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# Using Neutrinos as a Probe of the Strong Interaction

## **Neutrino / Anti-neutrino Deep-Inelastic Scattering off of Massive Nuclear Targets**

e-Nucleus XI  
Elba – June, 2010

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Fermilab

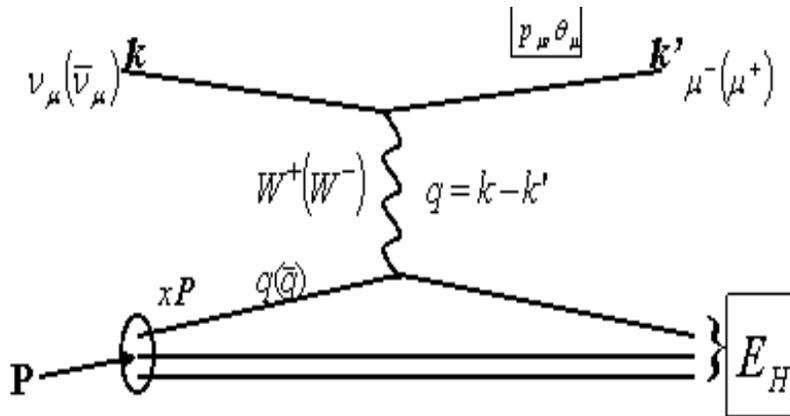
With thanks to the many publications and presentations of Martin Tzanov, U. Colorado

# Motivation for Studying $\nu$ DIS

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- ◆ Interacting with the weak current means a **much smaller interaction rate** than  $e/\mu$  scattering
  - ▼ Need huge, higher-A detectors and intense neutrino beams
  - ▼ The neutrino flux is difficult to predict and measure.
- ◆ However **can select which set of quarks involved in the interaction via  $\nu$  or  $\bar{\nu}$**
- ◆ While  $F_2$  is measured precisely by the charge lepton scattering,  **$xF_3$  is accessible by neutrino DIS.**
  - ▼  $\Delta xF_3$  yields increased sensitivity to the **valence quark distributions.**
  - ▼ However, through  $\Delta xF_3 = 4x(s-c)$ ,  $\Delta xF_3$  is **also sensitive to heavy quarks.**
- ◆ Speaking of heavy quarks, examining charm production with neutrinos also gives us insight into the **strange quark distribution.**
- ◆ **Electroweak physics** has been a rich neutrino subject for decades.
- ◆ **Finally, recent phenomenological / experimental work is indicating some interesting differences concerning nuclear effects with neutrinos compared to charged lepton scattering.**

# The Parameters of $\nu$ DIS



$Q^2 = 4E_\nu E_\mu \sin^2 \frac{\theta}{2}$	Squared 4-momentum transferred to hadronic system
$x = \frac{Q^2}{2ME_{HAD}}$	Fraction of momentum carried by the struck quark
$y = \frac{\nu}{E_\nu} = \frac{E_{HAD}}{E_\nu}$	Inelasticity

Differential cross section in terms of structure functions:

$$\frac{1}{E_\nu} \frac{d^2\sigma^{\nu(\bar{\nu})}}{dx dy} = \frac{G_F^2 M}{\pi(1 + Q^2/M_W^2)} \left[ \left( 1 - y - \frac{Mxy}{2E_\nu} + \frac{y^2}{2} \frac{1 + 4M^2 x^2/Q^2}{1 + R(x, Q^2)} \right) F_2^{\nu(\bar{\nu})} \pm \left( y - \frac{y^2}{2} \right) x F_3^{\nu(\bar{\nu})} \right]$$

Structure Functions in terms of parton distributions (for  $\nu$ -scattering)

$$F_2^{\nu(\bar{\nu})N} = \sum [xq^{\nu(\bar{\nu})N}(x) + x\bar{q}^{\nu(\bar{\nu})N}(x) + 2xk^{\nu(\bar{\nu})N}(x)]$$

$$xF_3^{\nu(\bar{\nu})N} = \sum [xq^{\nu(\bar{\nu})N}(x) - x\bar{q}^{\nu(\bar{\nu})N}(x)] = x(d_\nu(x) + u_\nu(x)) \pm 2x(s(x) - c(x))$$

$$R = \frac{\sigma_L}{\sigma_T}$$

# Neutrino Experiments have been studying QCD for about 40 years

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- ◆ For example, Gargamelle made one of the first measurements of  $\Lambda_{\text{ST}}$  in the early 1970's using sum rules and the  $x$ - $Q^2$  behavior of the structure functions  $F_2$  and  $xF_3$  measured off heavy liquids.
- ◆ BEBC followed with QCD studies using  $\nu + p$  and  $\nu + D$  scattering.

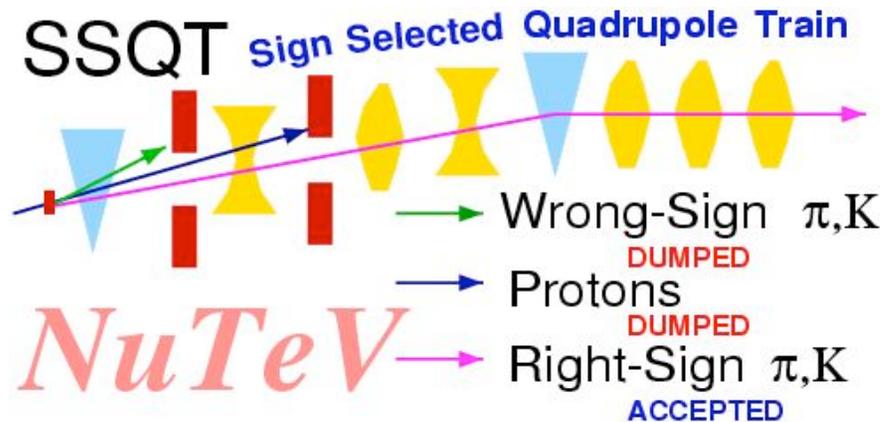
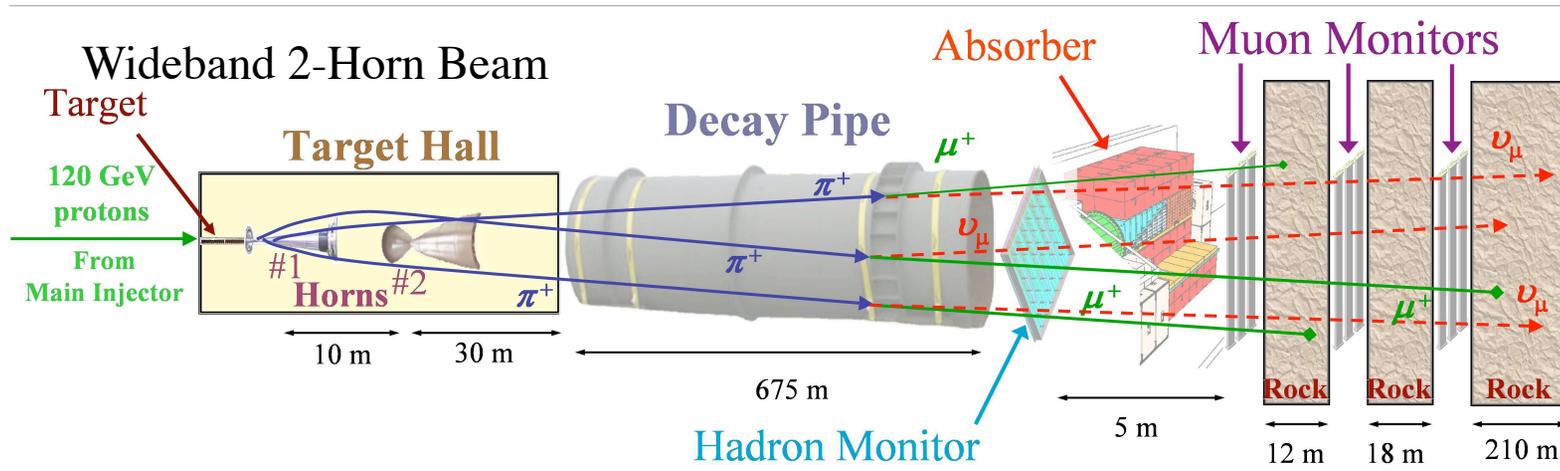
## Most “Recent” DIS Experiments

- ◆ There followed a long string of  $\nu$  scattering experiments with **increasing statistics** and **decreasing systematic errors** ....

	$E_\nu$ range ( $\langle E_\nu \rangle$ ) (GeV)	Run	Target A	$E_\mu$ scale	$E_{HAD}$ scale	Detector
<b>NuTeV (CCFR)</b>	<b>30-360(120)</b>	<b>96-97</b>	<b>Fe</b>	<b>0.7%</b>	<b>0.43%</b>	<b>Coarse</b>
<b>NOMAD</b>	<b>10-200(27)</b>	<b>95-98</b>	<b>Various (mainly C)</b>	<b>--</b>	<b>---</b>	<b>Fine- grained</b>
<b>CHORUS</b>	<b>10-200(27)</b>	<b>95-98</b>	<b>Pb</b>	<b>2%</b>	<b>5%</b>	<b>Fine- grained</b>
<b>MINOS</b>	<b>3-15</b>	<b>05-10</b>	<b>Fe</b>	<b>2.5%</b>	<b>5.6%</b>	<b>Coarse</b>

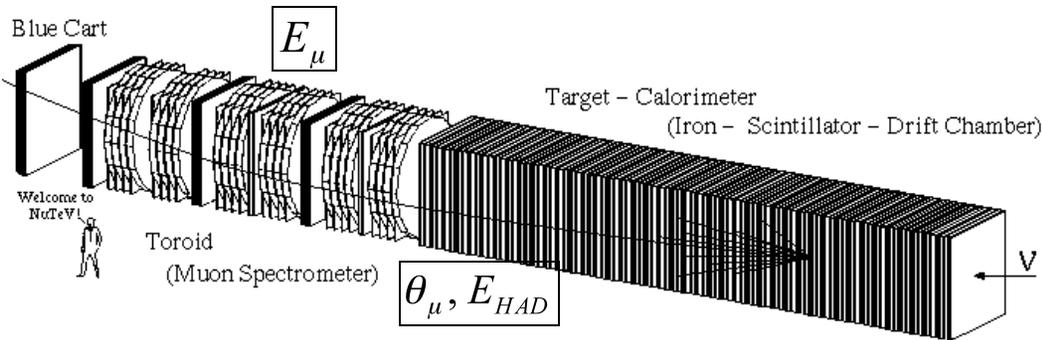
# Neutrino Beamlines

- ◆ Intense proton beam on a target. Collect  $\pi$  and K and steer into a decay area. Absorb hadrons and muons from beam leaving only neutrinos.

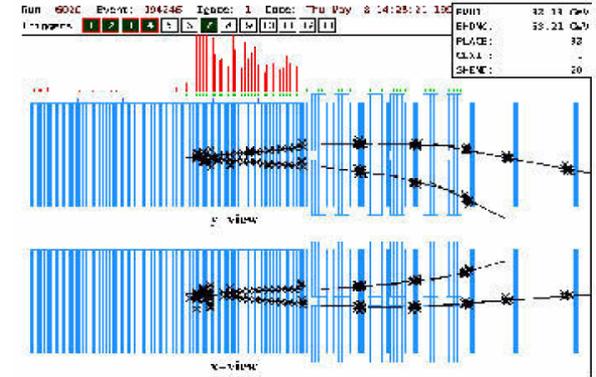


# The NuTeV Experiment: 800 GeV Protons

> 3 million neutrino/antineutrino events with  $20 \leq E_\nu \leq 400 \text{ GeV}$



Refurbished CCFR detector



Target Calorimeter:

- ◆ Steel-Scintillator Sandwich (10 cm)

$$\frac{\delta E}{E} \approx \frac{0.86}{\sqrt{E}} \text{ -resolution}$$

- ◆ Tracking chambers for muon track and vertex

◆ Muon Spectrometer:

