

High-dimensional surprises near
the glass and the jamming transitions:
connecting spin glasses with reality

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Dimensional collaborators

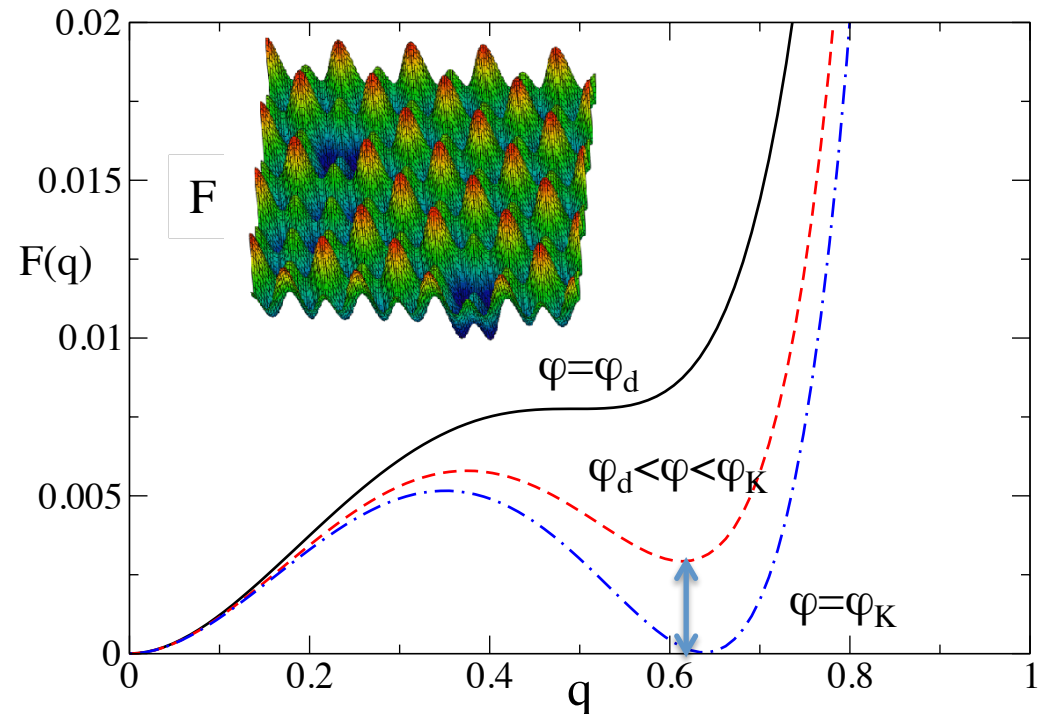
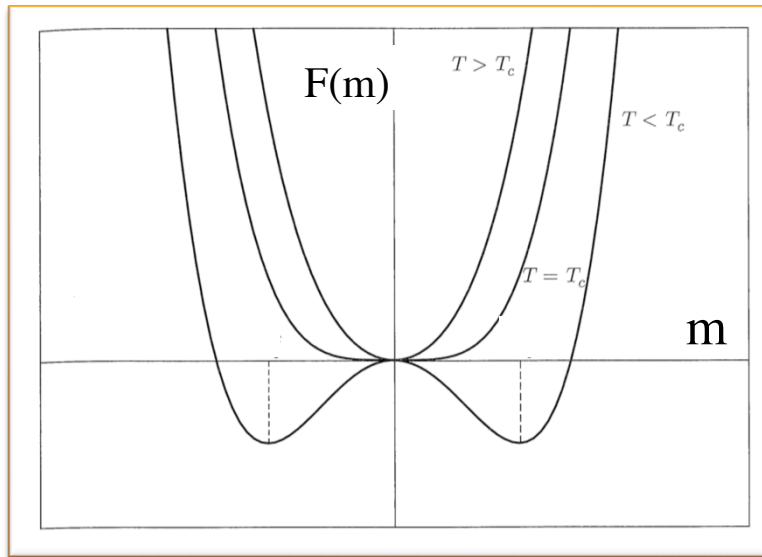
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Spin-based MF picture for glasses (RFOT)

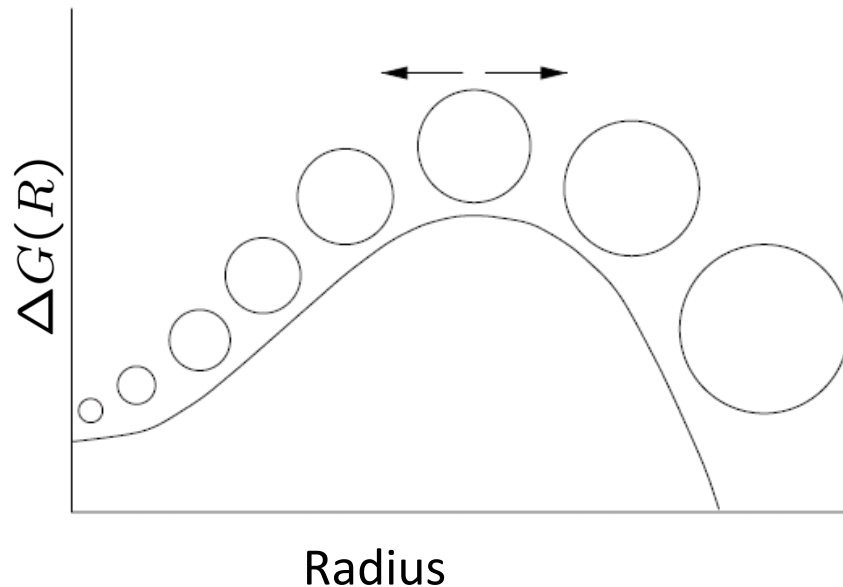


Exact in $d=\infty$.

The dynamics gets infinitely sluggish at the onset of the breakdown in states φ_d .

Kirkpatrick, Thirumalai, Wolynes (1987-1989); Parisi *et al.*
 Debenedetti and Stillinger, Nature (2001)
 Franz, Parisi (1997); KPZ (2012)

Problem 0: Crystal nucleation



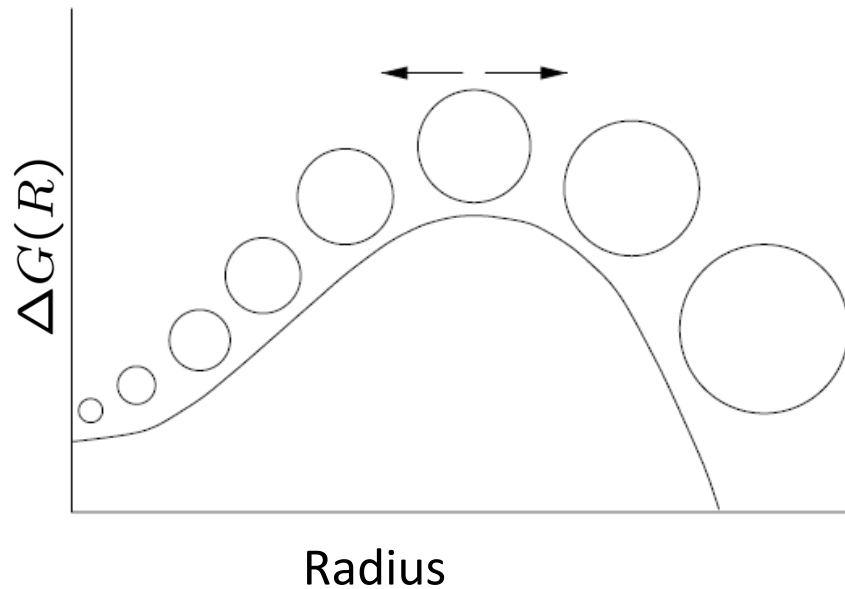
The crystal state is completely orthogonal to this description and it thought to interfere with its observation.

Typical solution: complex alloys
Other possible solution: increase spatial dimension

$$\Delta G(R) = \gamma S_{d-1} R^{d-1} - \Delta\mu \rho_s V_d R^d$$

$$\Delta G^\ddagger(R^*) \sim \frac{(2d\pi)^{d/2} \gamma^d}{\Delta\mu^{d-1}}$$

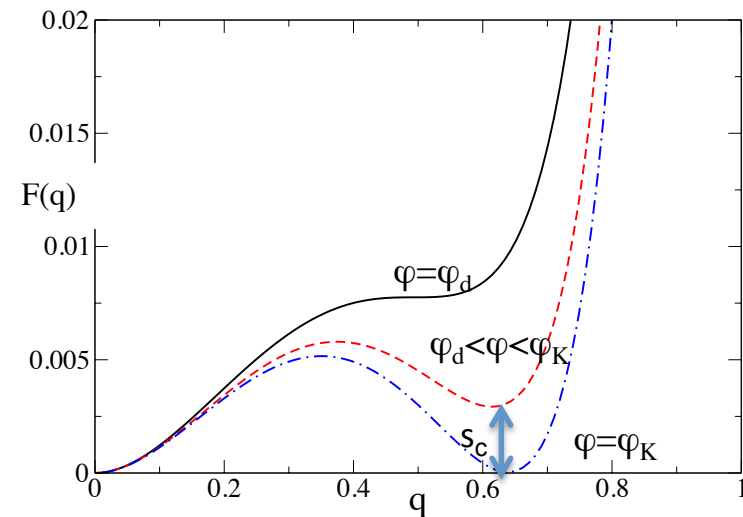
Problem 1: Glass-glass nucleation and (facilitated) hopping



