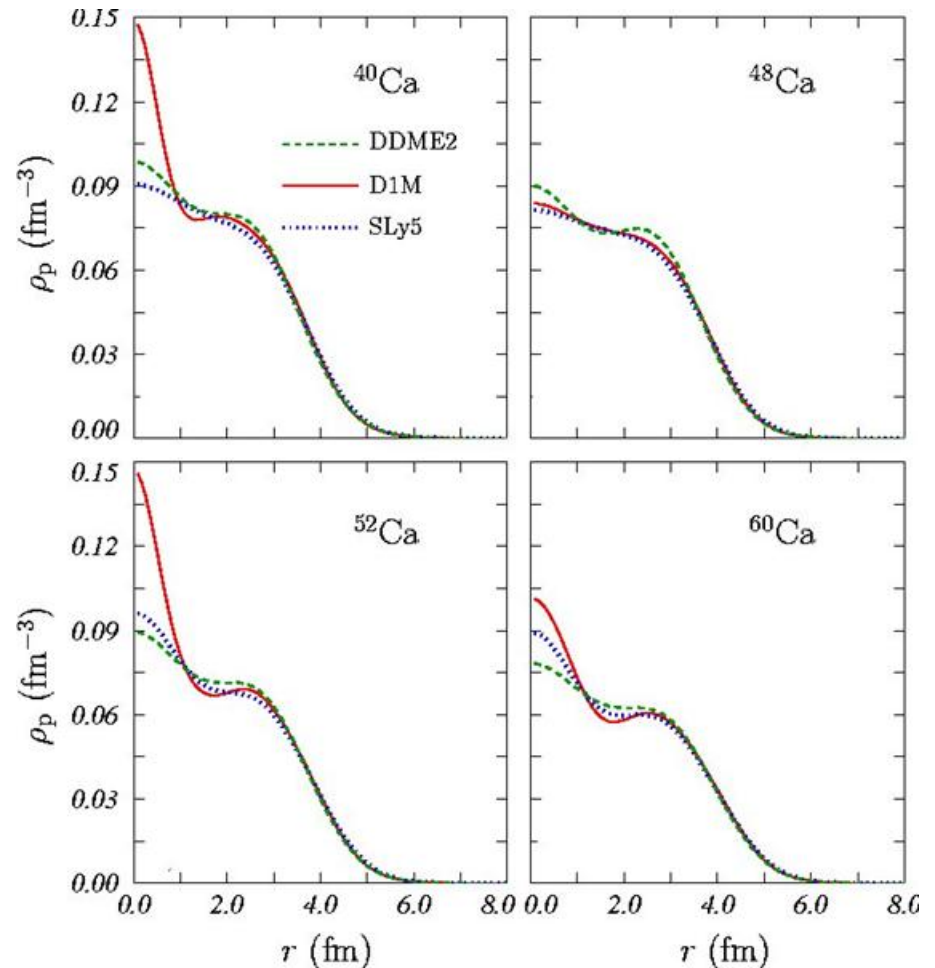
parallel kin.

G.Co', V. De Donno, P. Finelli, M. Grasso, M. Anguiano, A.M. Lallena, C. Giusti, A. Meucci, F.D. Pacati