Three toroidal iron magnets with five sets of drift chambers

$$\langle B_\varphi \rangle \approx 1.7T, p_t \approx 2.4 \text{ GeV} / c$$

$$\delta(1/p)/(1/p) \sim 11\% \text{ MCS dominated}$$

- ◆ Always focusing for **leading muon**

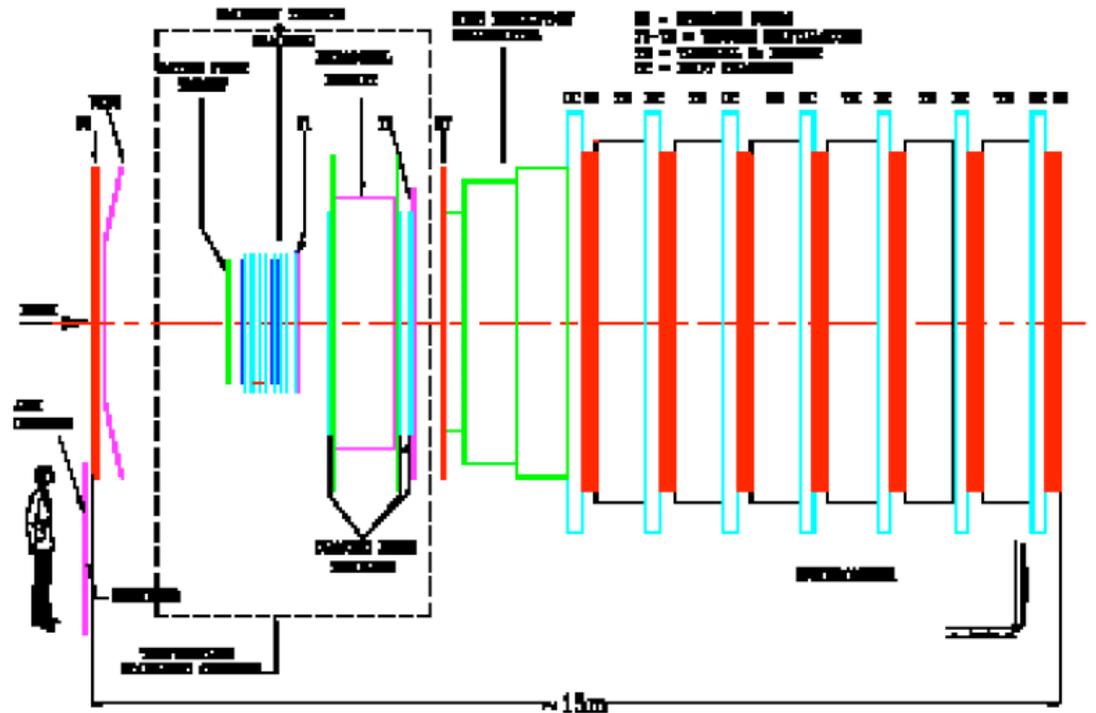
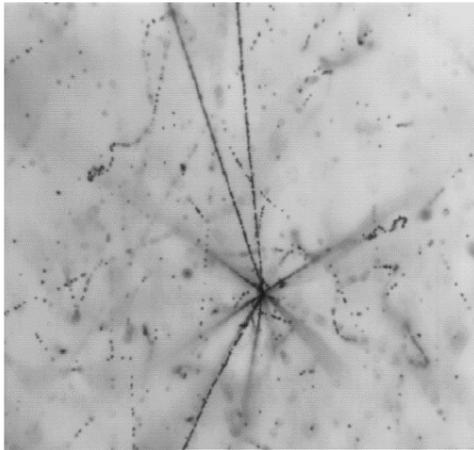
To confront leading systematic errors, there was a continuous calibration beam that yielded

**Hadron energy scale**  $\frac{\Delta E_{HAD}}{E_{HAD}} = 0.43\%$

**Muon energy scale**  $\frac{\Delta E_\mu}{E_\mu} = 0.7\%$

# CHORUS Experiment – nuclear emulsions

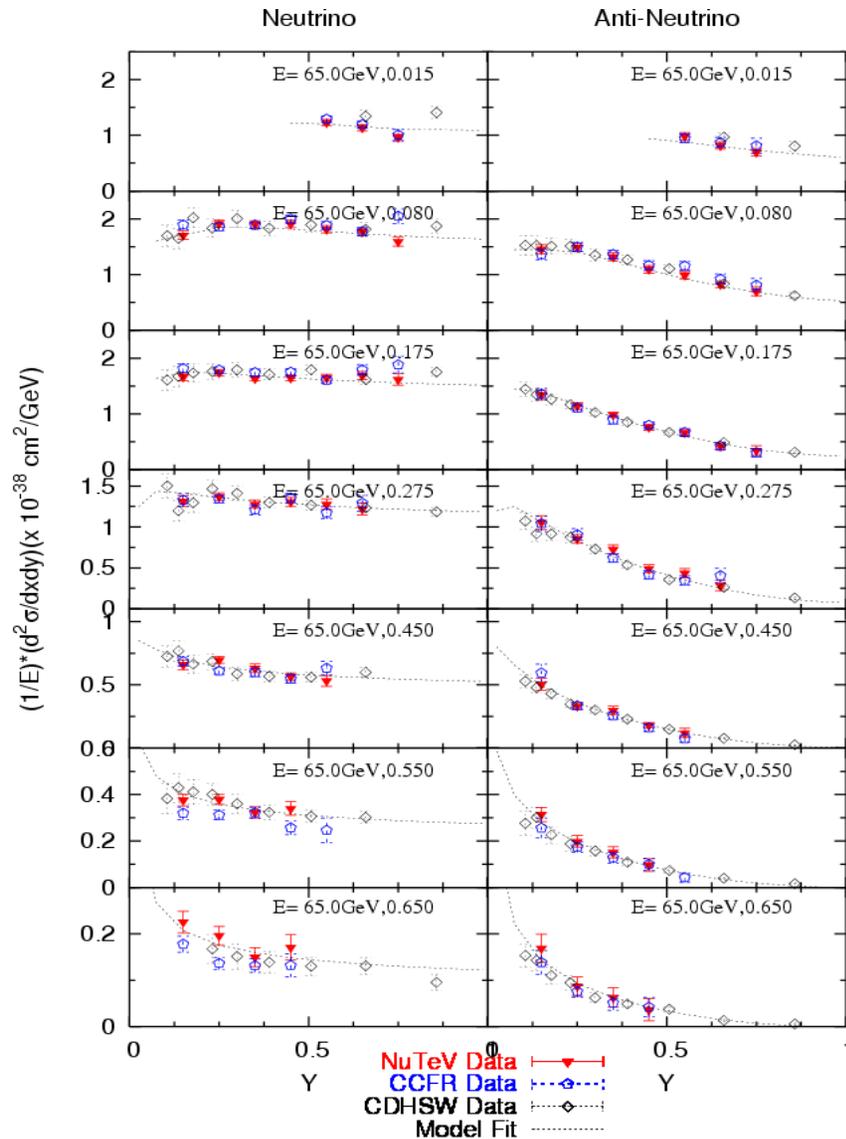
- ◆ 450 GeV protons  $\rightarrow$  10 – 200 GeV  $\nu$ , 6% wrong-sign background
- ◆ Nuclear Emulsion Target (Pb, Fe, Ca and C)
- ◆ Scintillating Fiber tracker



Muon energy scale – 2.5%  
Hadron Energy Scale - 5%  
(test beam exposure)

# NuTeV CC Differential Cross Section

## $d\sigma/dy$ for different $E_\nu$

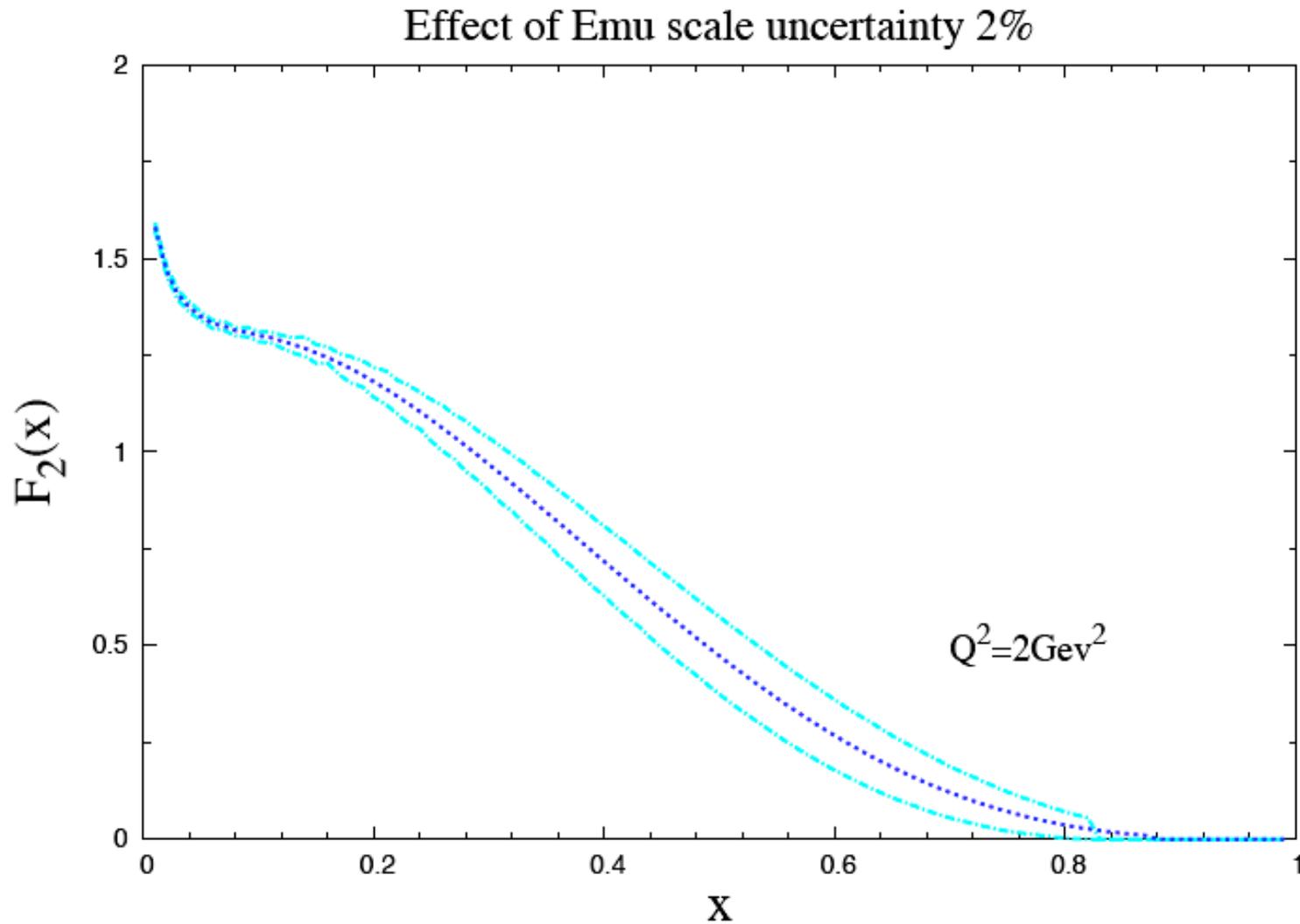


	$E_\mu$ scale	$E_{\text{HAD}}$ scale	$E_\nu$ range (GeV)
CDHSW	2%	2.5%	20-200
CCFR	1%	1%	30-360
NuTeV	0.7%	0.43%	30-360

- ◆ NuTeV has increased statistics compared to other  $\nu$ -Fe experiments.
- ◆ Significant reduction in the largest systematic uncertainties : -  $E_\mu$  and  $E_{\text{HAD}}$  scales

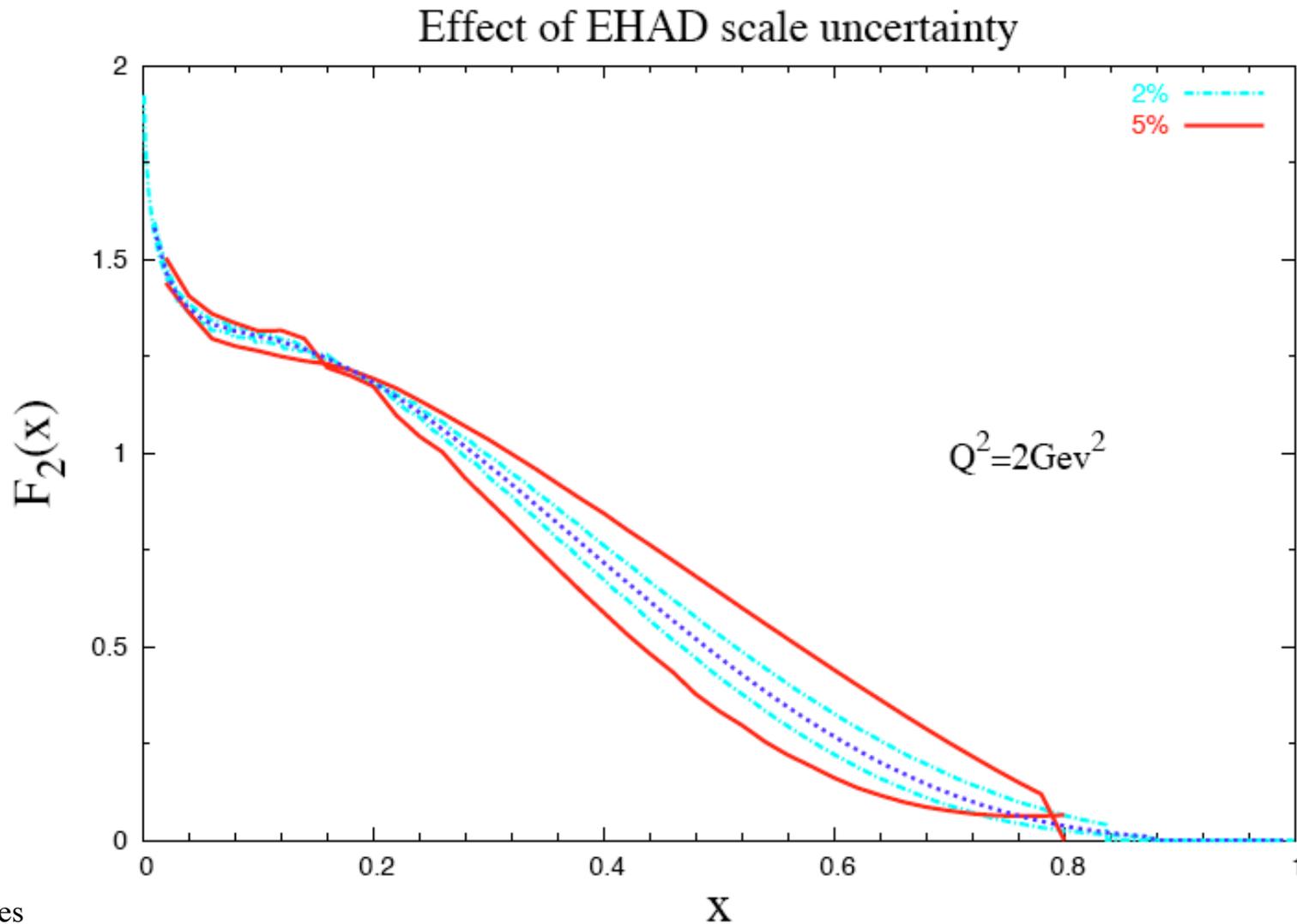
# Estimated systematic error: $E_\mu$ scale