$$\Delta G(R) = \gamma S_{d-1} R^{d-1} - s_c \rho V_d R^d$$

$$\Delta G^\ddagger(R^*) \sim \frac{(2d\pi)^{d/2} \gamma^d}{s_c^{d-1}}$$

Mean-field hopping: see Yuliang Jin's poster

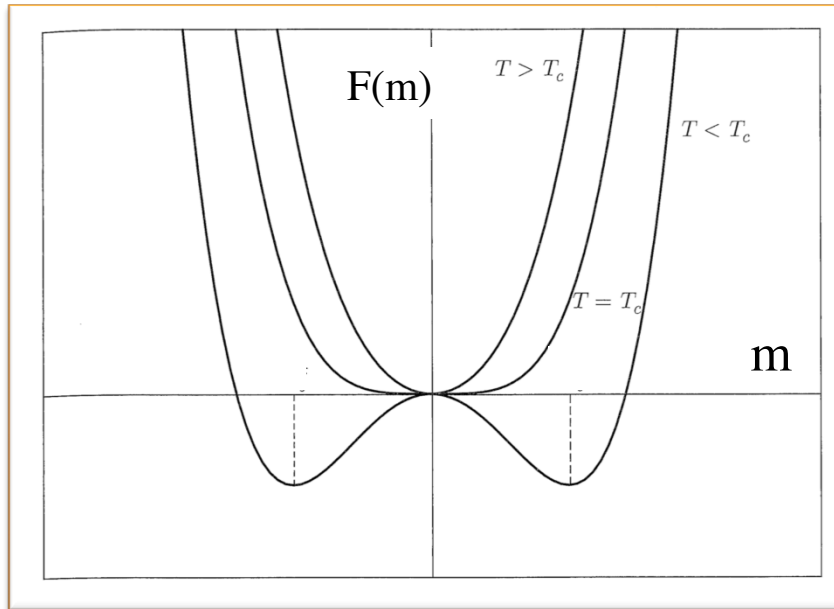


Once replica symmetry is broken, nucleation from one state to another is still possible.

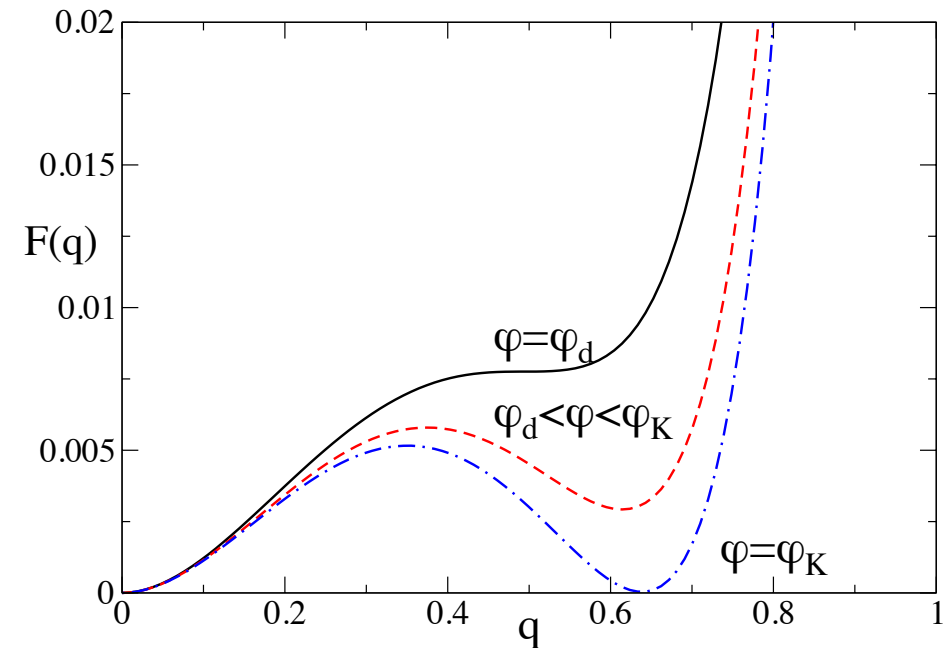
Typical solution: patch theory, look in crossover area

Other possible solution: increase spatial dimension

Problem 2: Below d_u fluctuations renormalize, or worse



For an Ising-type field theory $d_u=4$

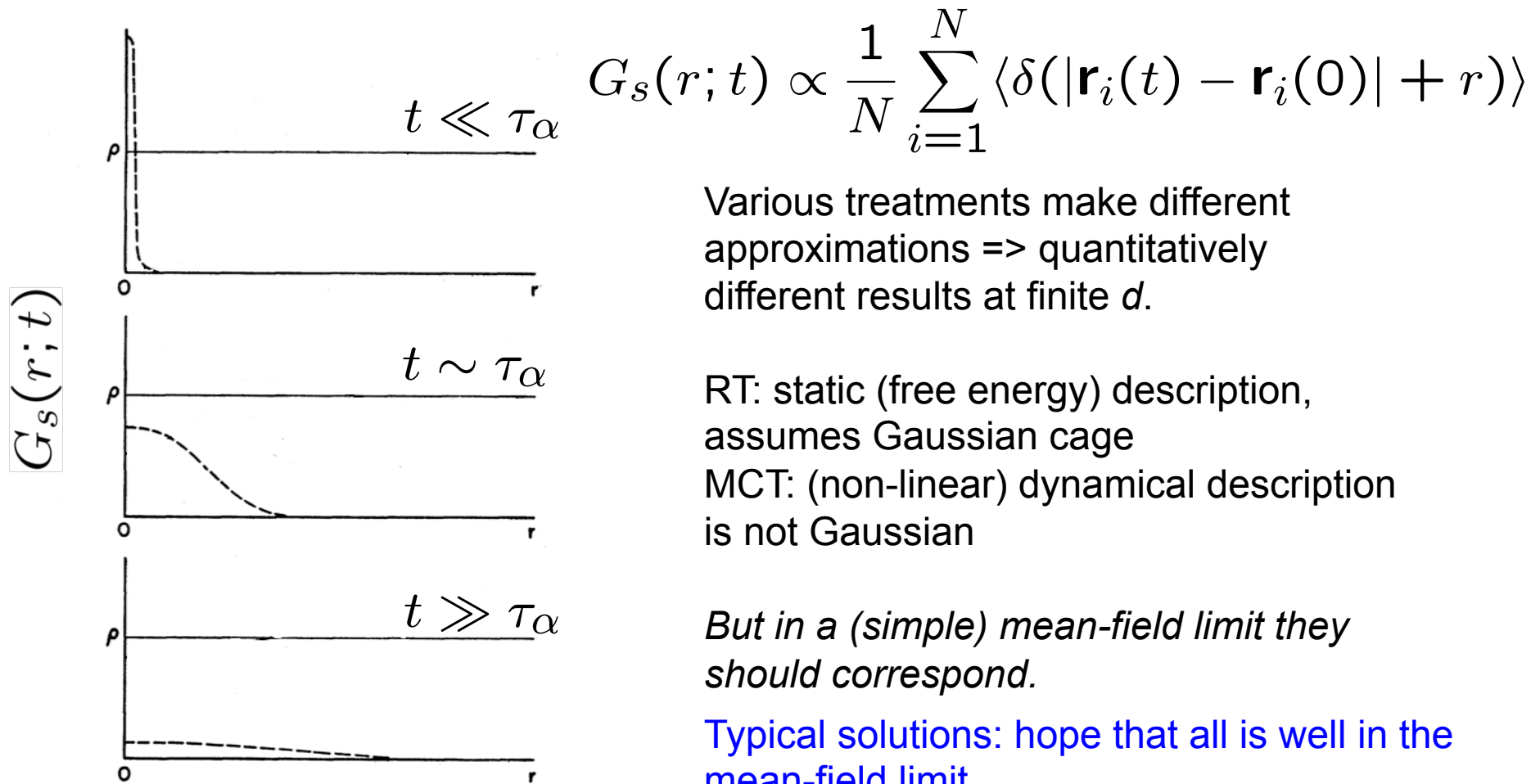


For the RFOT dynamical transition $d_u=8$

Typical solutions: wave hands or throw everything away

Other possible solution: increase spatial dimension

Problem 3: Caging order parameter is non-trivial to describe



$$G_s(r; t) \propto \frac{1}{N} \sum_{i=1}^N \langle \delta(|\mathbf{r}_i(t) - \mathbf{r}_i(0)| + r) \rangle$$

Various treatments make different approximations => quantitatively different results at finite d .

RT: static (free energy) description, assumes Gaussian cage

MCT: (non-linear) dynamical description is not Gaussian

But in a (simple) mean-field limit they should correspond.

Typical solutions: hope that all is well in the mean-field limit

Other possible solution: increase spatial dimension to check

Example 1: MF view of MCT

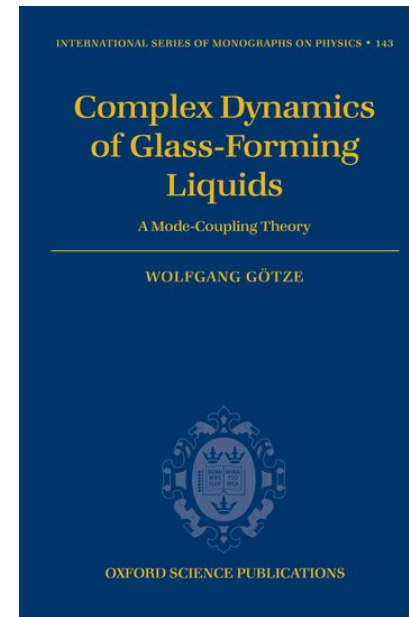
The Mode-Coupling Theory of supercooled liquids

Does it wear any clothes?

