

Interfacing Many-Body methods and EFT interactions

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June 13th 2017

Doctoral Training Program
ECT*

- ① Reminders on Nuclear interactions
 - Historical review of interactions
 - Generalities about EFT
 - Going to many-body observables

- ② Interfacing *ab initio* many-body methods with EFT
 - Self-Consistent Green's Function for nuclear matter
 - Recent results about neutron matter

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- 1935 : Yukawa potential
 - Birth of Meson Theory
 - 40s -50s : Pions Theories
 - Troubles with multi-pions, anti-nucleons diagrams, ...
 - 60s -00s : One-Boson-Exchange Model ($\rho, \sigma, \omega, \dots$)
 - Good data fitting
 - But no systematic
 - Need separate framework for 3-body interactions
- } No renormalization group invariance
- 90s - today : EFT based interactions [S.Weinberg 90 91]
 - *Pragmatic* view : Break RG invariance + estimate error
[Entem, Machleidt 03] [Epelbaum, Glöckle, Meißner 05]
 - *Canonical* view : Modify power counting and adapt Many-body methods
[Kaplan, Savage, Wise 96] [Nogga, Timmermans, van Kolck 05]

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The quantum field theory generated by the most general Lagrangian with some assumed symmetries will produce the most general S matrix incorporating quantum mechanics, Lorentz invariance, unitarity, cluster decomposition and the assumed symmetries. [S.Weinberg 79]

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- ▷ Degrees of Freedom
 - Nucleons
 - Pions
 - ...
- ▷ Symmetries
 - Space-time
 - Internal
- ▷ LECs → Fit to data

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Advantages

- Direct connection to underlying theory
- Hierarchy of terms in potential
- Same framework for all A -body interactions
- Estimations of uncertainty and range of validity

Regularization procedure needed

- Regulator = modify high energy physics
- Introduce an arbitrary cut-off scale Λ

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Matching procedure

- Compute observables $O(c_i(\Lambda), \Lambda)$ at a given order N
- Fit LECs to experimental data or to the underlying theory
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Renormalization features

- High energy physics fully included in LECs
- Independence of the observables from the regularization

- Low energy observables in nuclear systems : $Q \ll m_\pi$

Degrees of freedom : non-relativistic nucleons (+ photons, ...)

$$\rightarrow \mathcal{L}_{\pi EFT} = N^\dagger \left(i\partial_0 + \frac{\vec{\nabla}^2}{2m_N} \right) N + C_0 (N^\dagger N)^2 + D_0 (N^\dagger N)^3 + \dots$$

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Truncation scheme

→ What diagrams to compute for N^{th} order ?

- Separation of scale
 - Low energy observable Q
 - Breakdown scale M

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Consistency check

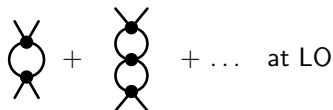
- Convergent observables in negative power of Λ
 - Observable \perp higher energy scale (except through LECs)
- Convergence of observable to experiment when N^{th} increases
 - Fail? Try another guess on LECs
 - No consistent power counting works? might be wrong D.o.F

- Guess on LECs size \rightarrow Proposed power counting for $\not\approx$ EFT

[Van Kolck 97] [Kaplan, Savage, Wise 98] [Bedaque, Hammer, van Kolck 98 99] ...

	2-body	3-body	4-body
LO			X
NLO		X	?
N ² LO	X		?

Un-natural scattering length a



- Consistency check : LO $A = 4, 6, 16$

[Contessi, Lovato, Pederiva, Roggero, Kirsner, van Kolck 17]

- & NLO up to $A = 3$

[Vanasse et al. 13][König et al. 16]

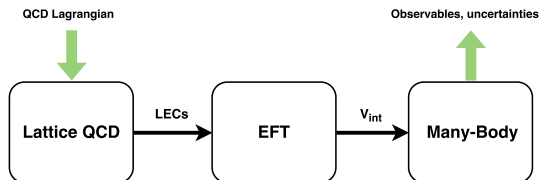
Need to extend consistency check to general A -body observables!