NuTeV achieved 0.7%



# Estimated systematic error: $E_{\text{had}}$ scale

NuTeV achieved 0.43%



# $F_2$ and $xF_3$ Measurement

## $F_2$

$$\left[ \frac{d^2\sigma^v}{dx dy} + \frac{d^2\sigma^{\bar{v}}}{dx dy} \right] \frac{\pi}{G_F^2 ME} =$$

$$= 2 \bar{F}_2 \left( 1 - y - \frac{Mxy}{2E} + \frac{y^2}{2} \frac{1 + 4M^2 x^2 / Q^2}{1 + R} \right) + y \left( 1 - \frac{y}{2} \right) \Delta xF_3$$

- ◆ Perform 1-parameter fit for  $F_2$
- ◆  $\Delta xF_3$  model
- ◆  $R_L$  model

## $xF_3$

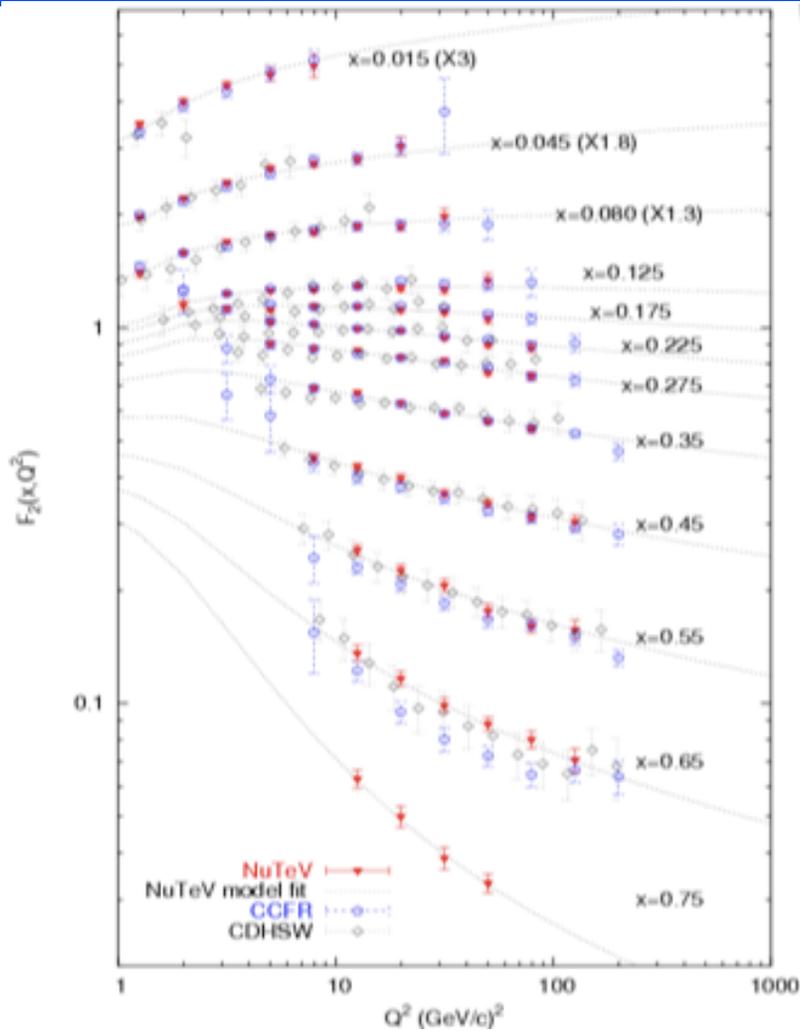
$$\left[ \frac{d^2\sigma^v}{dx dy} - \frac{d^2\sigma^{\bar{v}}}{dx dy} \right] \frac{\pi}{G_F^2 ME} =$$

$$= \Delta F_2 \left( 1 - y - \frac{Mxy}{2E} + \frac{y^2}{2} \frac{1 + 4M^2 x^2 / Q^2}{1 + R} \right) + 2 y \left( 1 - \frac{y}{2} \right) x\bar{F}_3$$

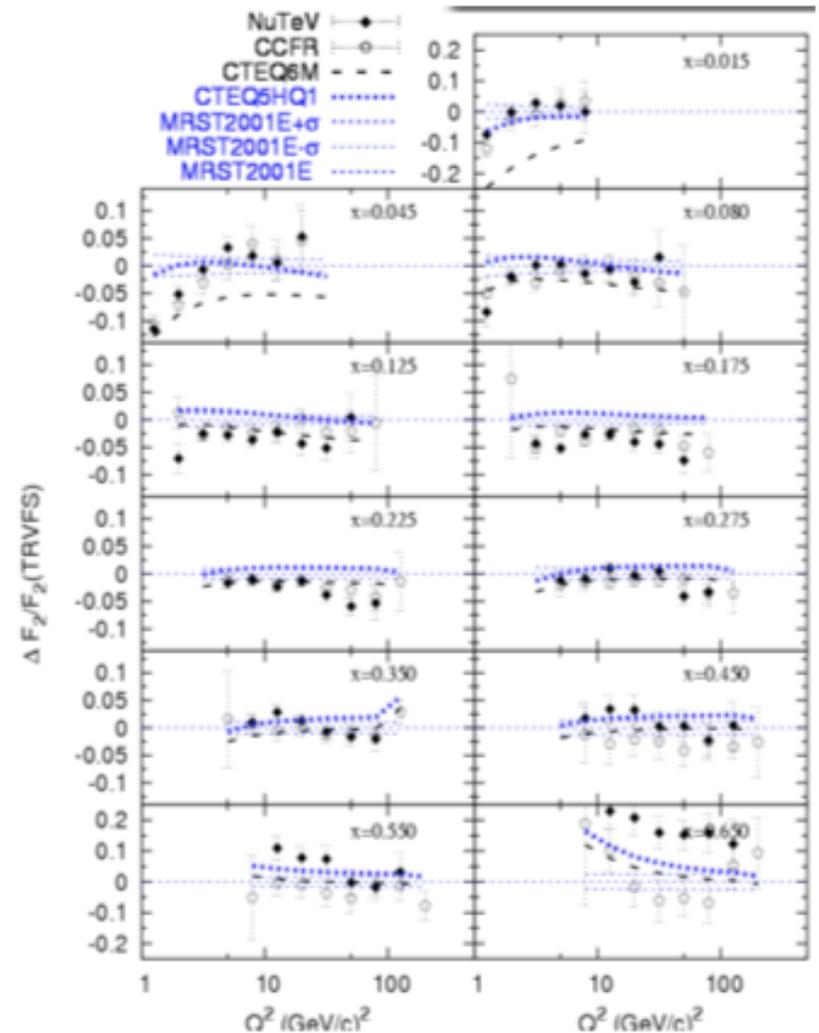
- ◆ Perform 1-parameter fit for  $xF_3$
- ◆  $\Delta F_2$  is very small and is neglected

- ◆ Radiative corrections applied
- ◆ Isoscalar correction applied

# NuTeV $F_2$ Measurement

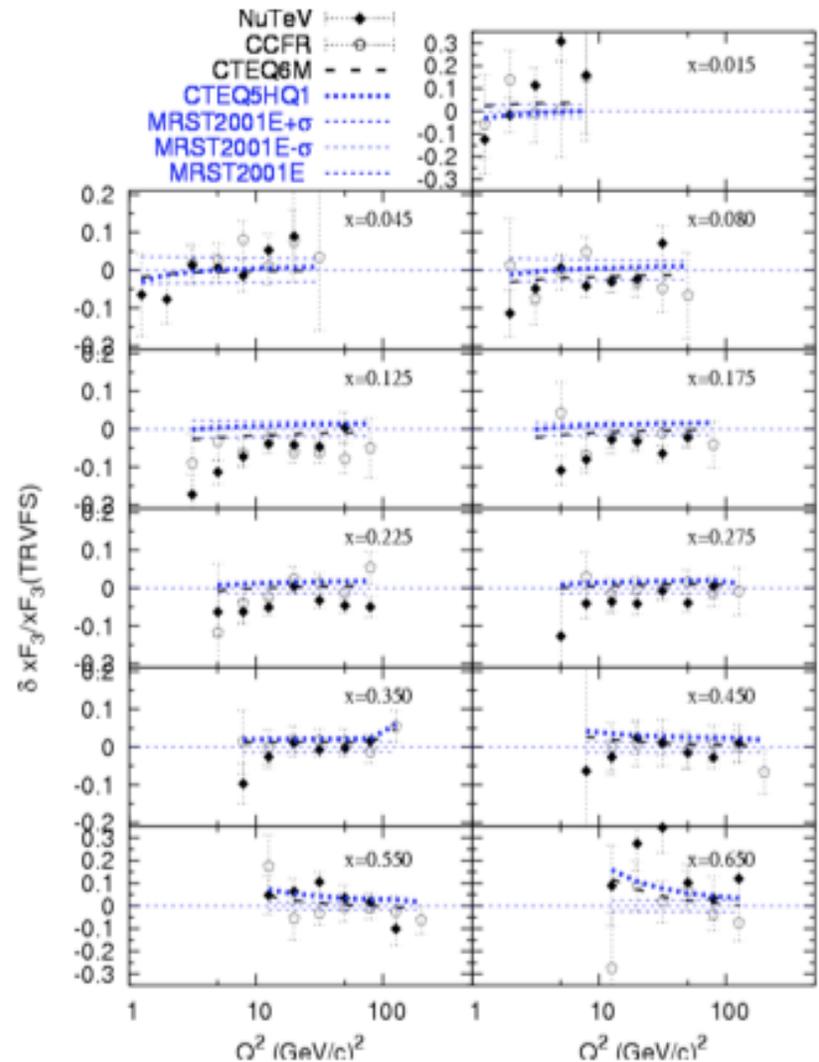
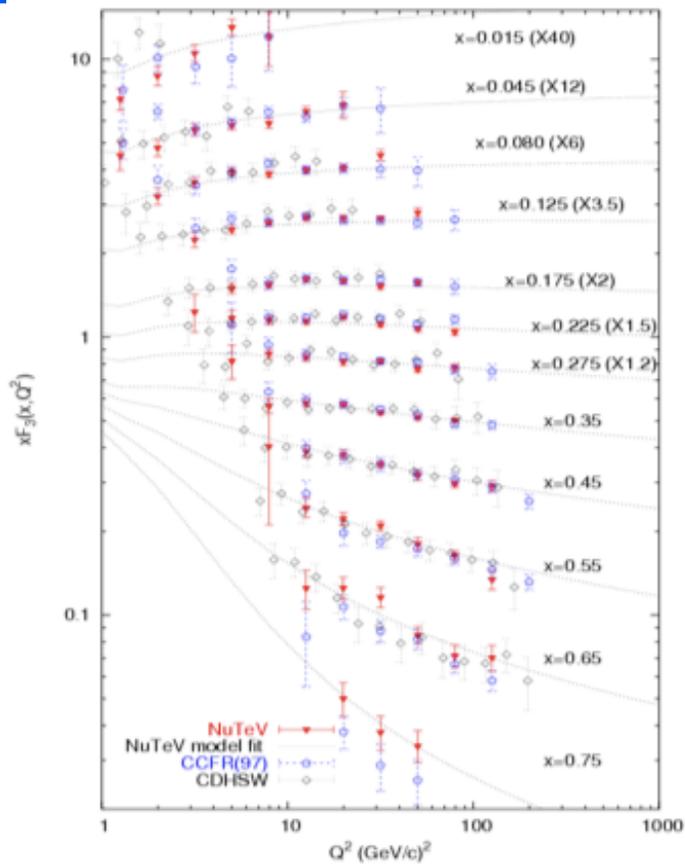