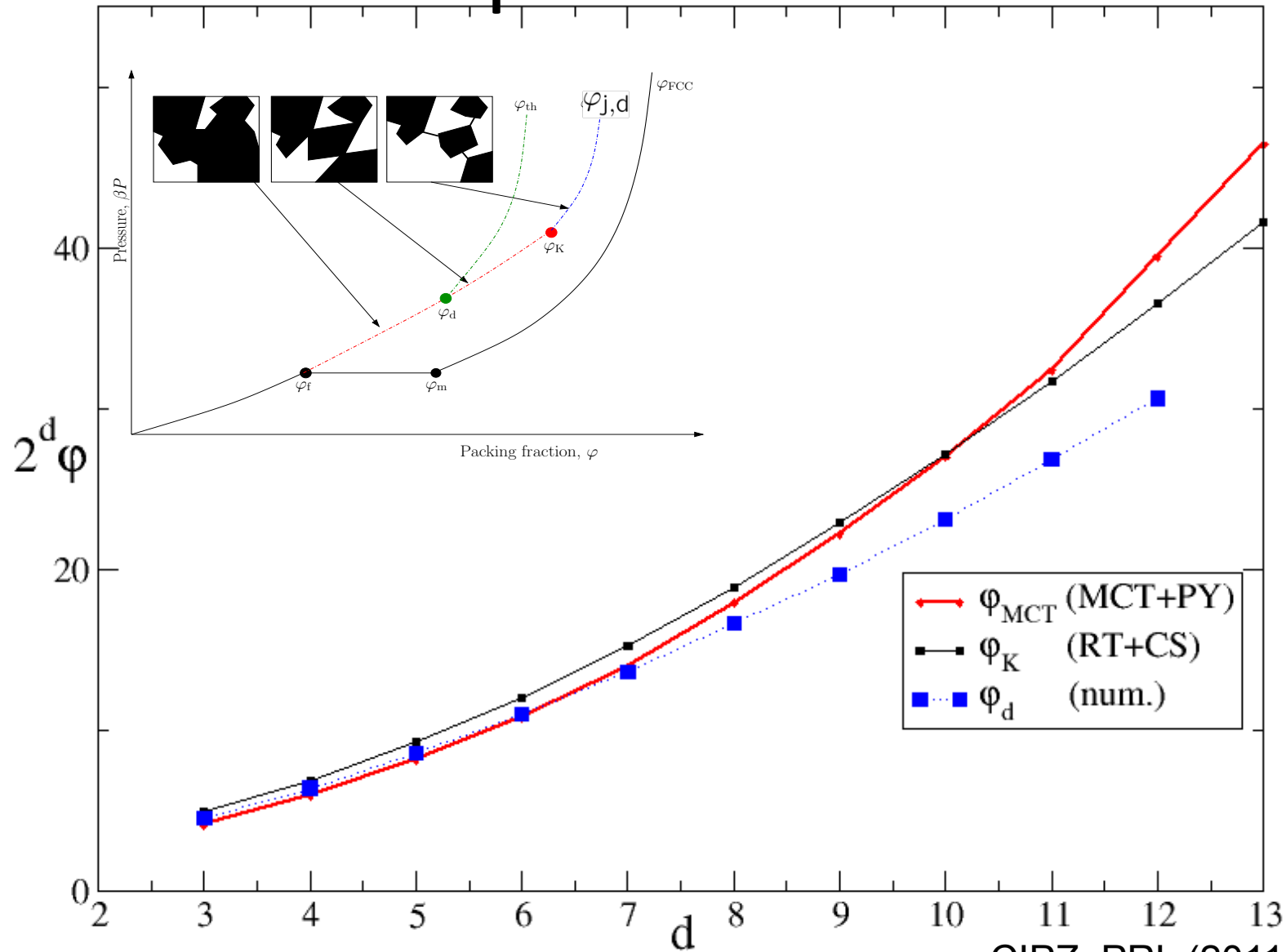
$$\varphi_{RT} \sim 4.8d2^{-d}$$
$$\varphi_{MCT} \sim 0.22d^22^{-d}$$

Glass transition of hard spheres in high dimensions,
Authors: Bernhard Schmid, Rolf Schilling
arXiv:1003.4559.

Mode-Coupling Theory as a Mean-Field Description of the Glass Transition
Authors: Atsushi Ikeda, Kunimasa Miyazaki
arXiv:1003.5472



MF problem with MCT

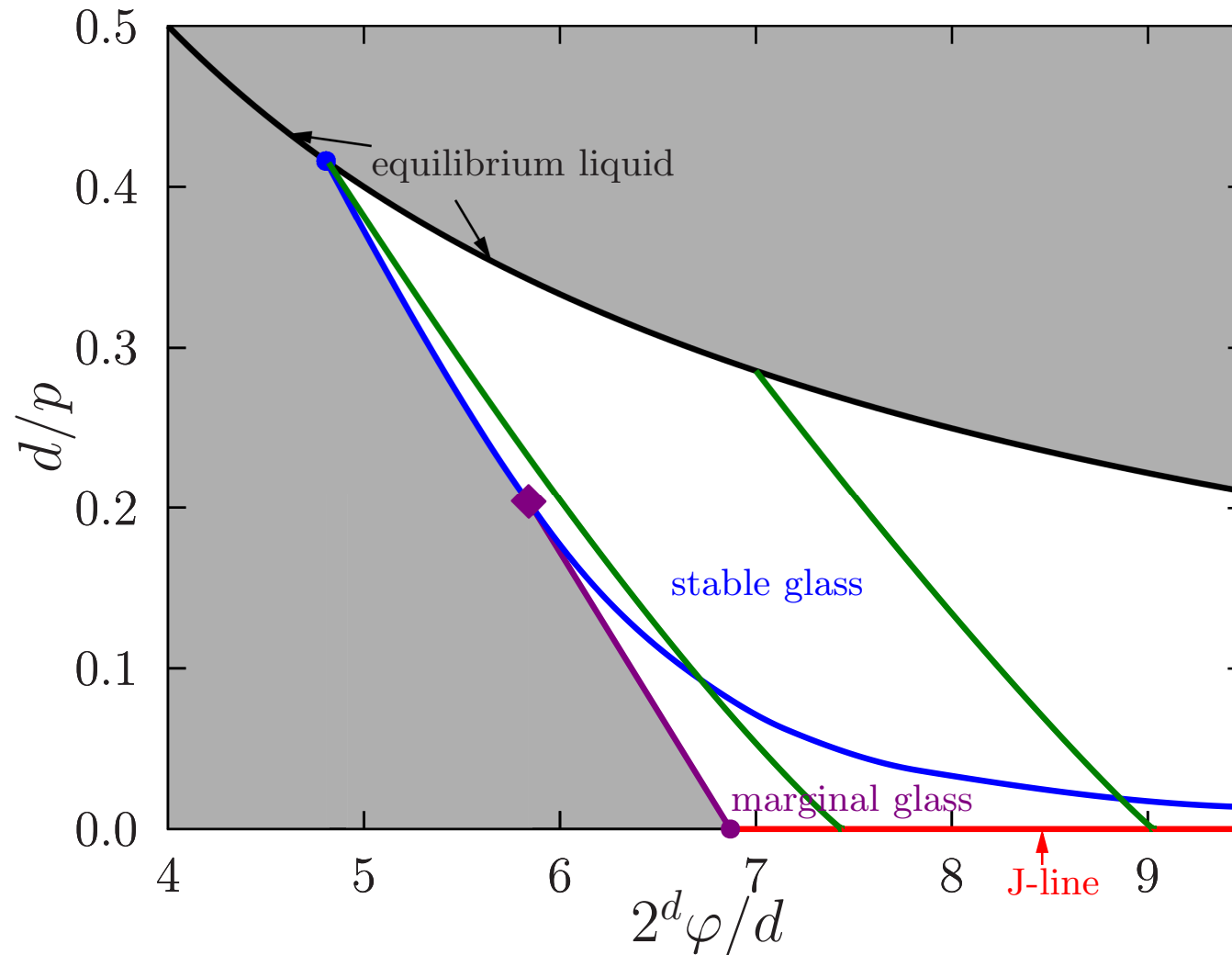


Example 2: Jamming

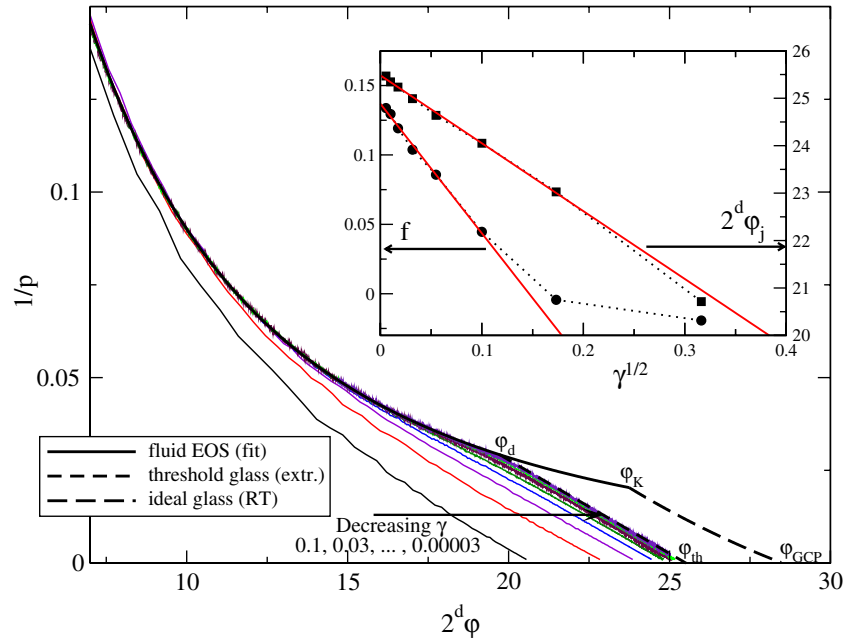
Out-of-equilibrium critical transition, hence describing it requires a good microscopic glass theory.

