Microscopic description of fission mass yields

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June 27, 2017 ECT* DTP, Trento

Motivation



K.H. Schmidt et. al., Nucl. Phys. A665 221 (2000)

- Fission fragment mass yield is a one of the basic, measurable observable
- The shape of observed fragment mass distribution allows to determine the type of fission (symmetric, asymmetric, bimodal)
- Accuracy of reproduction of the experimental mass yields is a test of the theoretical models
- The proper theoretical description of fission mass distribution is still pending



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$$\delta\left[\left\langle \Phi(\boldsymbol{Q}_{20},\boldsymbol{Q}_{30}) \mid \widehat{H} - \lambda_{Z}\widehat{Z} - \lambda_{N}\widehat{N} - \sum_{i}\lambda_{i}\widehat{Q}_{i} \mid \Phi(\boldsymbol{Q}_{20},\boldsymbol{Q}_{30})\right\rangle\right] = 0$$

The quadrupole moment:

$$\widehat{Q}_{20} = \sqrt{\frac{16\pi}{5}} \sum_{i=1}^{A} r_i^2 Y_{20}$$

The octupole moment:

$$\widehat{Q}_{30} = \sqrt{\frac{4\pi}{7}} \sum_{i=1}^{A} r_i^3 Y_{30}$$





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Kształty przedrozszczepieniowe





A_H / A_L = 131 / 121







A_H / A_L = 142 / 110



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Fission dynamics

$$\begin{split} \widehat{H}_{coll} &= -\frac{\hbar^2}{2\sqrt{\gamma}} \sum_{i,j=2}^3 \frac{\partial}{\partial q_{i_0}} \sqrt{\gamma} B_{ij}(Q_{20}, Q_{30}) \frac{\partial}{\partial q_{j_0}} + V(Q_{20}, Q_{30}) \\ \gamma &\equiv \det G(Q_2, Q_3), \qquad B_{ij} = \mathcal{M}_{ij}^{-1} \\ \widehat{H}_{coll}' g_n^{\pi}(Q_2, Q_3, t = 0) = E_n^{\pi} g_n^{\pi}(Q_2, Q_3, t = 0) \end{split}$$





 $\widehat{H}_{\text{coll}}g(Q_2, Q_3, t) = i\hbar \frac{\partial g(Q_2, Q_3, t)}{\partial t}$



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Probability current density

$$egin{aligned} ec{J}(Q_{20},Q_{30},t) &= rac{\hbar}{2i}\sqrt{\gamma}B(Q_{20},Q_{30}) imes \ [g^*(Q_{20},Q_{30},t)
abla g(Q_{20},Q_{30},t)
abla g^*(Q_{20},Q_{30},t)] \end{aligned}$$

$$P(Q_{20}^{sc}, Q_{30}^{sc}) = \int_{t=0}^{t=T^{propag}} \vec{J}(Q_{20}^{sc}, Q_{30}^{sc}, t) \cdot \vec{n} dt$$







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Pre-scission configuration





M. Warda, A. Zdeb, Physica Scripta 90 114003 (2015).



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Neck rupture probability



 $P(z) \sim \exp[-\sigma(z)]$

Brosa U., Phys. Rev. C38 1944 (1988).



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Random Neck Rupture Mechanism

$$\sigma(z) = 2\pi \int_{0}^{\infty} r_{\perp} \rho(z, r_{\perp}) dr_{\perp}$$

°_ 0

$$P(A_1/A_2) = \exp\left[\frac{-2\gamma\sigma(z)}{T}\right] \implies$$





Q₂=238 b



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Brosa U., Phys. Rev. C38 1944 (1988).

The tunneling probability



A. Zdeb, A. Dobrowolski, M. Warda, Phys. Rev. C95 054608 (2017).

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Time of evolution



A. Zdeb, A. Dobrowolski, M. Warda, Phys. Rev. C95 054608 (2017)



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A. Zdeb, A. Dobrowolski, M. Warda, Phys. Rev. C95 054608 (2017)

Parity dependence



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A. Zdeb, A. Dobrowolski, M. Warda, Phys. Rev. C95 054608 (2017)

Mixed states



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A. Zdeb, A. Dobrowolski, M. Warda, Phys. Rev. C95 054608 (2017)

Conclusions and perspectives

- The peaks positions are well reproduced
- There is no strong correlation between mass distribution and initial conditions
- Theoretical mass yields are too narrow in comparison to the experimental ones
- To improve the accuracy of data reproduction model extensions are required



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Third dimension - Q_{40}

The hexadecapole moment:

$$\widehat{Q}_{40} = \sqrt{rac{4\pi}{9}}\sum_{i=1}^{A}r_{i}^{4}Y_{40}$$



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M. Warda, K. Pomorski, J.L. Egido, L.M. Robledo, J. Phys. G: Nucl. Part. Phys. 31, S1555 (2005).

Thank you!







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Searching for pre-scission points





$$\mathcal{D}_{
ho
ho^{,}}=\int\limits_{0}^{\infty}|
ho(ec{r})-
ho'(ec{r})|oldsymbol{d} au$$

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