◆ Comparison of NuTeV  $F_2$  with global fits



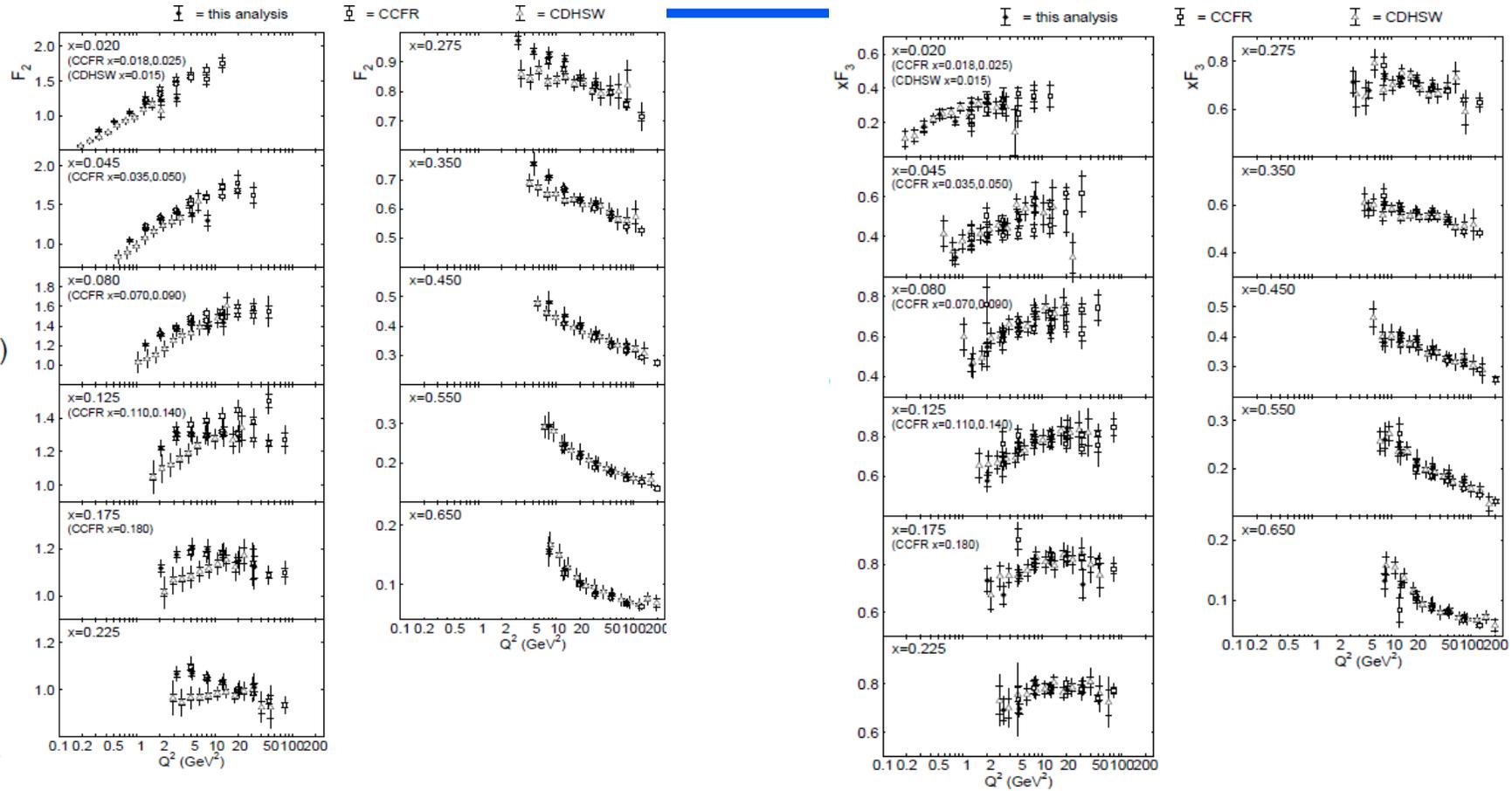
◆ At  $x > 0.4$  NuTeV is systematically above CCFR

# NuTeV $x F_3$ Measurement



- ◆ At  $x > 0.5$  NuTeV is systematically above CCFR
- ◆ NuTeV  $F_2$  agrees with theory for medium  $x$ .
- ◆ At low  $x$  different  $Q^2$  dependence.
- ◆ At high  $x$  ( $x > 0.5$ ) NuTeV is systematically higher.

# CHORUS Structure Functions: $\nu$ Pb



- ◆ First  $\nu$ -Pb differential cross section and structure functions
- ◆ CHORUS measurement favors CCFR over NuTeV
- ◆ Much larger systematic errors than the NuTeV experiment

# Parton Distribution Functions: What Can We Learn With All Six Structure Functions?

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**Recall Neutrinos have the ability to directly resolve flavor of the nucleon's constituents:  
 $\nu$  interacts with  $d$ ,  $s$ ,  $\bar{u}$ , and  $\bar{c}$  while  $\bar{\nu}$  interacts with  $u$ ,  $c$ ,  $d$  and  $s$ .**

**Using Leading order expressions:**

$$F_2^{\bar{\nu}N}(x, Q^2) = x[u + \bar{u} + d + \bar{d} + 2s + 2c]$$

$$F_2^{\nu N}(x, Q^2) = x[u + \bar{u} + d + \bar{d} + 2s + 2\bar{c}]$$

$$xF_3^{\bar{\nu}N}(x, Q^2) = x[u + d - \bar{u} - \bar{d} - 2s + 2c]$$

$$xF_3^{\nu N}(x, Q^2) = x[u + d - \bar{u} - \bar{d} + 2s - 2\bar{c}]$$

**Taking combinations of the Structure functions**

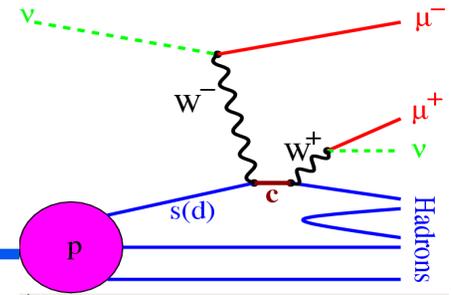
$$F_2^{\nu} - xF_3^{\nu} = 2(\bar{u} + \bar{d} + 2\bar{c})$$

$$F_2^{\bar{\nu}} - xF_3^{\bar{\nu}} = 2(\bar{u} + \bar{d} + 2\bar{s})$$

$$xF_3^{\nu} - xF_3^{\bar{\nu}} = 2[(s + \bar{s}) - (\bar{c} + c)]$$

# Charm Production by Neutrinos

## a direct look at strange sea.



- ◆ Charm quark is produced from CC neutrino interaction with s(d) quark in the nucleon. d-quark interaction is CKM suppressed
- ◆ Detect charm via the semi-leptonic decay which yields a very clear signature – two opposite sign muons
- ◆ It is sensitive to  $m_c$  through  $E_\nu$  dependence.
- ◆ With high-purity  $\nu$  and  $\bar{\nu}$  beams, NuTeV made high statistics separate s and  $\bar{s}$  measurements: 5163  $\nu$  and 1380  $\bar{\nu}$
- ◆ Could then make a measurement of s –  $\bar{s}$ .

# Strange Sea Asymmetry

$$s^- = (s - \bar{s})$$

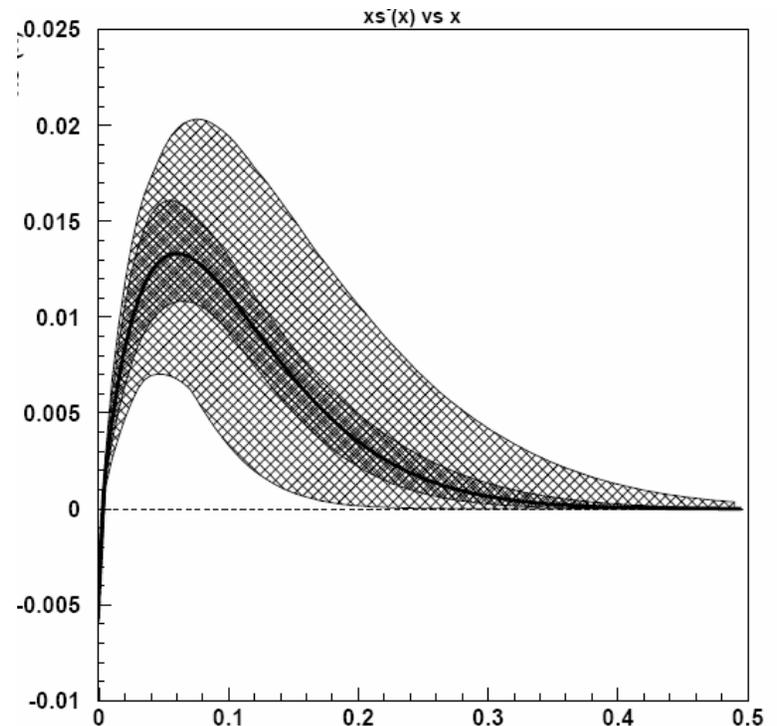
$$S^- = \int_0^1 x s^-(x) dx = 0.00196 \pm 0.00046 \text{ (stat)} \pm 0.00045 \text{ (syst)} \pm 0.00128 \text{ (external)}$$

- ◆ CTEQ inspired NLO model,
- ◆ in the fit net strangeness of the nucleon is forced to 0.

$$m_c = 1.41 \pm 0.10 \text{ (stat)} \pm 0.008 \text{ (syst)} \pm 0.12 \text{ (ext)} \text{ GeV}/c^2$$

**This is an analysis of strange quarks in an Fe nucleus!**

**Are  $\nu$  nuclear effects known? Are they the same for  $\nu$  and  $\bar{\nu}$ ?**



## Summary $\nu$ Scattering Results – NuTeV

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NuTeV accumulated over 3 million neutrino / antineutrino events with  $20 \leq E_\nu \leq 400$  GeV.

NuTeV considered 23 systematic uncertainties.

NuTeV  $\sigma$  agrees with other  $\nu$  experiments and theory for medium  $x$ .

*At low  $x$  different  $Q^2$  dependence.*

*At high  $x$  ( $x > 0.6$ ) NuTeV is systematically higher.*

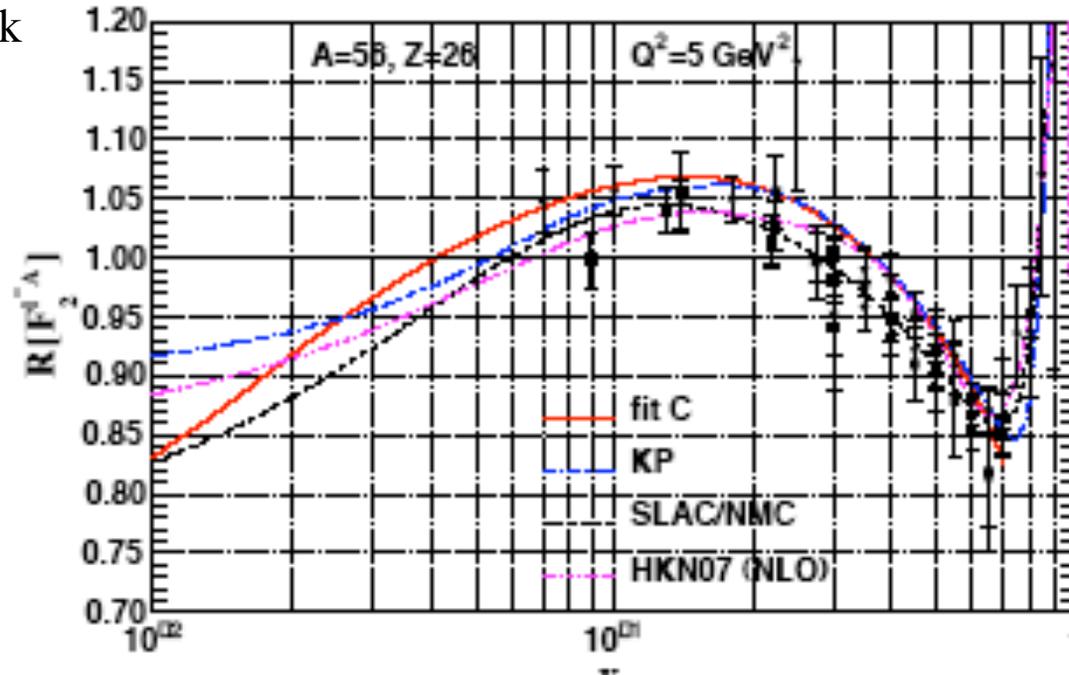
NuTeV extracts the **strange quark** distribution via charm production using both  $\nu$  and  $\bar{\nu}$  and gets a value of  $S^-(x)$

**All of the NuTeV Results are for  $\nu - \text{Fe}$  interactions and where necessary have assumed the nuclear corrections for neutrino interactions are the same as  $\text{I}^\pm$ . Is this really the case?**

# Nuclear Structure Function Corrections

## $Q^\pm$ (Fe/D<sub>2</sub>)

See yesterday's talk  
by Solvignon !



- ◆  $F_2$  / nucleon changes as a function of  $A$ . Measured in  $\mu/e - A$ , not in  $\nu - A$
- ◆ Good reason to consider nuclear effects are DIFFERENT in  $\nu - A$ .
  - ▼ Presence of axial-vector current.
  - ▼ Different nuclear effects for valance and sea --> different shadowing for  $xF_3$  compared to  $F_2$ .