->Stringent test of glass theories.

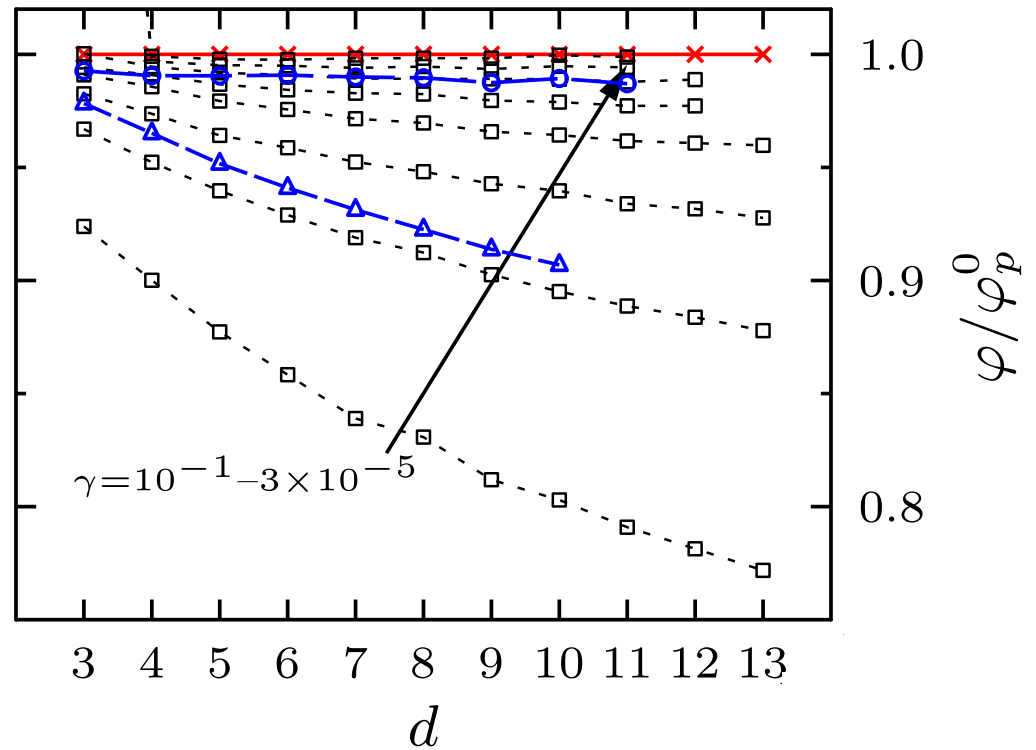
HS phase diagram (GP's talk tomorrow)



Evolving Density



Consistent with $d=3$ results
Chaudhuri, Berthier, Sastry PRL (2010)



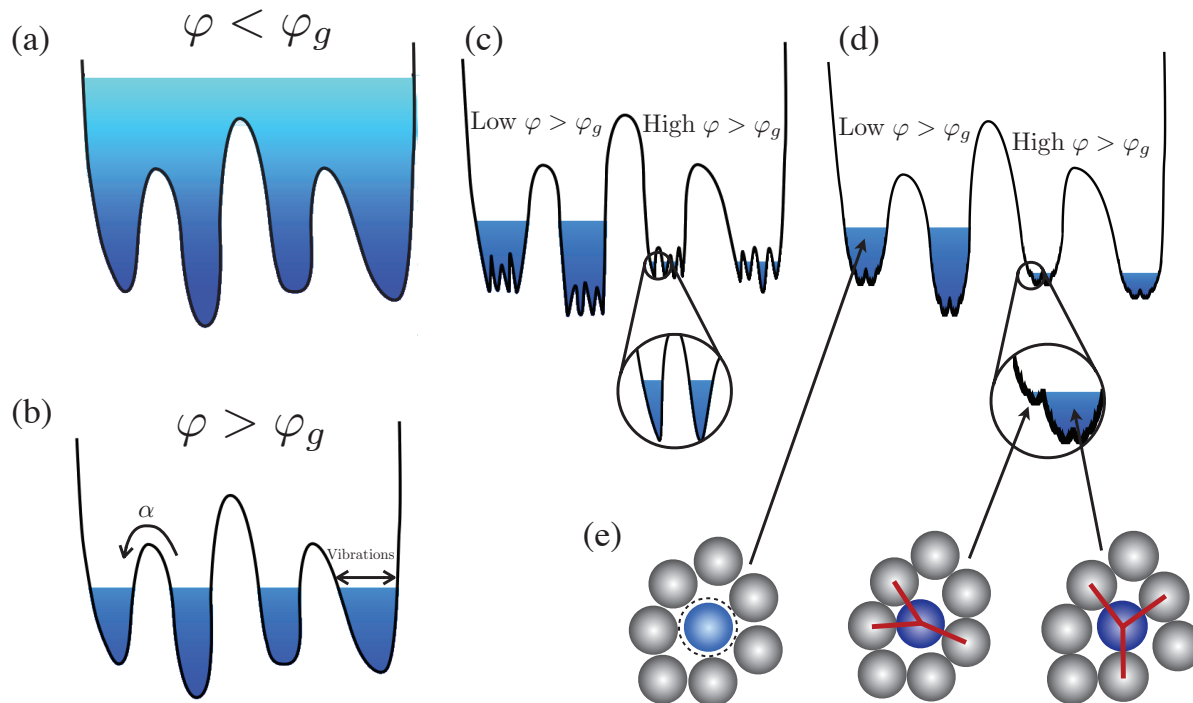
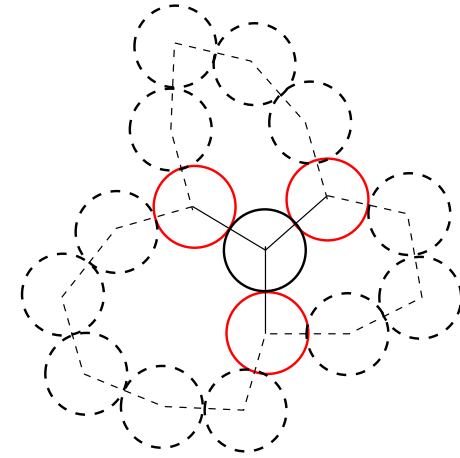
CIPZ PRL (2011); CCPZ PRL (2012)

Cage Collapse

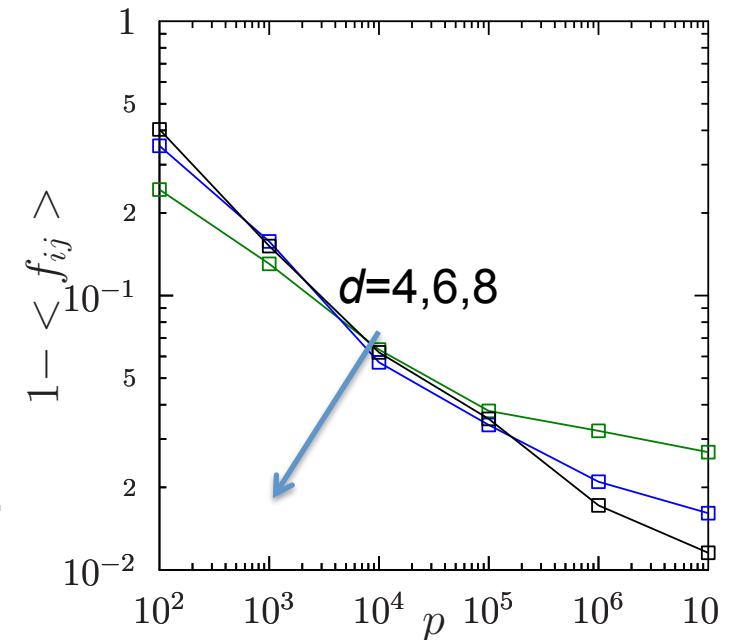
Between the liquid and jamming, something must happen, because