## wave functions $2s_{1/2}$

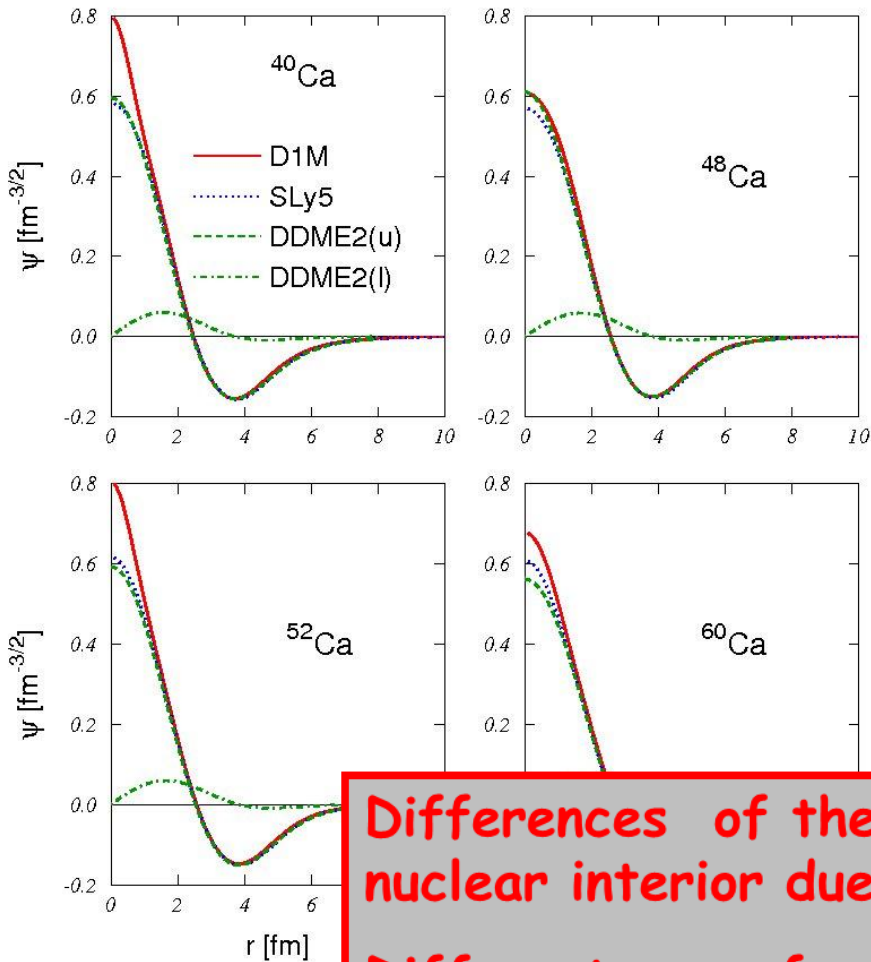


## proton distributions

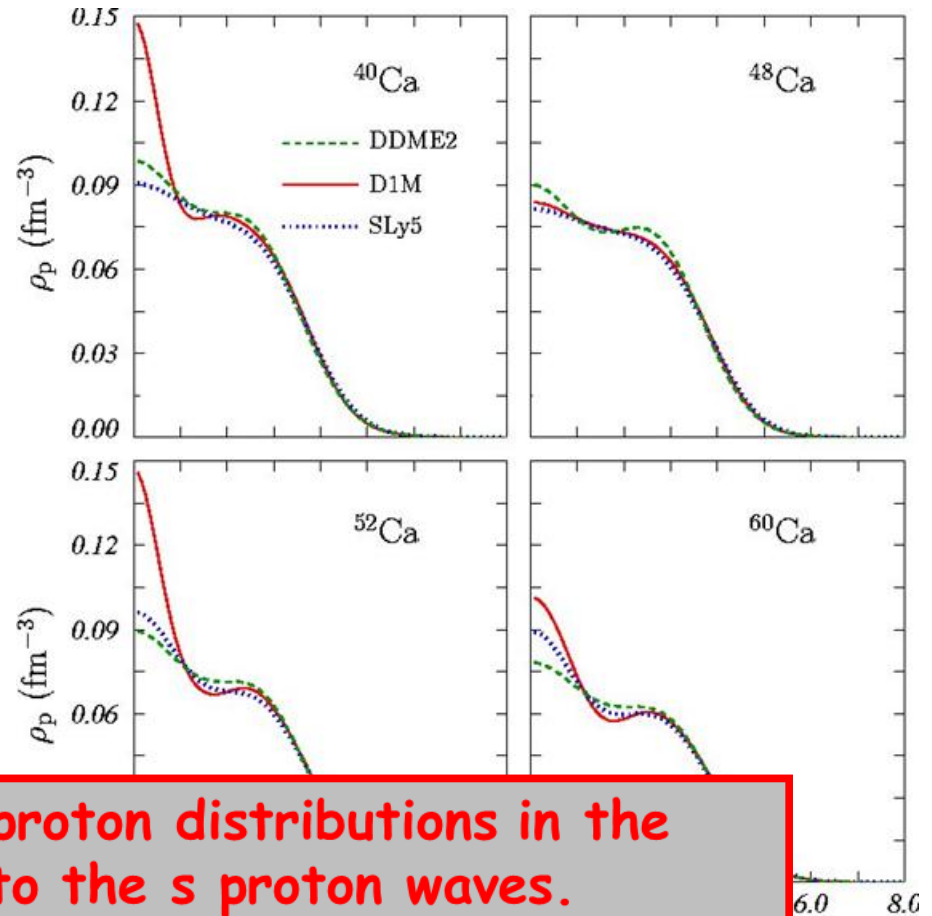


G.Co', V. De Donno, P. Finelli, M. Grasso, M. Anguiano, A.M. Lallena, C. Giusti, A. Meucci, F.D. Pacati

## wave functions $2s_{1/2}$



## proton distributions



Differences of the proton distributions in the nuclear interior due to the s proton waves.

Different wave functions produce similar (e,e'p) cross sections. The considered kinematics unable to emphasize the differences in the s.p wave functions.



## CONCLUSIONS (I)

- evolution of nuclear properties with models of proven reliability in stable isotopes (DWIA-WS DWIA-HF RDWIA)
- all the considered models give good and similar description of the available  $(e,e'p)$  data on  $^{40}\text{Ca}$  and  $^{48}\text{Ca}$
- general behavior of the cross sections with respect to the increasing N/Z asymmetry is analogous for all the three models: the reduced cross sections are larger and narrower for the lighter isotopes and evolve by lowering and widening increasing N
- the behavior of the s.p. bound-state wave functions shows different trends for the different models
- the dependence of the w.f. on N/Z is responsible for only a part of the differences in the calculated cross sections, an important and crucial contribution is given also by FSI which are described in the calculations by phenomenological optical potentials
- the optical potential is an important ingredient of the model, affects the size and the shape of the cross section in a way that strongly depends on kinematics
- the dependence of the optical potential on N/Z deserves careful investigation

## CONCLUSIONS (II)

- spectroscopic factors and correlations: recent exp. and theor. studies indicate that the s.f. depend on  $N/Z$ , in general the quenching of the quasi-hole states becomes stronger increasing the separation energy (increasing  $N$ )
- $(e,e'p)$  measurements on nuclei with neutron excess would offer a unique opportunity for studying the dependence of the properties of bound protons on  $N/Z$
- the present results can serve as a useful reference for future experiments
- comparison with data can confirm or invalidate the predictions of the models and will test the ability of the established nuclear theory in the domain of exotic nuclei