# CTEQ study: The Impact of new neutrino DIS and Drell-Yan data on large-x parton distributions

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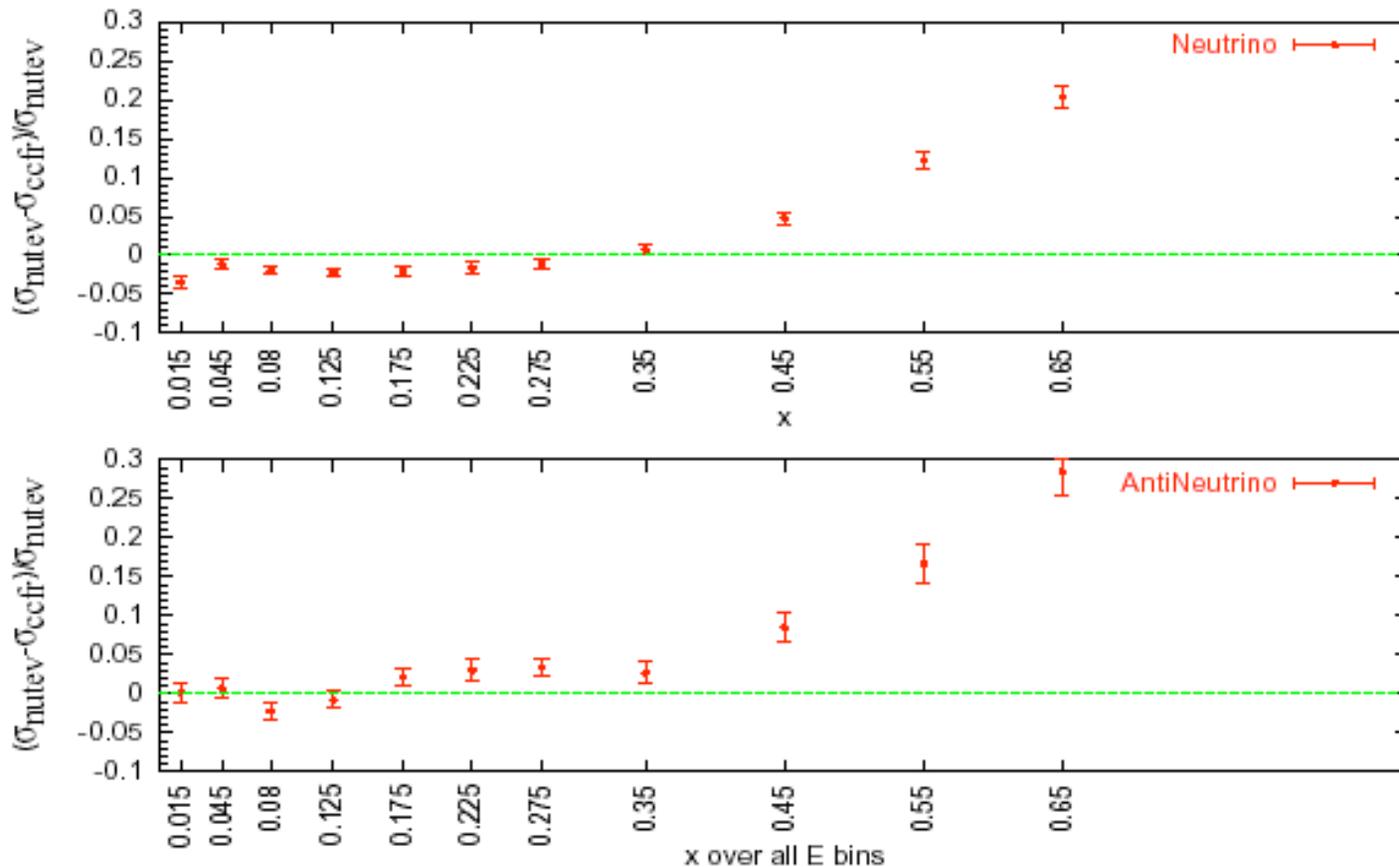
Joey Huston - MSU, Cynthia Keppel - Hampton, Steve Kuhlmann - ANL,  
JGM - Fermilab, Fred Olness - SMU, Jeff Owens - Florida State,  
Jon Pumplin and Dan Stump - MSU

Published in **Phys.Rev.D75:054030,2007.**  
e-Print: **hep-ph/0702159**

Had to use  $l^\pm$ -Fe correction factors to combine NuTeV  $\nu$ -Fe results with  
E866 p-H and p-D Drell-Yan results.  
Tension between NuTeV and E866 started us on a rather convoluted path  
to extracting **nuclear effects from neutrino interactions.**

NuTeV ( $\nu$ -Fe) Compared to CCFR (in PDF fits).  
 At High- $x$  NuTeV Indicates Effect **Opposite** to E866 D-Y.  
 (CHORUS ( $\nu$ -Pb) in between CCFR and NuTeV at high  $x$ )

Is the tension between NuTeV and E866 coming from applying  $I^{\pm}$ -Fe nuclear corrections to the NuTeV  $\nu$ -Fe measurements?



# CTEQ High-x Study: nuclear effects

No high-statistics D2 data – “make it” from PDFs

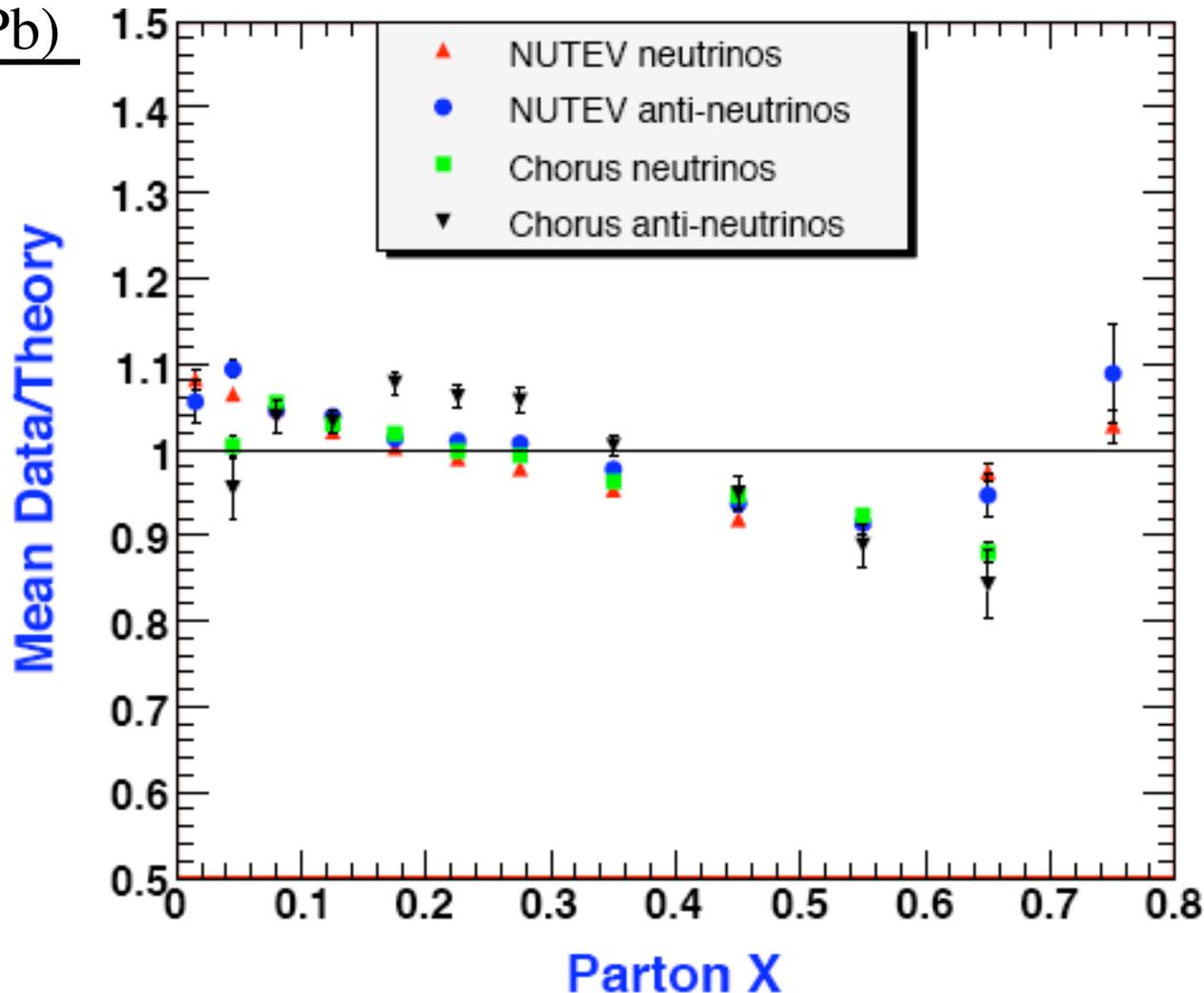
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- ◆ Form reference fit mainly nucleon (as opposed to nuclear) scattering results:
  - ▼ BCDMS results for  $F_2^p$  and  $F_2^d$
  - ▼ NMC results for  $F_2^p$  and  $F_2^d/F_2^p$
  - ▼ H1 and ZEUS results for  $F_2^p$
  - ▼ CDF and DØ result for inclusive jet production
  - ▼ CDF results for the W lepton asymmetry
  - ▼ E-866 results for the ratio of lepton pair cross sections for pd and pp interactions
  - ▼ E-605 results for dimuon production in pN interactions.
  
- ◆ Correct for deuteron nuclear effects

# NuTeV(Fe) and CHORUS (Pb) $\nu$ scattering (unshifted) $\sigma$ results compared to reference fit

**no nuclear corrections**

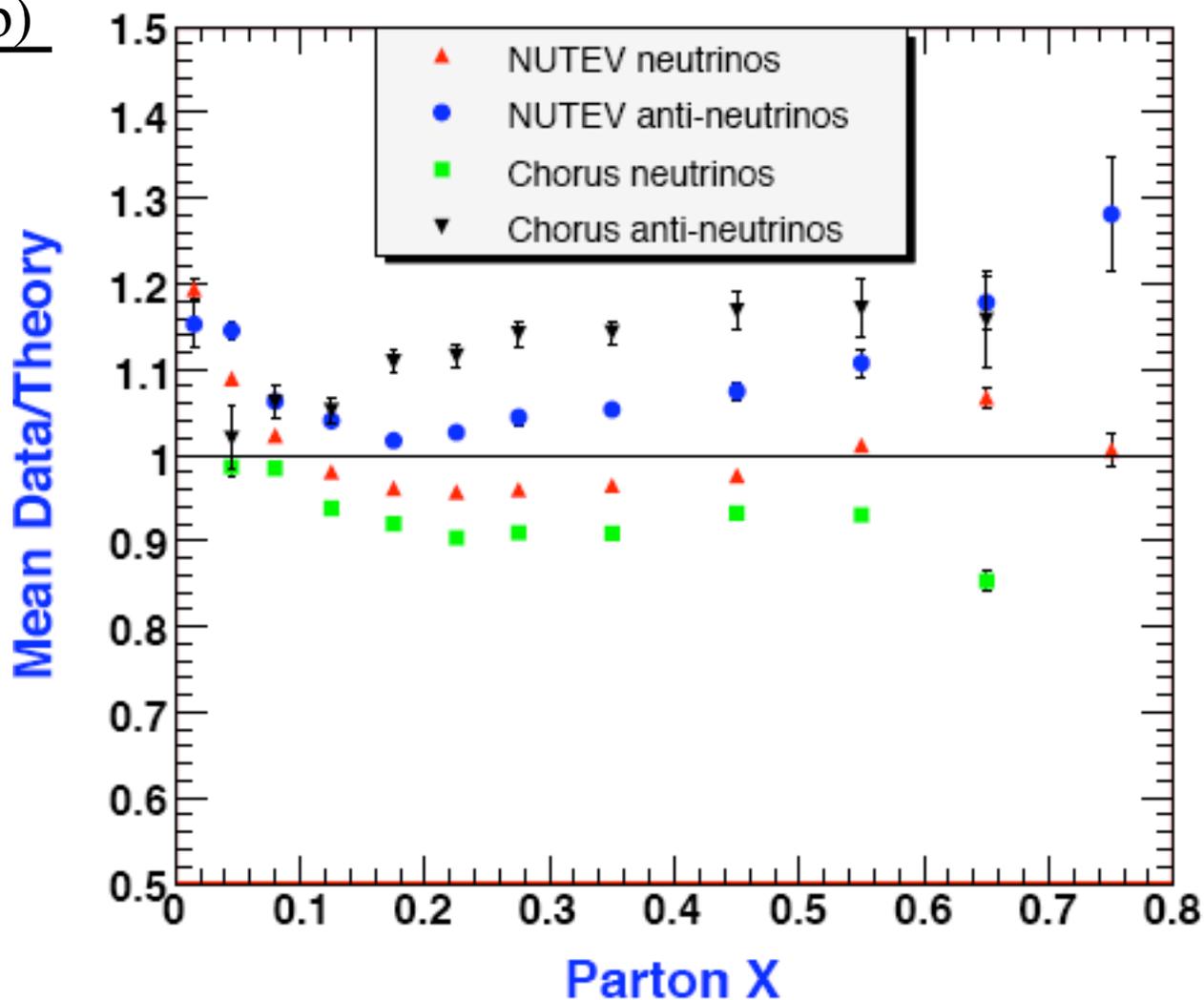
$$\frac{\sigma(\nu\text{Fe or } \nu\text{Pb})}{\sigma(\nu\text{D}_2)}$$