$$\bar{z} = 2d \text{ (isostaticity)}$$

$$\bar{z} \sim e^d \text{ (liquid shell)}$$



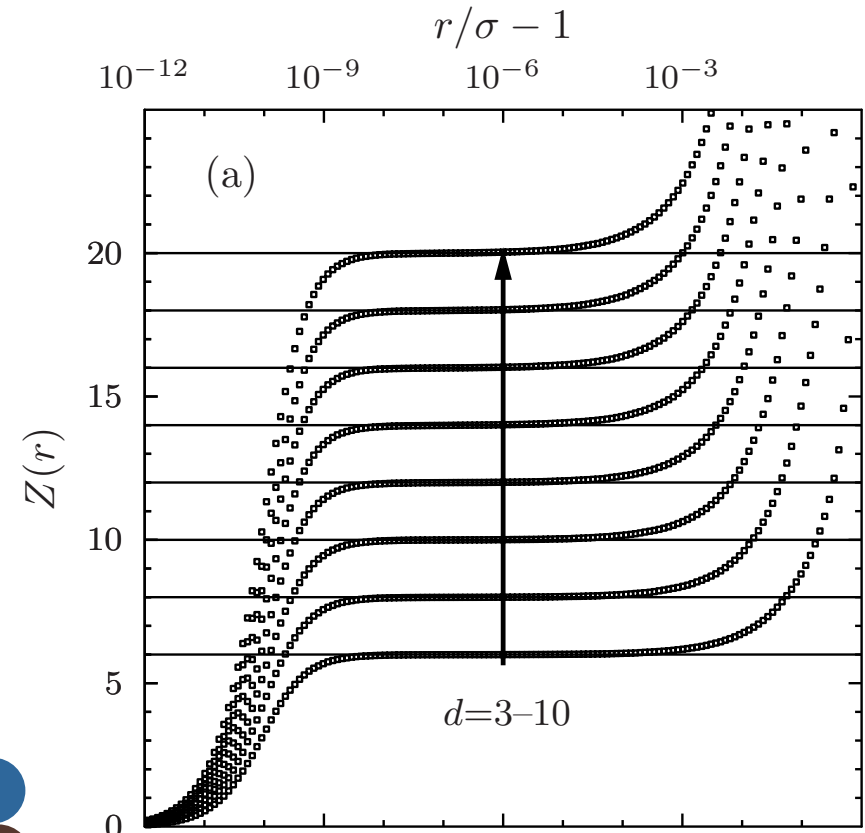
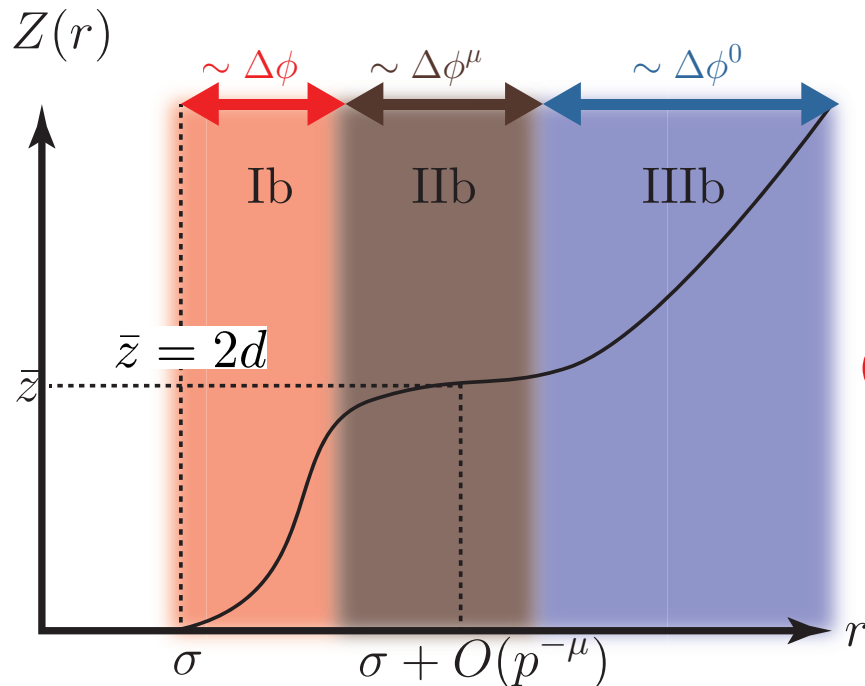
How determined are the force contacts (and the rattlers)?



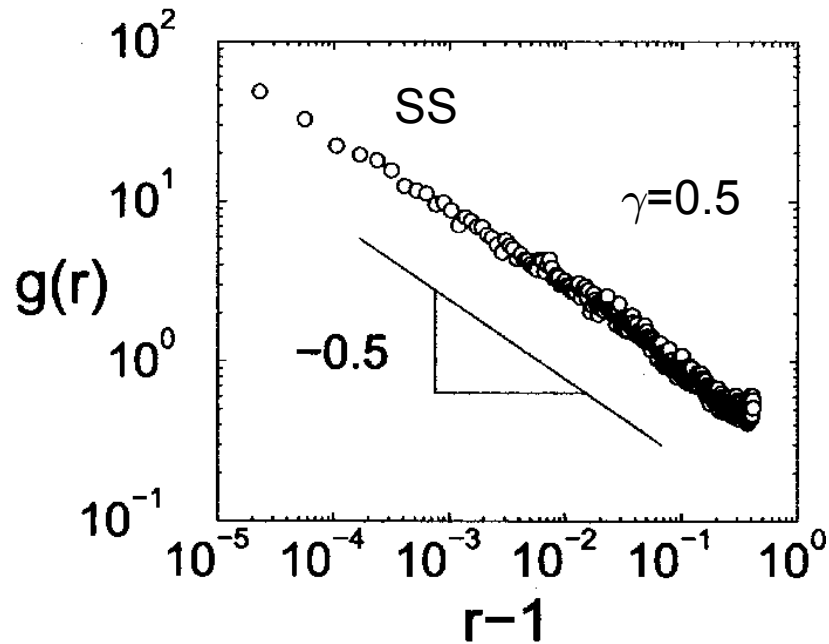
Jamming (marginally stable) cage

As proposed by Wyart PRL (2012), two critical regimes follow from marginality.

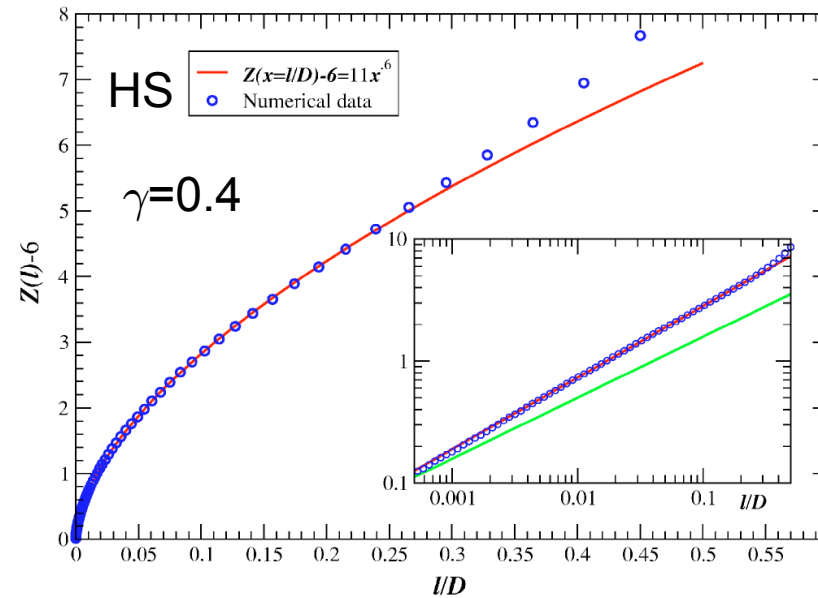
1RSB solution does not: power-law exponents are 0.