This is where *ab initio* many-body methods enter the game

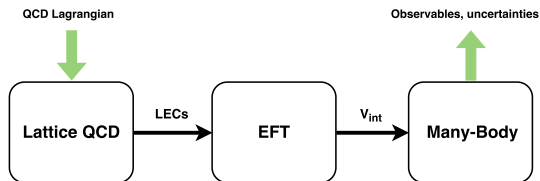
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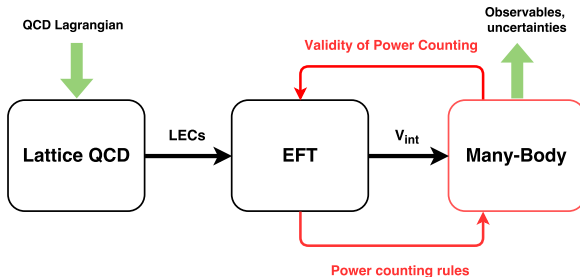
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- Adapt many-body scheme to assess proposed power counting



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Choose a split of the hamiltonian $H \equiv H_0 + V$ and consider (un)correlated states, related by the Lippmann-Schwinger equation

$$\begin{aligned} H_0|\phi\rangle &\equiv E|\phi\rangle \\ H|\psi\rangle &\equiv E|\psi\rangle \\ V|\psi\rangle &\equiv T|\phi\rangle \end{aligned} \quad \Rightarrow \quad \begin{aligned} |\psi\rangle &= |\phi\rangle + \frac{1}{E - H_0} V|\psi\rangle \\ T &= V + \frac{1}{E - H_0} T \end{aligned}$$

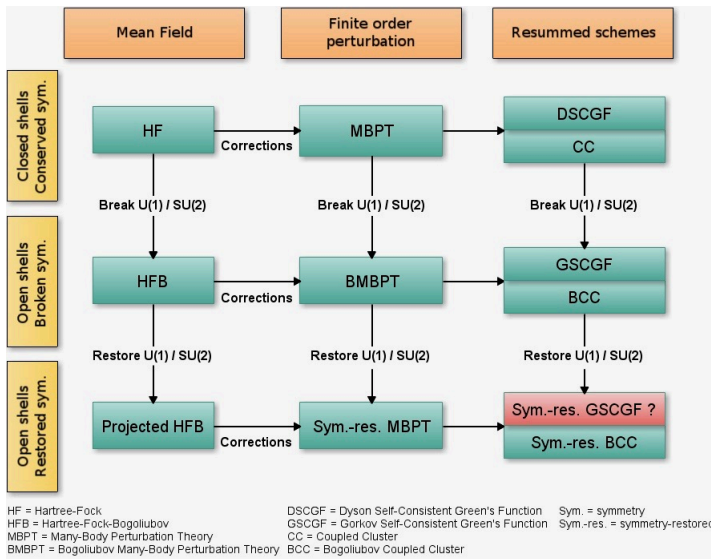
Thus we get the perturbation expansion

$$T = V + V \frac{1}{E - H_0} V + V \frac{1}{E - H_0} V \frac{1}{E - H_0} V + \dots$$

Main features

- Choice of $H_0 \sim$ From what state to start the expansion
- Order of expansion

Brief overview of many-body perturbation methods



Courtesy of P. Arthuis, T. Duguet

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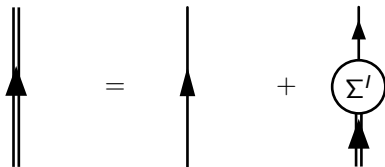
$$G_{\alpha\beta}(t, t') \equiv -i \frac{\langle \Psi_0^N | \mathcal{T} [a_\alpha(t) a_\beta^\dagger(t')] | \Psi_0^N \rangle}{\langle \Psi_0^N | \Psi_0^N \rangle}$$

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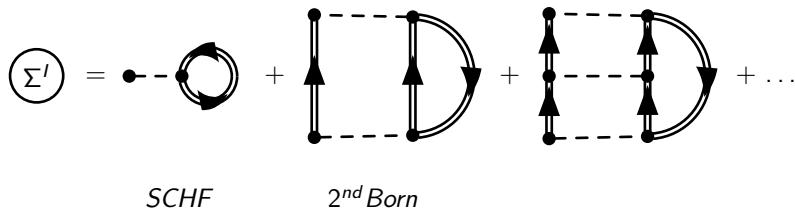
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Dyson equation to be solved self consistently

