# Recent developments in Bogoliubov Many-Body Perturbation Theory 

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## Outline

(1) On ab initio methods and symmetry breaking
(2) On Bogoliubov Many-Body Perturbation Theory
(3) Recent progress

- Validation of the formalism
- First calculations
- To higher orders


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## Evolution of the ab initio reach



Courtesy of V. Soma, T. Duguet
"Exact" ab initio methods

- Since the 80 's
- GFMC, NCSM, FY


## Evolution of the ab initio reach



Courtesy of V. Soma, T. Duguet
Ab initio approaches for closed-shell nuclei

- Since the 2000's
- DSCGF, CC, IMSRG


## Evolution of the ab initio reach



Courtesy of V. Soma, T. Duguet
Non-perturbative ab initio approaches for open-shell nuclei

- Since the 2010's
- GSCGF, BCC, MR-IMSRG


## Evolution of the ab initio reach



Courtesy of V. Soma, T. Duguet
Ab initio shell model

- Since 2014
- Effective interaction via CC/IMSRG


## What makes a method ab initio

(1) Consider point-like nucleons as appropriate degrees of freedom
(2) Use interactions rooted in underlying theory (i.e. QCD)
(3) Expand the many-body Schrödinger equation systematically
(4) Truncate at a given order and solve using computational methods
(5) Estimate systematic error

## On symmetry breaking

Symmetry breaking helps incorporating non-dynamical correlations:

- Superfluid character: $U(1)$ (particle number)
- Deformations: $\operatorname{SU}(2)$ (angular momentum)

But nuclei carry good quantum numbers (e.g. number of particles)
$\Rightarrow$ Symmetries must eventually be restored


## Quantum many-body methods



Expansion methods around unperturbed product state

## Quantum many-body methods



MBPT: Recently (re)implemented with SRG-evolved $\mathrm{H}_{\text {[Tichai et al. 2016] }}$ GSCGF, BCC: Recently proposed and implemented [Somà et al. 2011, Signoracci et al. 2014] Sym.-res. BCC \& sym.-res. BMBPT: Recently proposed [Duguet 2015, Duguet \& Signoracci 2016]

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## Bogoliubov Many-Body Perturbation Theory

(1) Use a Bogoliubov vacuum $|\Phi\rangle$ with $\beta_{k}|\Phi\rangle=0$ for all $k$
(2) Define grand potential operator $\Omega$ from chiral interaction

$$
\Omega \equiv H-\lambda A
$$

then normal-order and split: $\Omega=\Omega_{0}+\Omega_{1}$
(3) Define evolved state in imaginary time

$$
|\Psi(\tau)\rangle \equiv \mathcal{U}(\tau)|\Phi\rangle=e^{-\tau \Omega_{0}} T e^{-\int_{0}^{\tau} d \tau \Omega_{1}(\tau)}|\Phi\rangle
$$

(4) Expand and truncate the grand potential kernel $\Omega(\tau) \equiv\langle\Psi(\tau)| \Omega|\Phi\rangle$...
...and the norm kernel $N(\tau) \equiv\langle\Psi(\tau) \mid \Phi\rangle$
(5) Extract ground state energy via

$$
\mathrm{E}_{0}=\lim _{\tau \rightarrow \infty} \frac{\Omega(\tau)}{N(\tau)}=\lim _{\tau \rightarrow \infty} \omega(\tau)
$$

## Expansion of the grand potential kernel

Inserting the operator $\Omega$ at time 0 and expanding

$$
\begin{aligned}
\mathrm{E}_{0}= & \lim _{\tau \rightarrow \infty} \frac{\langle\Psi(\tau)| \Omega|\Phi\rangle}{\langle\Psi(\tau) \mid \Phi\rangle} \\
= & \langle\Phi|\left\{\Omega(0)-\int_{0}^{\infty} d \tau_{1} \mathrm{~T}\left[\Omega_{1}\left(\tau_{1}\right) \Omega(0)\right]\right. \\
& \left.+\frac{1}{2!} \int_{0}^{\infty} d \tau_{1} d \tau_{2} \top\left[\Omega_{1}\left(\tau_{1}\right) \Omega_{1}\left(\tau_{2}\right) \Omega(0)\right]+\ldots\right\}|\Phi\rangle_{c}
\end{aligned}
$$

Then expressing the grand potential in the qp basis

$$
\Omega=\Omega^{00}+\frac{1}{1!} \sum_{k_{1} k_{2}} \Omega_{k_{1} k_{2}}^{11} \beta_{k_{1}}^{\dagger} \beta_{k_{2}}+\frac{1}{2!} \sum_{k_{1} k_{2}}\left\{\Omega_{k_{1} k_{2}}^{20} \beta_{k_{1}}^{\dagger} \beta_{k_{2}}^{\dagger}+\Omega_{k_{1} k_{2}}^{02} \beta_{k_{2}} \beta_{k_{1}}\right\}+\ldots
$$