1st power law: near contacts

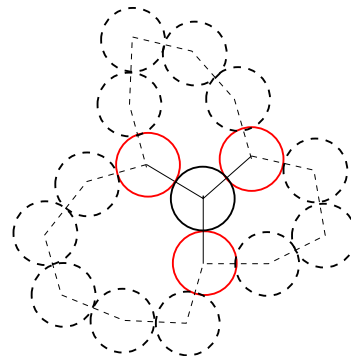
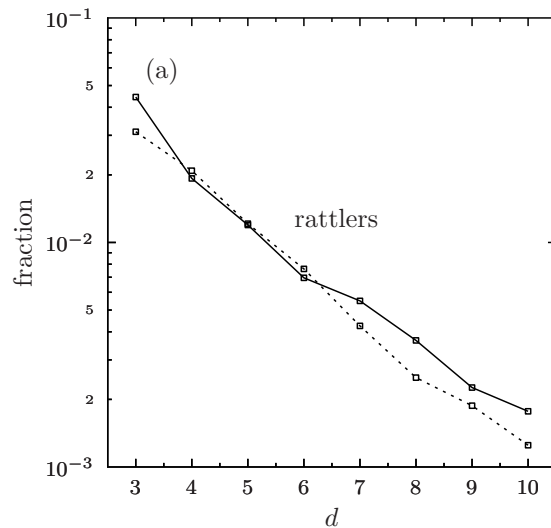
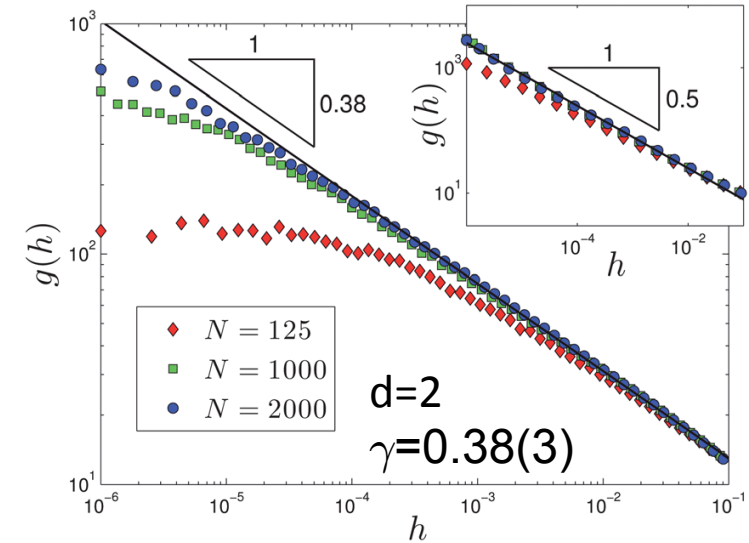
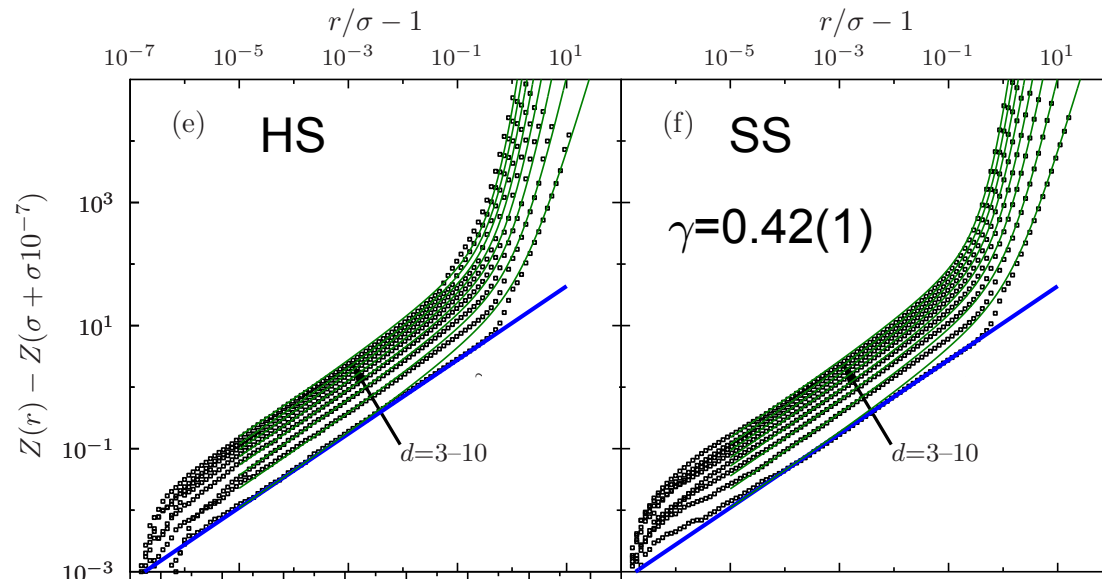


Silbert, Liu, Nagel, PRE (2005)



Donev, Stillinger, Torquato, PRE (2005);
Skoge *et al.* PRE (2006)

Near contacts: remove rattlers

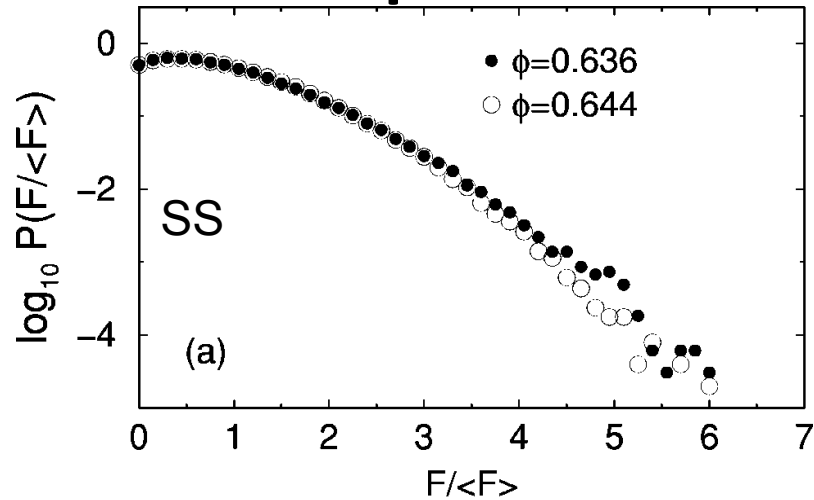


fullRSB $\gamma=0.41269$

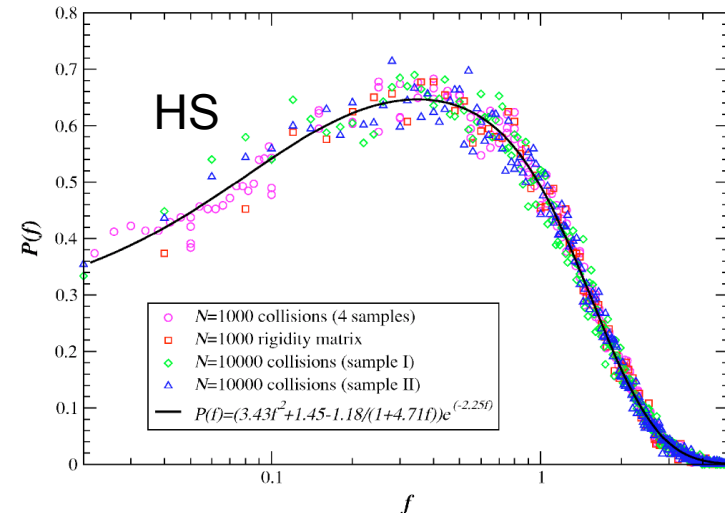
Suggest $d_u=d_l=2$ (?)

Agrees with finite-size scaling arguments of Goodrich and Liu PRL (2012).

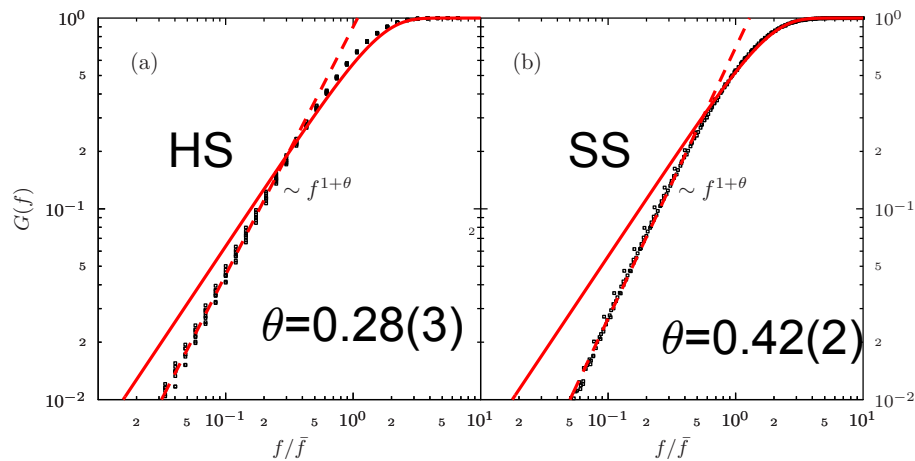
2nd power law: force contacts



O'Hern et al. PRE (2003)

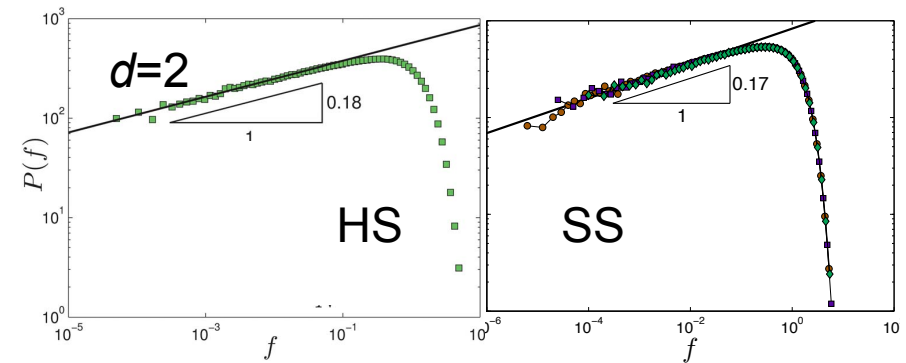


Donev, Stillinger, Torquato, PRE (2005)