# NuTeV $\sigma(\text{Fe})$ & CHORUS $\sigma(\text{Pb})$ $\nu$ scattering (un-shifted) results compared to reference fit

**Kulagin-Petti nuclear corrections**

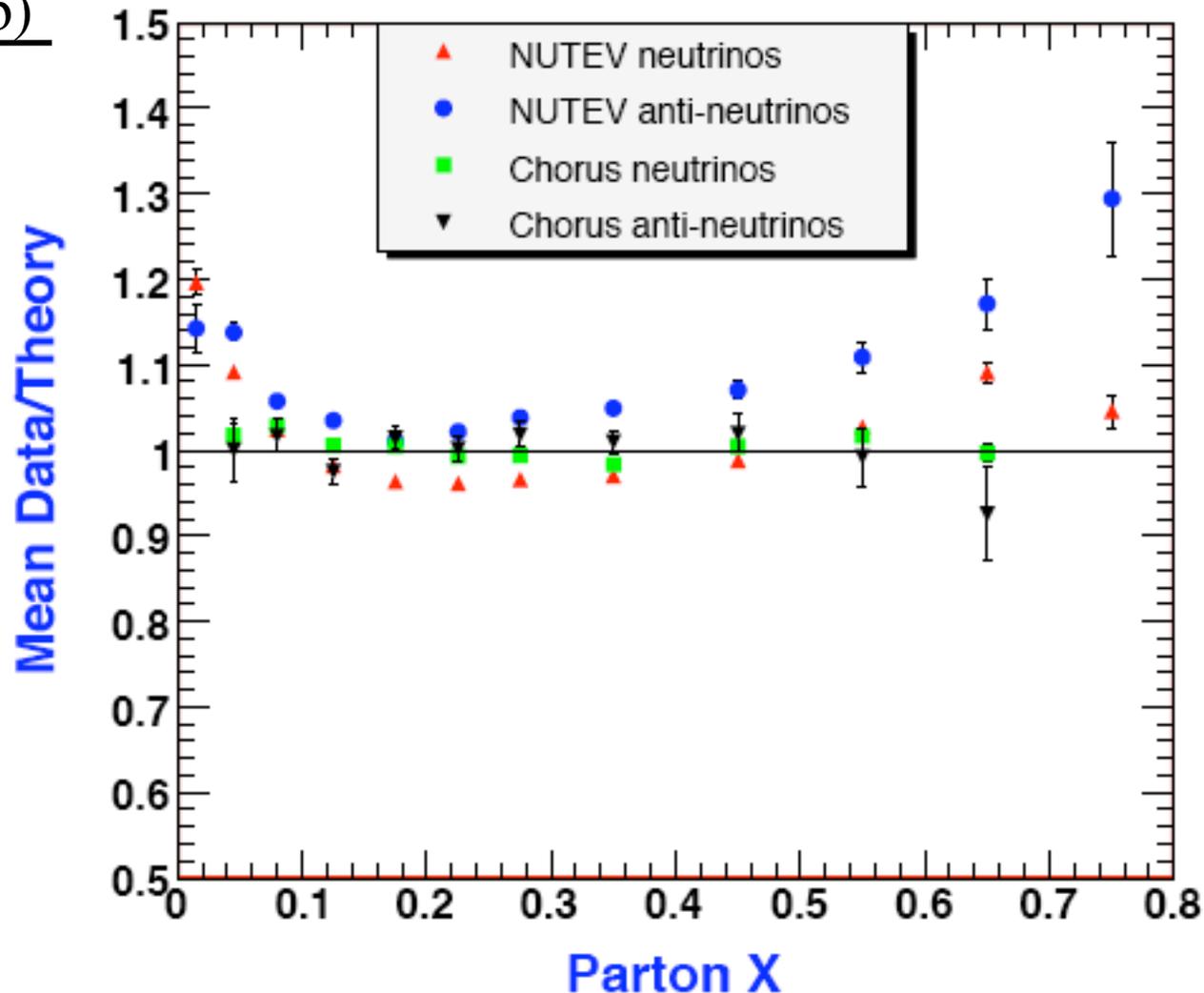
$$\frac{\sigma(\text{Fe or Pb})}{\sigma(\text{D}_2)}$$



# NuTeV $\sigma(\text{Fe})$ & CHORUS $\sigma(\text{Pb})$ $\nu$ scattering (shifted) results compared to reference fit

**Kulagin-Petti nuclear corrections**

$$\frac{\sigma(\text{Fe or Pb})}{\sigma(\text{D}_2)}$$



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# **Nuclear PDFs from neutrino deep inelastic scattering**

**I. Schienbein (SMU & LPSC-Grenoble, J-Y. Yu (SMU)  
C. Keppel (Hampton & JeffersonLab) J.G.M. (Fermilab),  
F. Olness (SMU), J.F. Olness (Florida State U)**

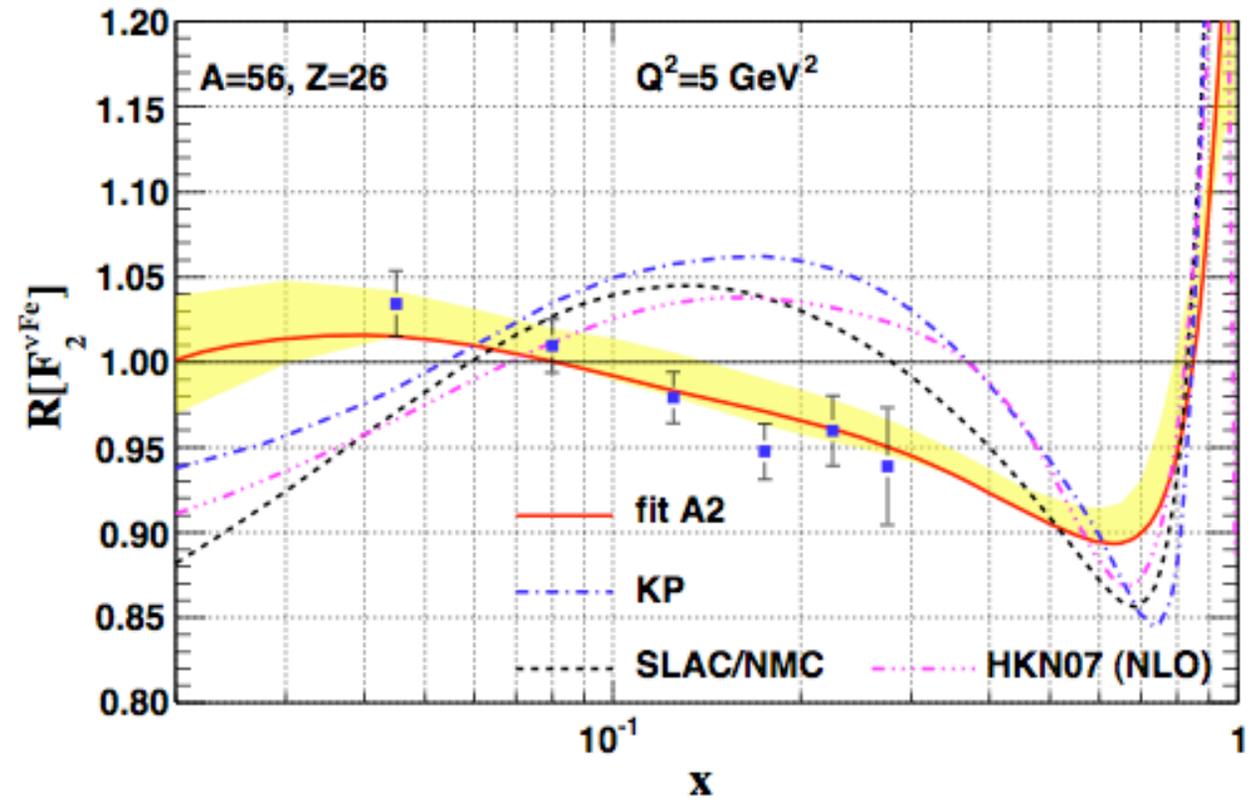
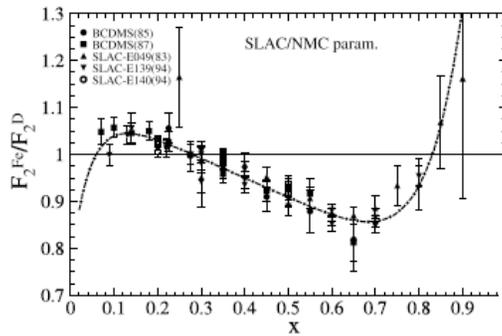
**e-Print: [arXiv:0710.4897](https://arxiv.org/abs/0710.4897) [hep-ph]**

# Same Reference Fit as Earlier Analysis

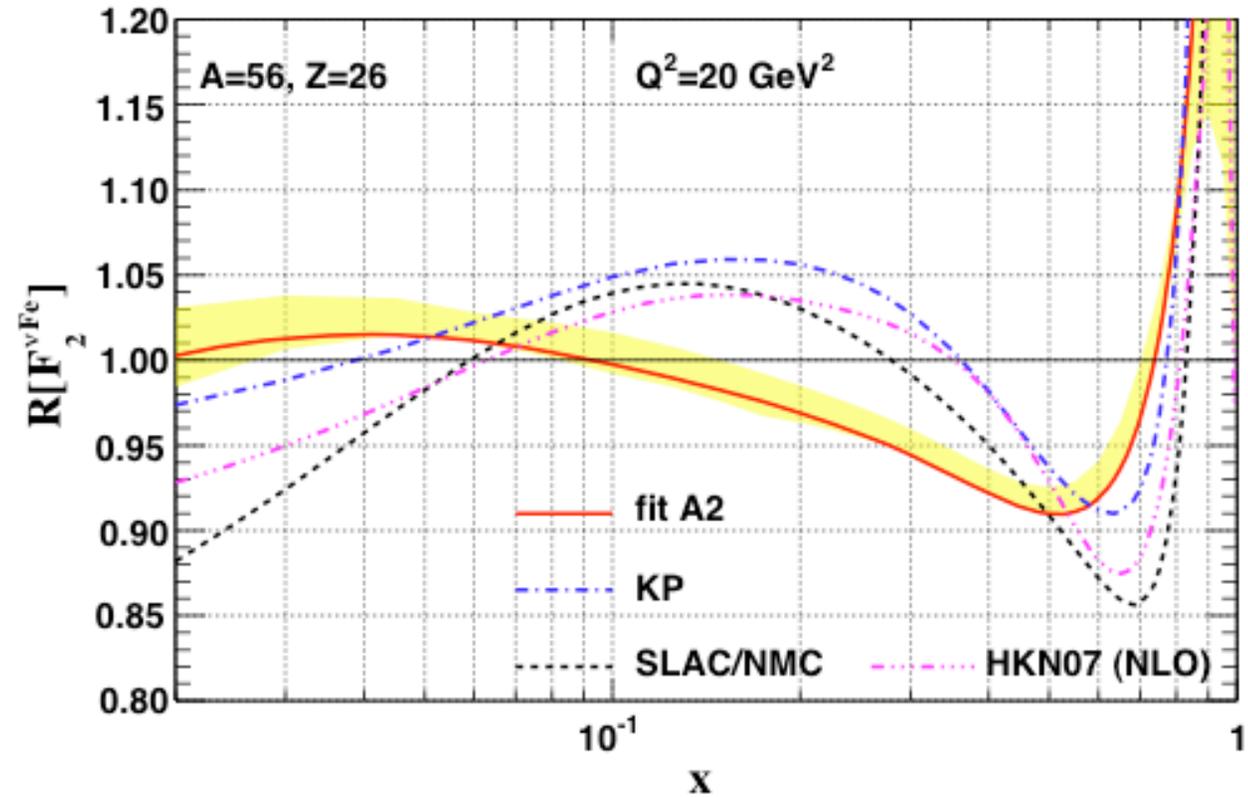
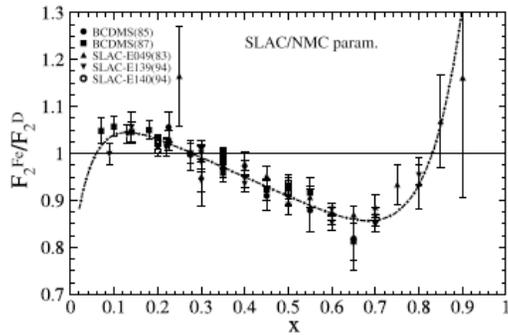
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- ◆ Form reference fit mainly nucleon (as opposed to nuclear) scattering results:
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- ◆ Correct for deuteron nuclear effects