$$G = G^0 + G^0 \Sigma' G$$



Ladder ansatz for the self-energy

$$\Sigma' = \text{SCHF} + \text{2}^{\text{nd}} \text{Born} + \dots$$


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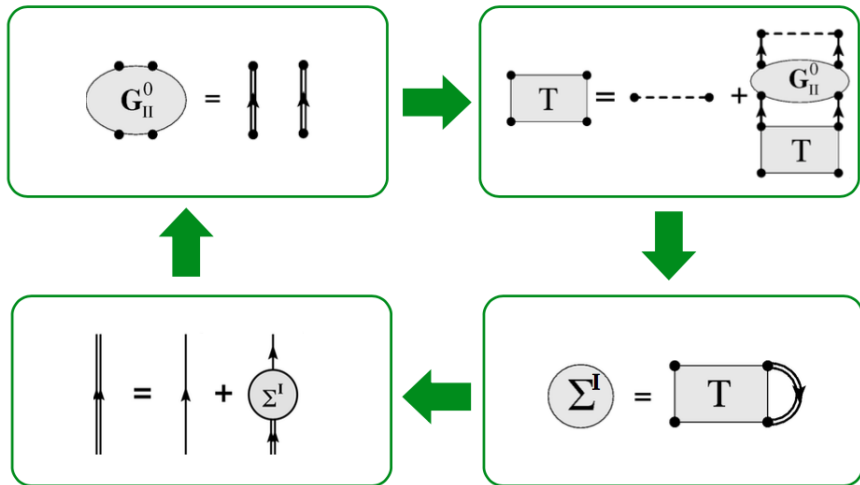
$$\Sigma' = \text{SCHF} + \text{2}^{\text{nd}} \text{Born} + \dots$$

The diagram shows the self-energy Σ' as a sum of terms. The first term is a self-energy loop (Schwinger-Dyson equation) labeled *SCHF*. The second term is a ladder diagram with two rungs, labeled *2nd Born*. The third term is a ladder diagram with three rungs. The terms are separated by plus signs and followed by an ellipsis.

The self-consistent equation (T-matrix formulation)

$$\Sigma' = T = \text{SCHF} + T$$

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 - Two different types of regulators : δ -shell and Gaussian

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δ -shell

$$V(k, k'; r_c) = C_0(r_c) \frac{\sin(kr_c)}{kr_c} \frac{\sin(k'r_c)}{k'r_c}$$

$$\frac{1}{C_0(r_c)} = \frac{m_N}{4\pi} \left(\frac{1}{a} - \frac{1}{r_c} \right)$$

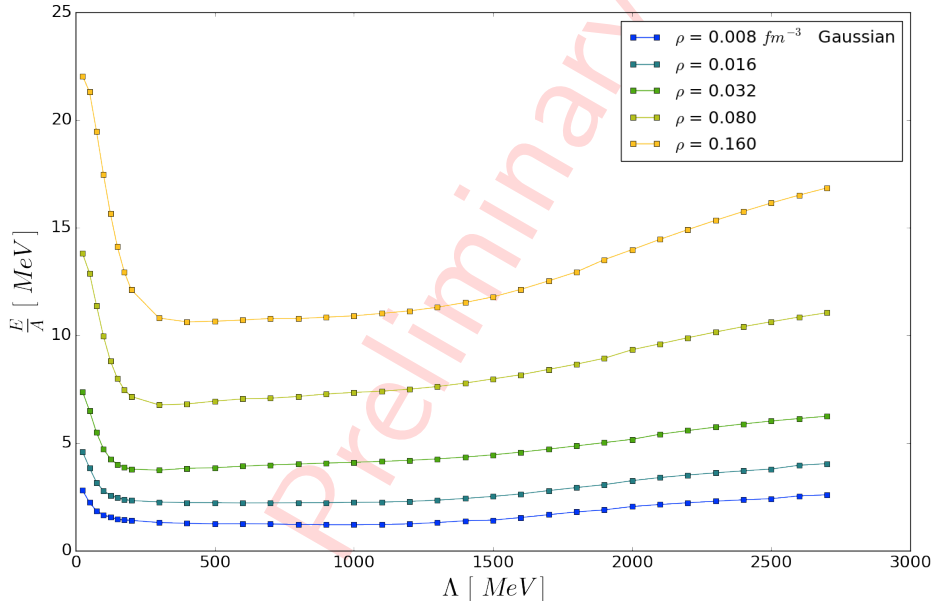
Gaussian

$$V(k, k'; \Lambda) = C_0(\Lambda) \exp\left(-\frac{k^2}{\Lambda^2}\right) \exp\left(-\frac{k'^2}{\Lambda^2}\right)$$