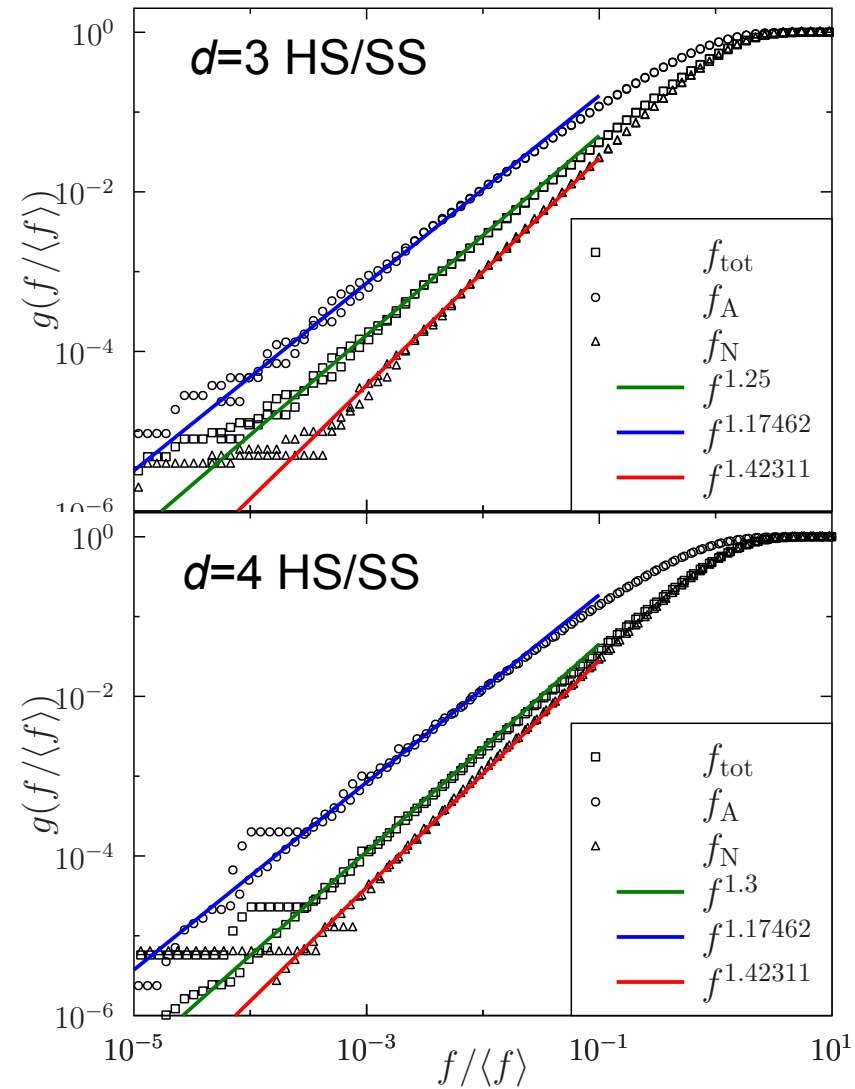
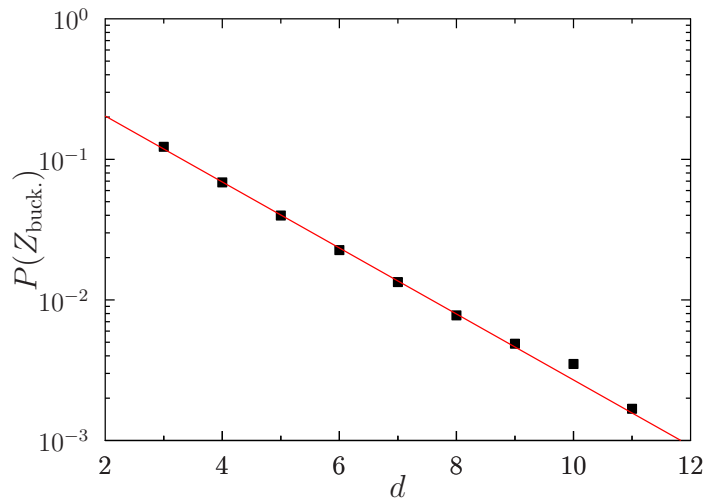
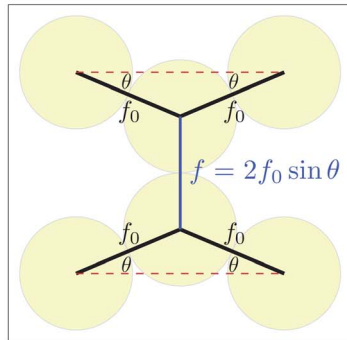
CCPZ PRL (2012)

Embarrassingly different...

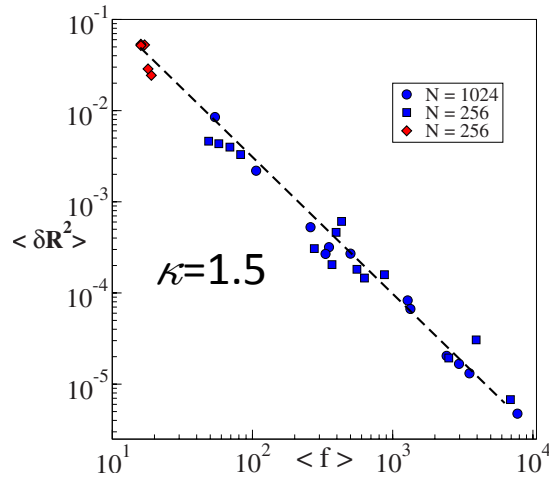


Lerner, Duering, Wyart Soft Matter (2013); DeGiuli et al. arXiv:1402.3834

Force contacts: remove bucklers



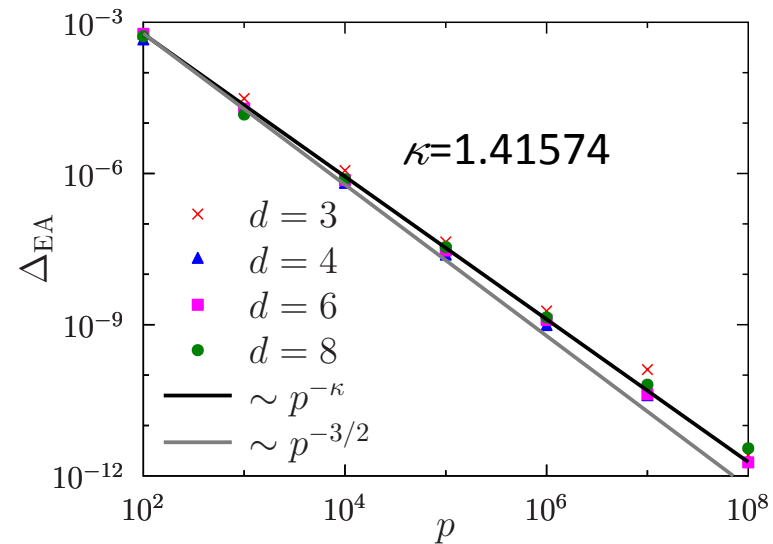
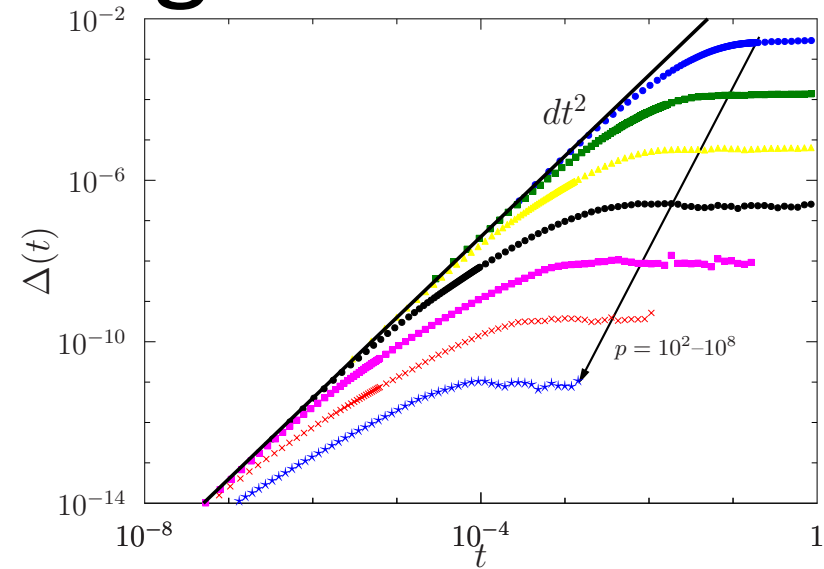
3rd power law: cage evolution



Brito and Wyart J. Chem. Phys. (2009)

1RSB solution gives 0.

High pressure numerical challenges:
 -eliminate rattlers
 -network and thus rattlers reorganize with time



CKUPZ Nat. Comm. (2014)

Problem 4: Can only approach dynamical transition from one side

Typical solutions: get out of equilibrium and pray

Other possible solution: High dimensions can't help much, so... plant a glass instead!

Example 3: Mari-Kurchan (MK) model

HS interacting via shifted potential

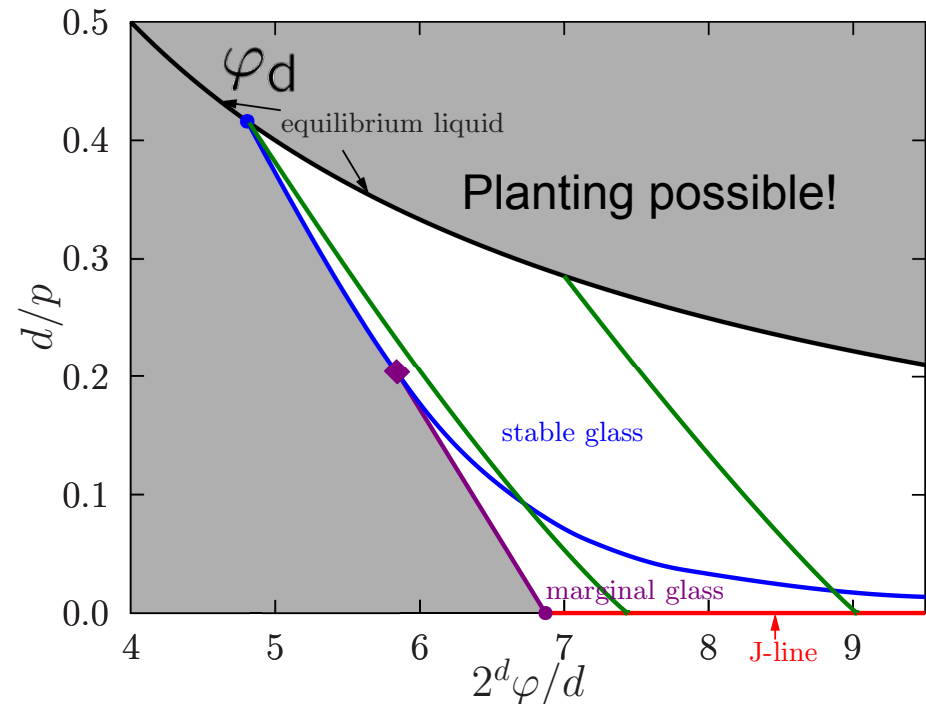
$$H_\Lambda = \sum_{j>i=1,N} V(|x_i - x_j + \Lambda_{ij}|)$$