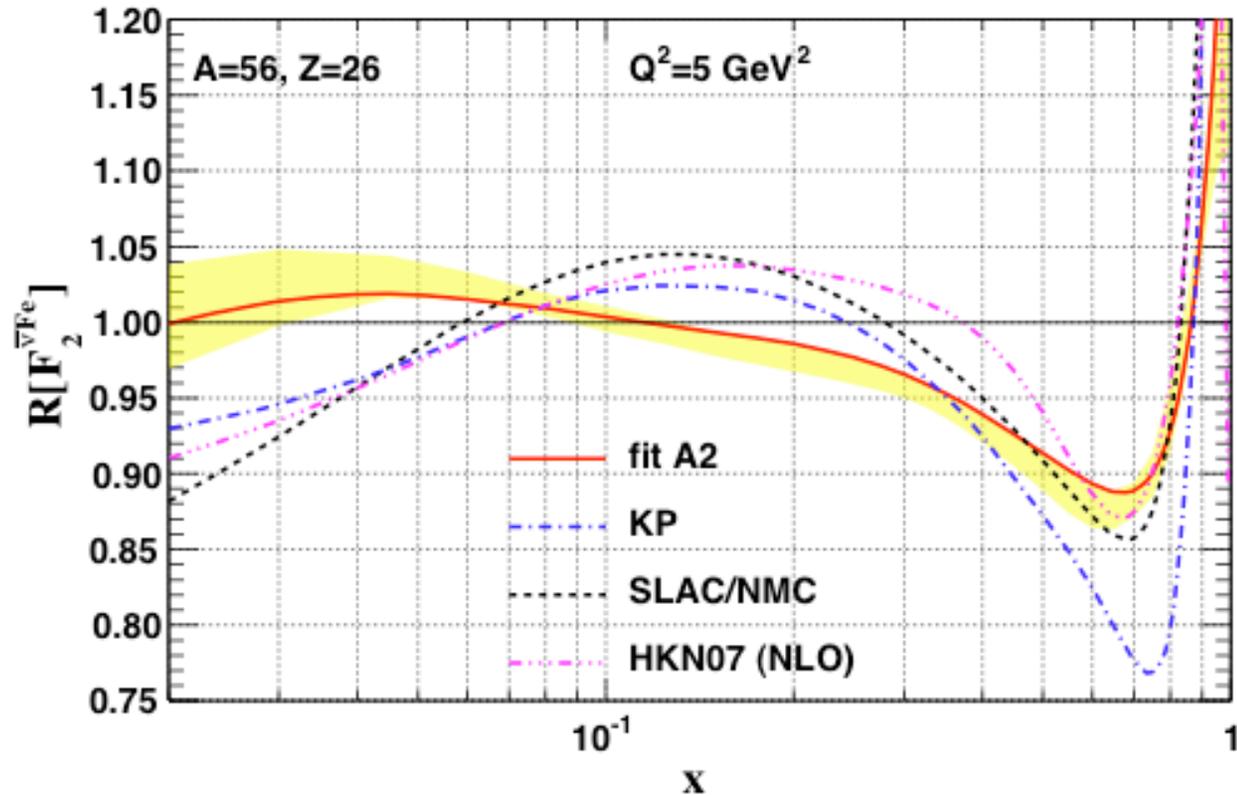
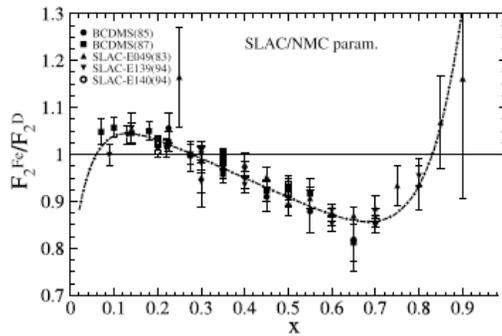
# $F_2$ Structure Function Ratios: $\nu$ -Iron



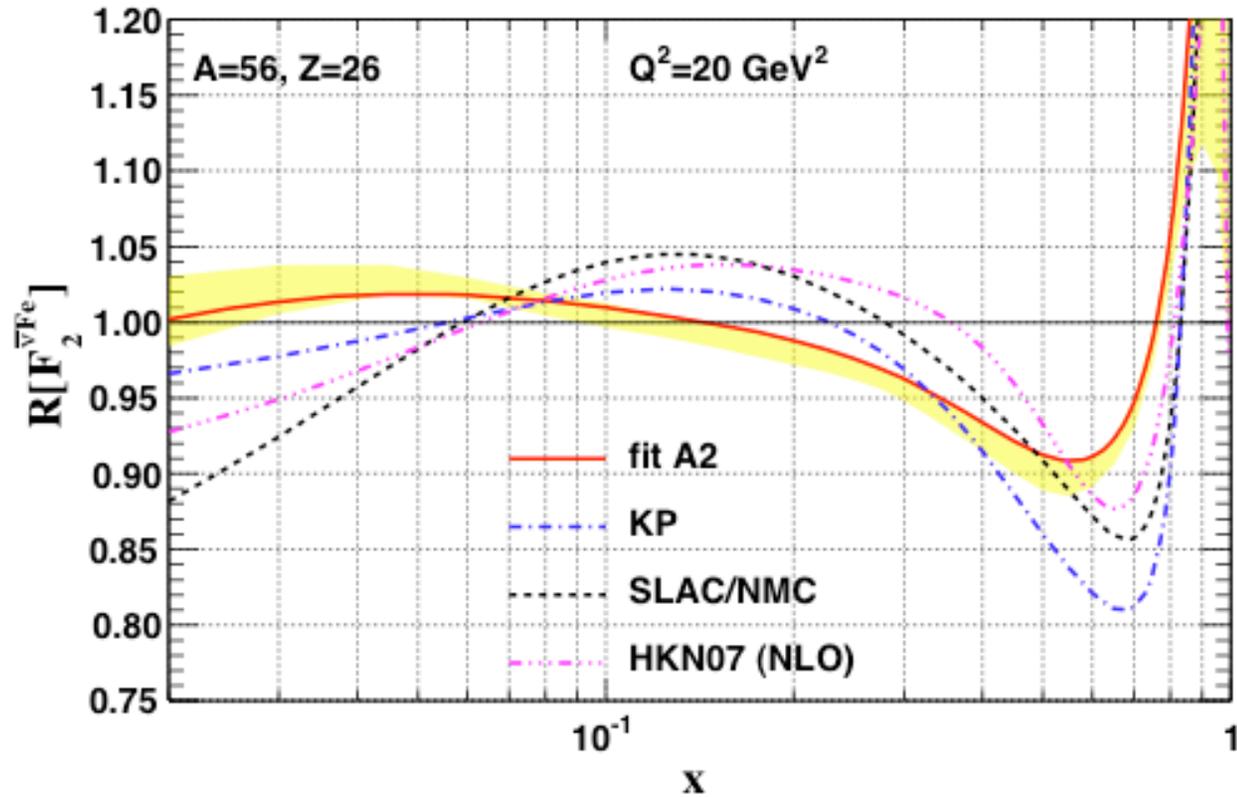
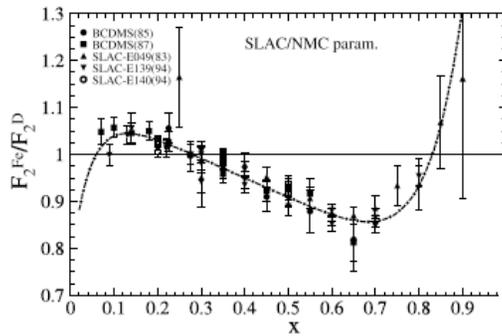
# $F_2$ Structure Function Ratios: $\nu$ -Iron



# $F_2$ Structure Function Ratios: $\bar{\nu}$ -Iron



# $F_2$ Structure Function Ratios: $\bar{\nu}$ -Iron



# First Conclusions

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- ◆ All high-statistics neutrino data is off nuclear targets. Need nuclear correction factors to include data off nuclei in fits with nucleon data.
- ◆ Nuclear correction factors ( $R$ ) different for neutrino-Fe scattering compared to charged lepton-Fe.
- ◆ Results from one experiment on one nuclear target... careful.
- ◆ If we combine  $\nu$ -nucleus with charged  $l^\pm$ -nucleus results and D-Y in a single global fit can we find a common description acceptable to both?

# Combined Analysis of $\nu A$ , $\ell A$ and DY data

Work in progress: **Kovarik**, Yu, Keppel, Morfin, Olness, Owens, **Schienbein**, Stavreva

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- ◆ Take an earlier analysis of  $\ell^\pm A$  data sets (built in A-dependence)
  - ▼ Schienbein, Yu, Kovarik, Keppel, Morfin, Olness, Owens,
  - ▼ PRD80 (2009) 094004
- ◆ For  $\ell^\pm A$  take  $F_2(A) / F_2(D)$  and  $F_2(A) / F_2(A')$  and DY  $\sigma(pA) / \sigma(pA')$ 
  - ▼ 708 Data points with  $Q > 2$  and  $W > 3.5$
- ◆ Use **8 Neutrino data sets**
  - ▼ NuTeV cross section data:  $\nu Fe$ ,  $\bar{\nu} Fe$
  - ▼ NuTeV dimuon off Fe data
  - ▼ CHORUS cross section data:  $\nu Pb$ ,  $\bar{\nu} Pb$
  - ▼ CCFR dimuon off Fe data
- ◆ Initial problem, with standard CTEQ cuts of  $Q > 2$  and  $W > 3.5$  neutrino data points (3134) far outnumber  $\ell^\pm A$  (708).

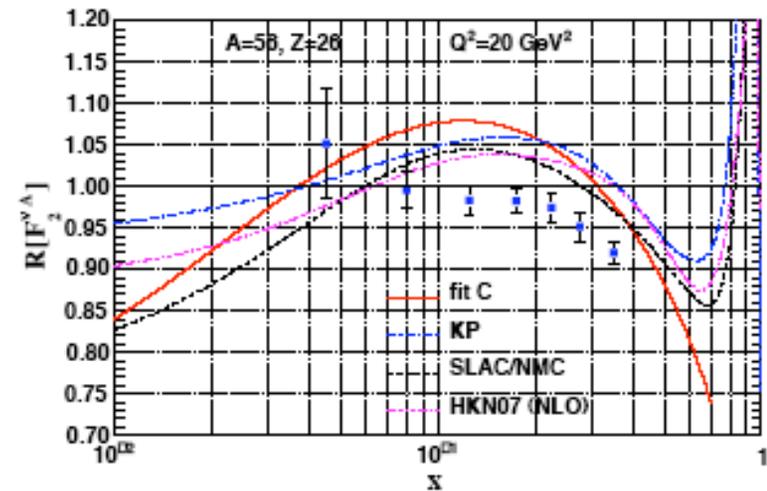
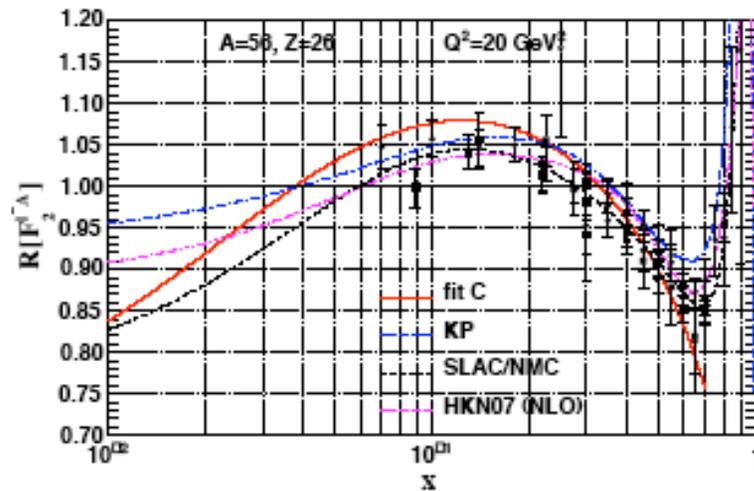
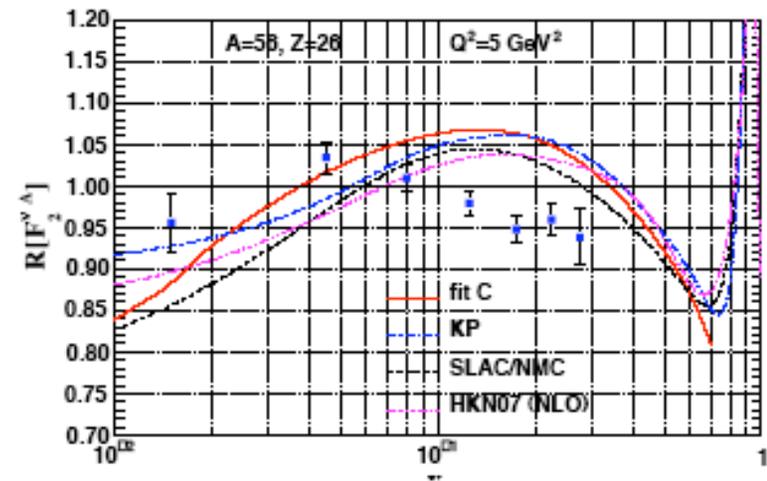
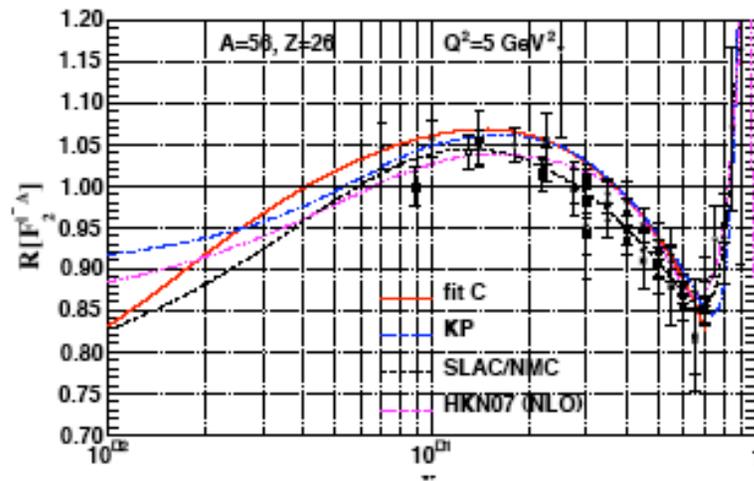
Use the usual procedure of observing the behavior of the fits as you adjust the