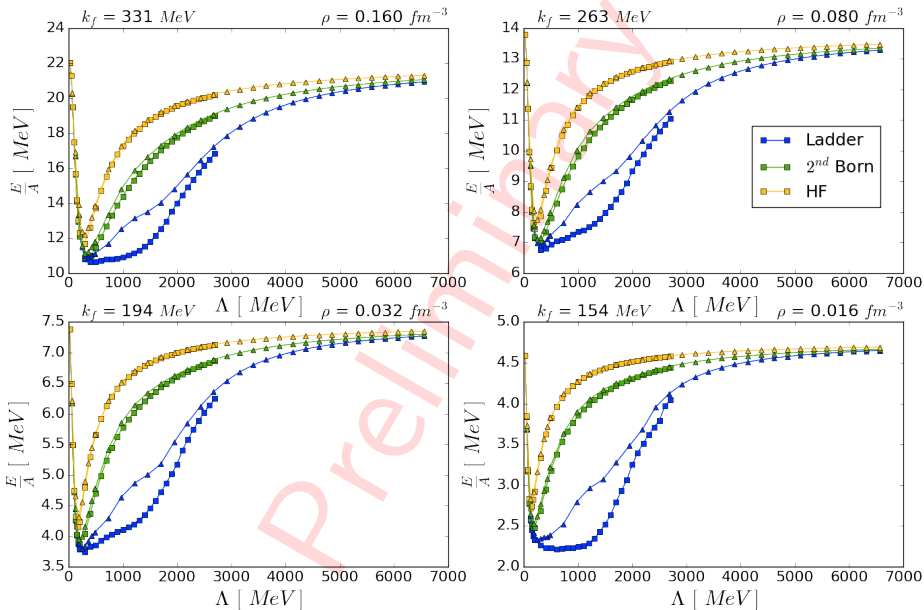
$$\frac{1}{C_0(\Lambda)} = \frac{m_N}{4\pi} \left(\frac{1}{a} - \frac{2}{\pi} \Lambda \int_0^\infty e^{-2x^2} dx \right)$$

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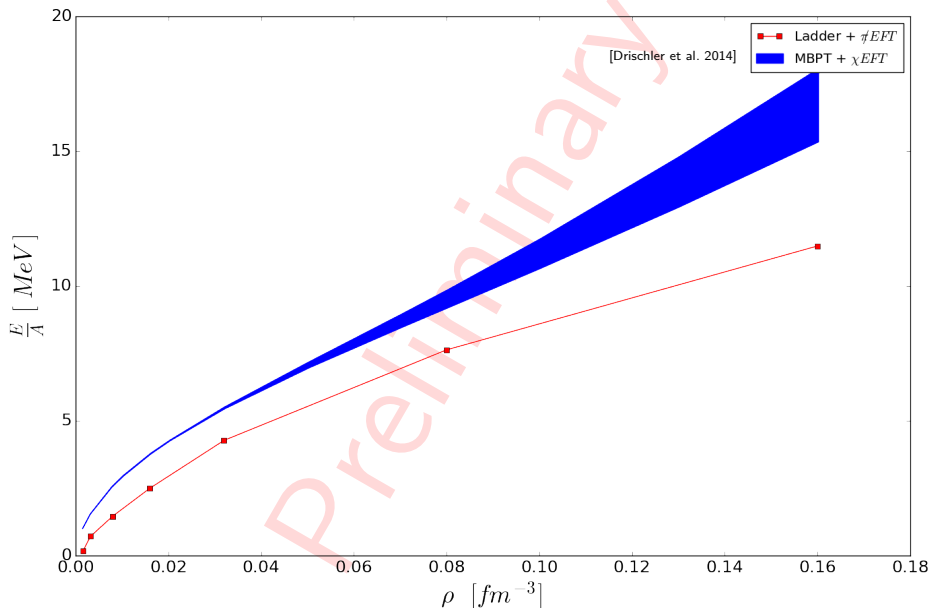
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Comparison of the many-body approaches



Equation of state of neutron matter



▷ Summary

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- Analytical developments
 - Better understand the apparent failure of Ladder approximation
 - Is there a consistent (non-complete) subset of diagrams?
- Exact calculation to benchmark preliminary results

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 - Better understand the apparent failure of Ladder approximation
 - Is there a consistent (non-complete) subset of diagrams?
- Exact calculation to benchmark preliminary results
- Extend study to symmetric matter/NLOs
- What about error estimations?
 - *Still a debated issue*

Thank you !



Advisor : V. Somà

Co-advisor : T. Duguet

B. Bally and P. Arthuis



Collaborators : U. Van Kolck, M. Pavon Valderrama
J. Yang