$\Lambda_{ij} \in [0, L]^d$ with a flat distribution. **If A is near B and B is near C, A is not necessarily near C.** In the liquid phase

$$g(x) = \exp(-\beta V(x))$$

i.e., $p = \beta P / \rho = 1 + \beta B_2 \rho$

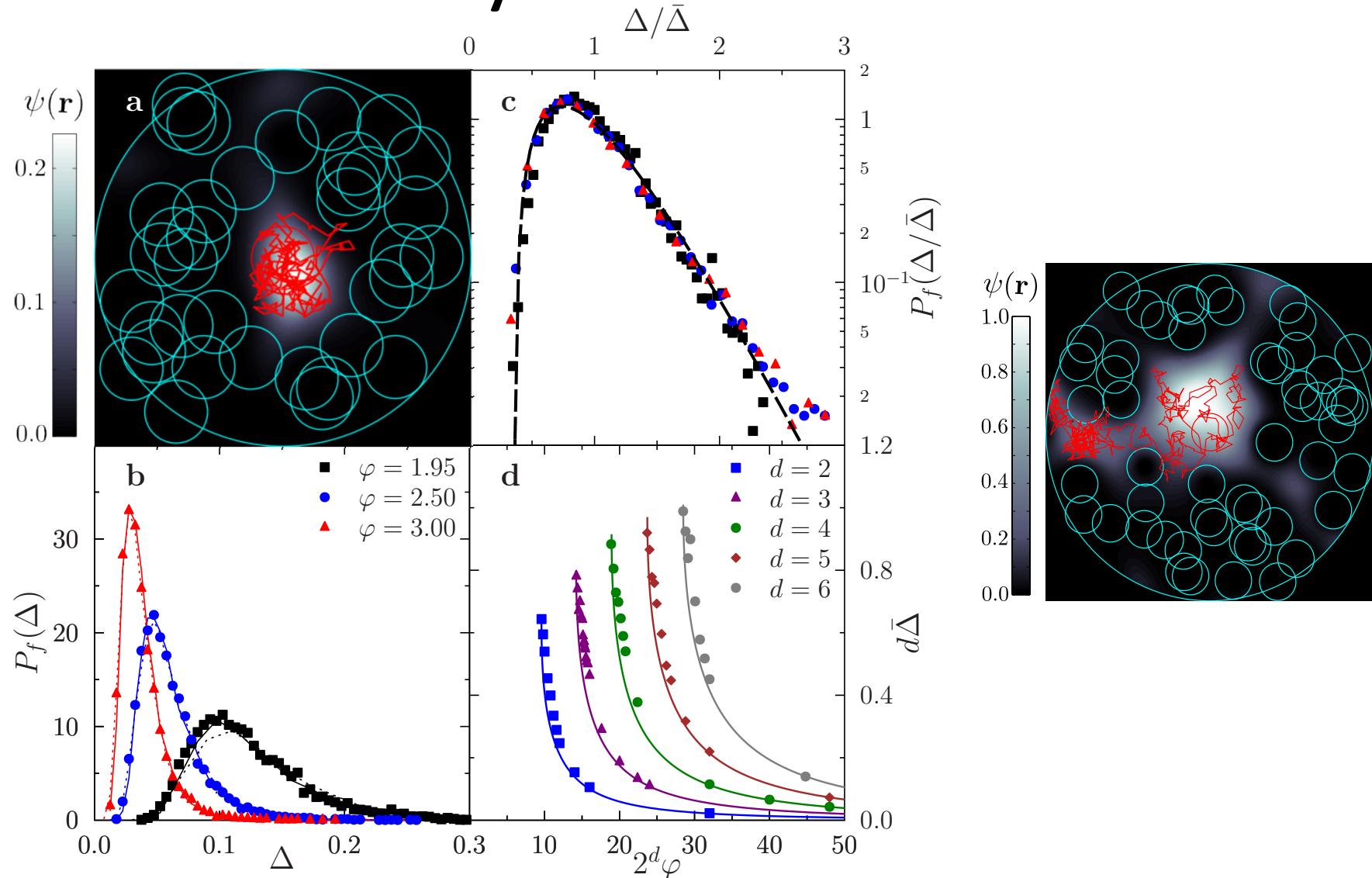
- Same high d partition function as HS.
- Mean-field critical behavior.
- Nucleation suppressed.
- Critical fluctuations don't couple with displacements.
- No spatial notion \rightarrow limited facilitation.
- **Ideal off-lattice finite d MF model!**



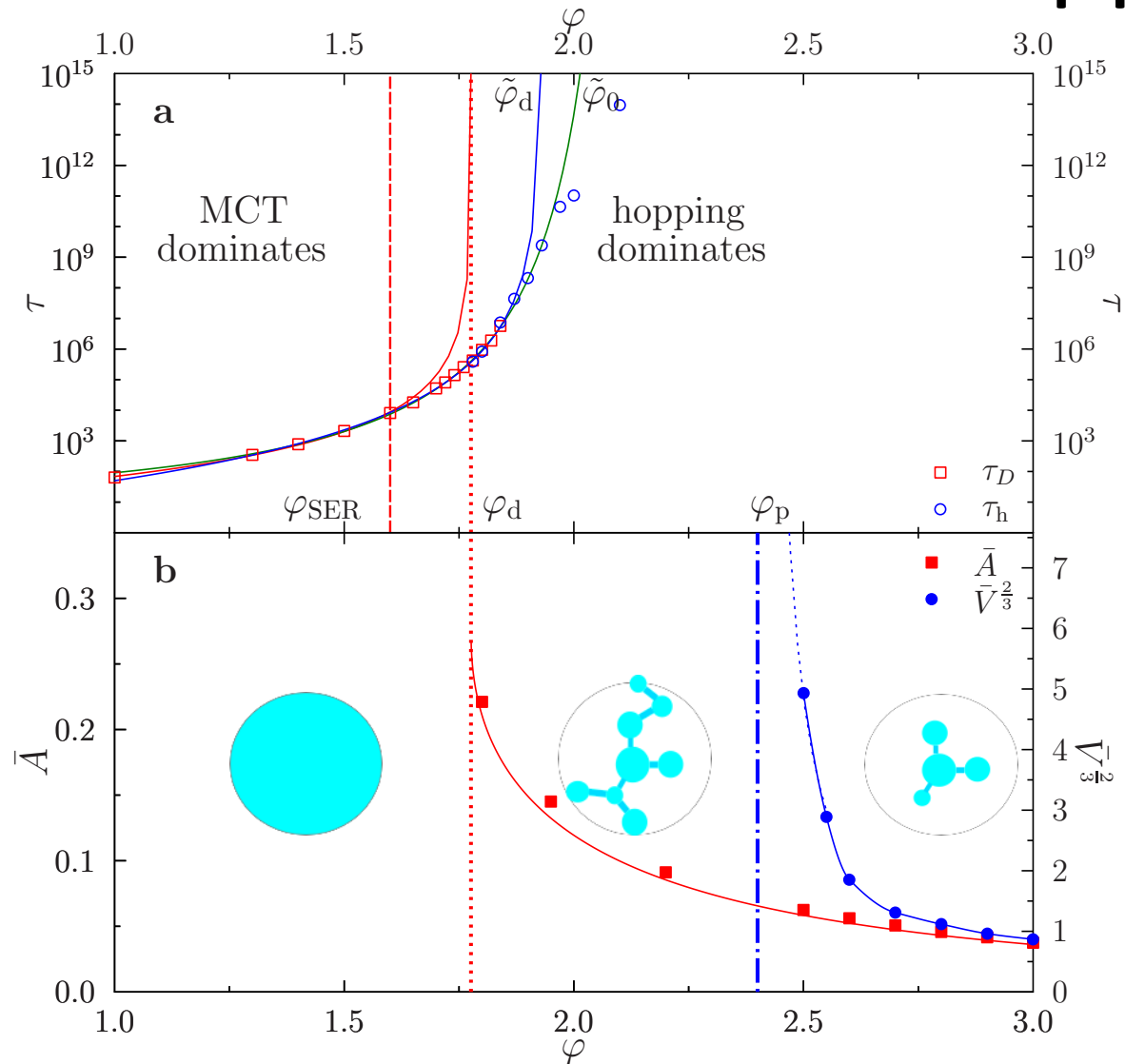
MK, PRL (2009); JCP (2011)

Mézard, Parisi, Tarzia, Zamponi, J. Stat. Phys. (2011)

Cavity reconstruction



Finite d MF \rightarrow RFOT + hopping



Conclusions

- Some qualitative *and quantitative* features of the mean-field fullRSB solution persist all the way down to $d=2$. How ubiquitous is fullRSB?
- Refinement of critical exponents and of deviations is ongoing (robust?).
- Gardner transition may be experimentally testable (ask Yuliang or Bea).
- Much to do about the dynamical transition and connection with void percolation (Jin and PC, arXiv:1409.0688)