## Expansion of the grand potential kernel

$$
\begin{aligned}
& \mathrm{E}_{0}=\sum_{p=0}^{\infty} \frac{(-1)^{p}}{p!} \sum_{i_{0}+j_{0}=2,4} \int_{0}^{\infty} d \tau_{1} \ldots d \tau_{p} \\
& i_{p}+j_{\rho}=2,4
\end{aligned}
$$

$$
\begin{aligned}
& \begin{array}{c}
I_{1} \ldots I_{i_{p}} \\
I_{i_{p}+1} \ldots I_{i_{p}+j_{p}}
\end{array} \\
& \times\langle\Phi| \mathrm{T}\left[\beta_{k_{1}}^{\dagger}\left(\tau_{1}\right) \ldots \beta_{k_{k_{1}}}^{\dagger}\left(\tau_{1}\right) \beta_{k_{k_{1}+j_{1}}}\left(\tau_{1}\right) \ldots \beta_{k_{k_{1}+1}}\left(\tau_{1}\right) \ldots\right. \\
& \ldots \beta_{l_{1}}^{\dagger}\left(\tau_{p}\right) \ldots \beta_{l_{p}}^{\dagger}\left(\tau_{p}\right) \beta_{l_{i_{p}+j_{p}}}\left(\tau_{p}\right) \ldots \beta_{l_{p}+1}\left(\tau_{p}\right) \\
& \left.\times \beta_{m_{1}}^{\dagger}(0) \ldots \beta_{m_{i 0}}^{\dagger}(0) \beta_{m_{i_{0}+j_{0}}}(0) \ldots \beta_{m_{i_{0}+1}}(0)\right]|\Phi\rangle_{c}
\end{aligned}
$$

All contributions computable algebraically and diagramatically

## First- and second-order diagrams

Diagrammatic representation of the grand potential $\Omega$


Extracting and applying diagrammatic rules

$$
\mathrm{E}_{0}^{(1+2)}=\quad \begin{gathered}
0 \\
\Omega^{00}
\end{gathered}
$$



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## Third-order diagrams





Derivation of all diagrams up to third order

## Validation of formal derivation

BMBPT must match standard MBPT in Slater determinant limit
$\rightarrow$ Matching must be true at each order
$\rightarrow$ Proof of consistent formalism for BMBPT

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BMBPT(3) diagrams match MBPT(3) ones exactly
Canonical HF-MBPT diagrams were recovered from only one BMBPT


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## Proof of principle calculations

First BMBPT(2) proof of principle calculation of ${ }^{20} \mathrm{O}$ :

using NN SRG-evolved chiral interaction
On MCPT:

- Multi-configurational MBPT
- Alternative method for open-shell nuclei


## Isotopic chains calculations at second order

First $\mathrm{BMBPT}(2)$ calculations on $\mathrm{O}, \mathrm{Ca}, \mathrm{Ni}$ and Sn isotopic chains

using NN and 3N SRG-evolved chiral interaction
Same chains under investigation at third order at the moment

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## Numerical derivation of higher orders

Have your computer do the diagrammatic work for you

- Produce the diagrams automatically
- Diagrams are associated with adjacency matrices
- Diagrammatic rules constrain the form of the matrices
- Have your code generate all matrices, hence diagrams
- Extract their expression automatically as well
- Read your diagrams (vertices, propagators, etc.)
- Extract useful information from the structure
- Retrieve the exact expression


## Numerical derivation of higher orders

Produce higher orders diagrams

- 59 diagrams at order 4
- 568 diagrams at order 5

Extend to three-body diagrams

- 15 diagrams at order 3
- 337 diagrams at order 4
- 10148 diagrams at order 5



## Prospects

- Go up to fourth order
$\rightarrow$ Even better than other ab initio methods?
$\rightarrow$ Test for computational cost
- Push BMBPT to heavier nuclei
$\rightarrow$ Can go further than other ab initio methods
$\rightarrow$ Good test for the computational cost
- Implement particle-number restored BMBPT for the first time
$\rightarrow$ Required for precise study of open-shell nuclei
$\rightarrow$ Proof of concept of symmetry-restored BMBPT / BCC
- Ab initio driven EDF method [T. Duguet et al. (2015)]
$\rightarrow$ Safe/correlated/improvable off-diagonal EDF kernels
$\rightarrow$ Based on PNR-BMBPT
- MBPT and BMBPT are special among ab initio methods
$\checkmark$ Computationally friendlier
$\checkmark$ Potentially as precise as others when using SRG-evolved H
- BMBPT has been formulated and is being implemented
$\checkmark$ First derivation up to fourth order
$\checkmark$ First calculations up to third order
$\checkmark$ Appropriate framework to tackle open-shell nuclei
$\checkmark$ Systematic studies at third and fourth order to come
- Symmetry-restored BMBPT is the next step


## Our collaborators

BMBPT Project

P. Arthuis
T. Duguet
J.-P. Ebran

On broader aspects

$\begin{array}{ll}\text { TECHNISCHE } & \text { A. Tichai } \\ \text { UNIVERSIAAT } \\ \text { DARMSTADT } & \text { R. Roth }\end{array}$

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