“weight” of the dominant data sample

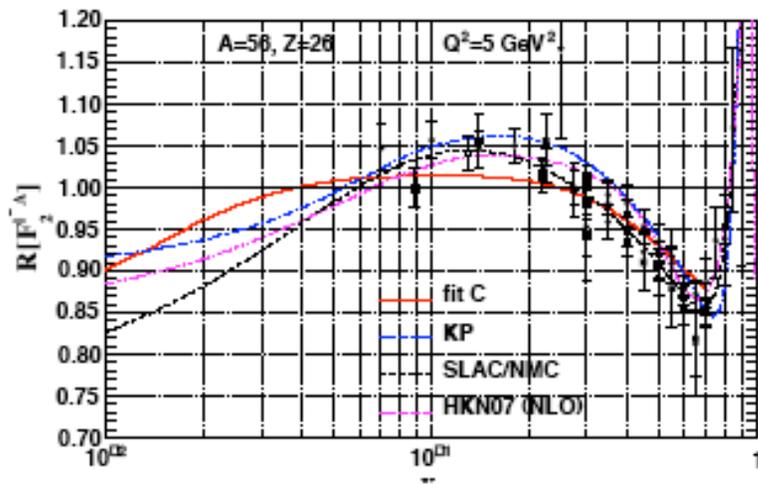
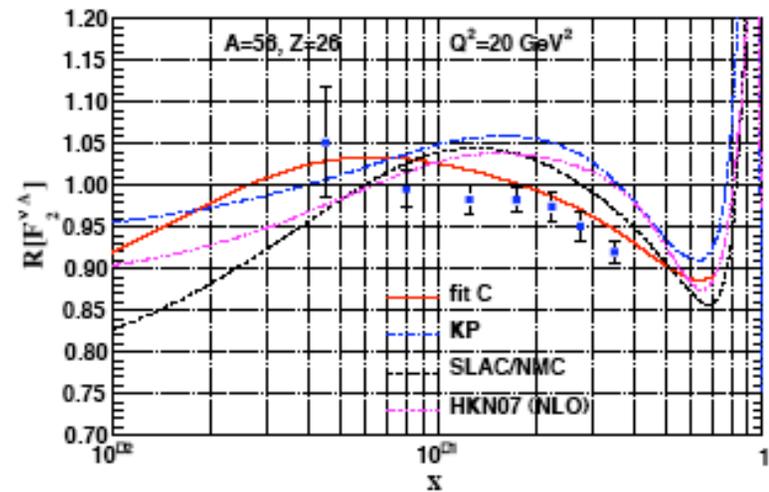
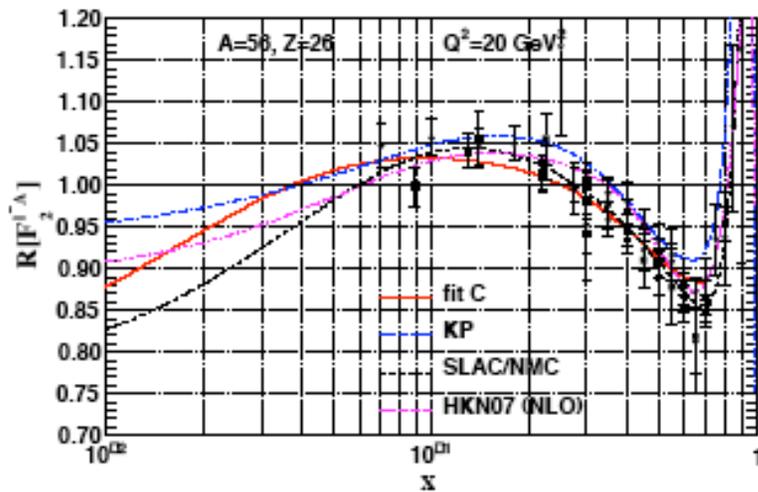
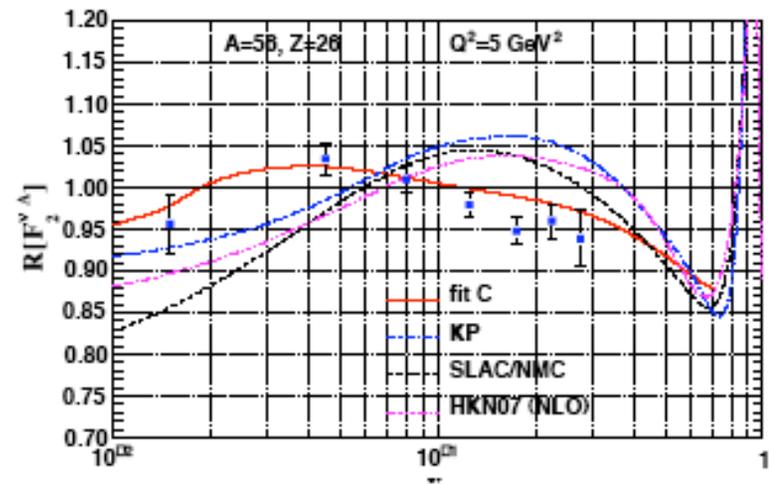
$$W = 0$$

$R[F_2(\ell^\pm \text{ Fe})]$

$R[F_2(\nu \text{ Fe})]$

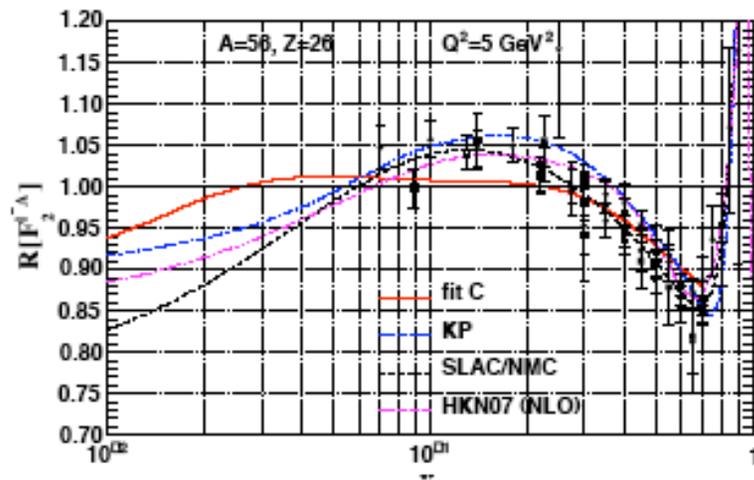


$$W = 1/2$$

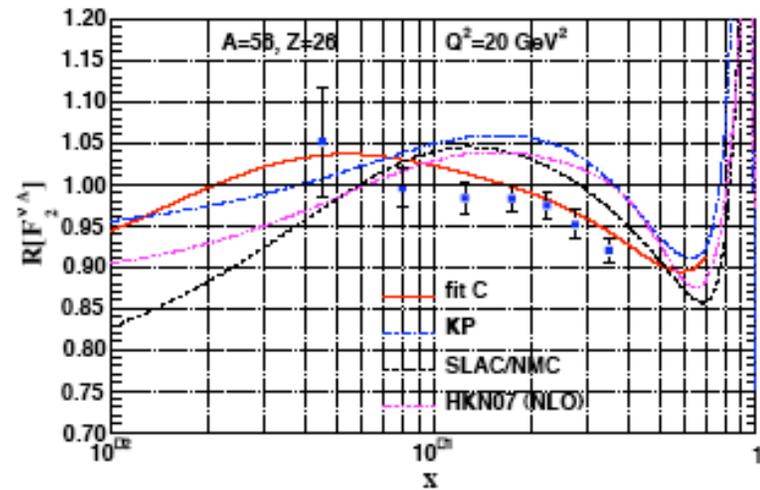
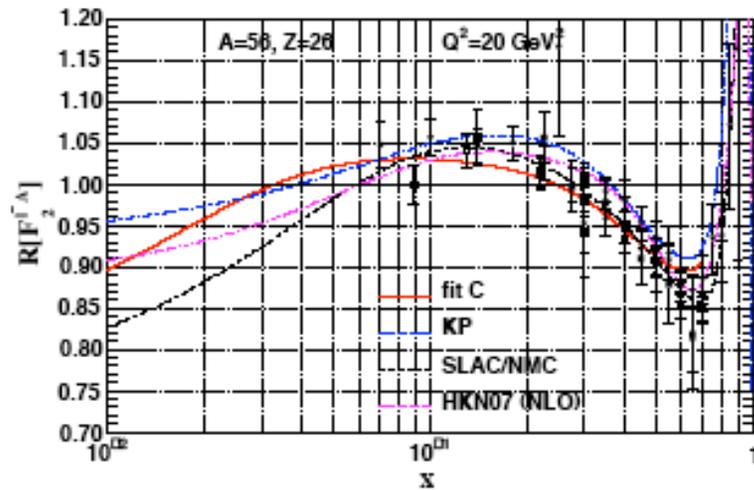
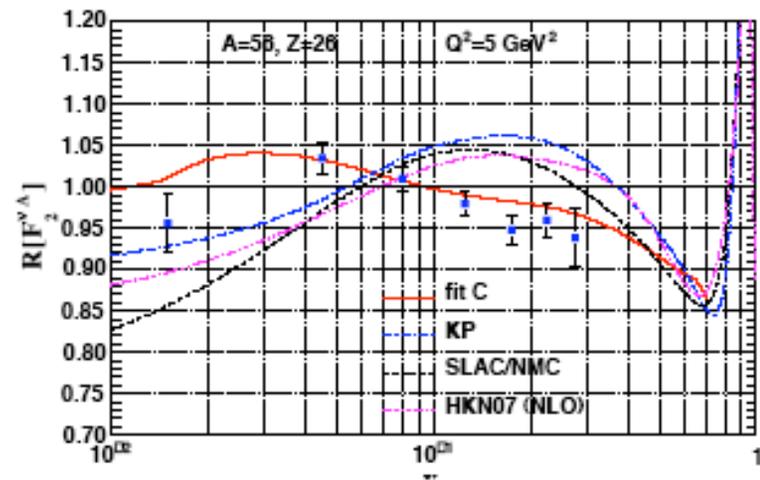
 $R[F_2(\ell^\pm \text{ Fe})]$  $R[F_2(\nu \text{ Fe})]$ 

W = 1

R[F<sub>2</sub>(ℓ<sup>±</sup> Fe)]

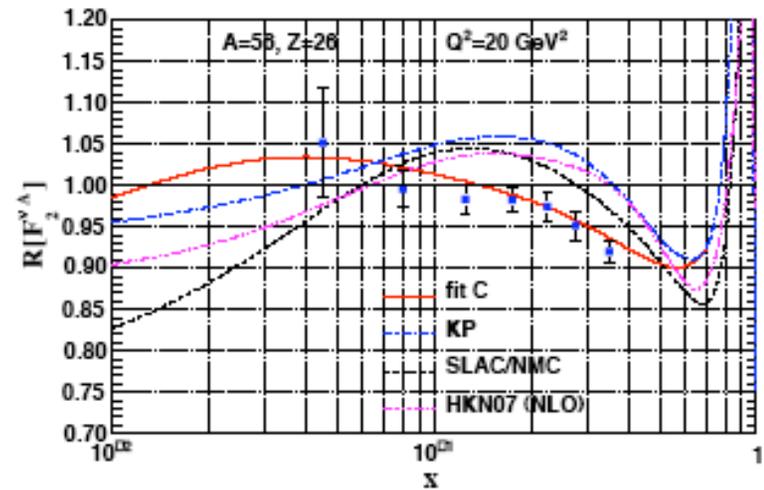
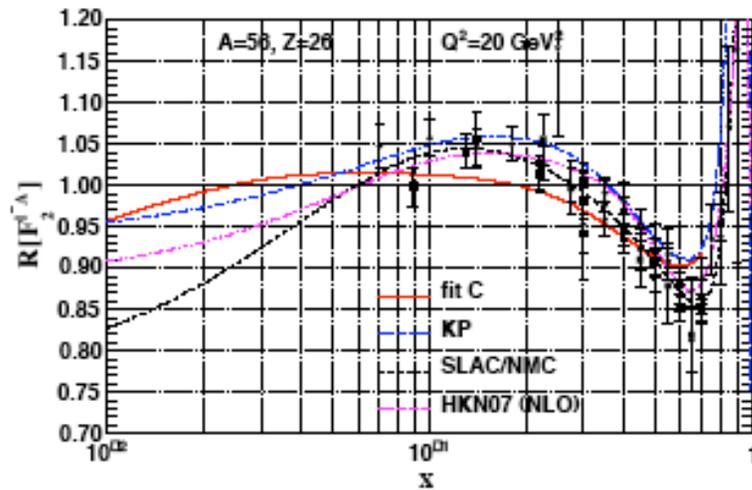
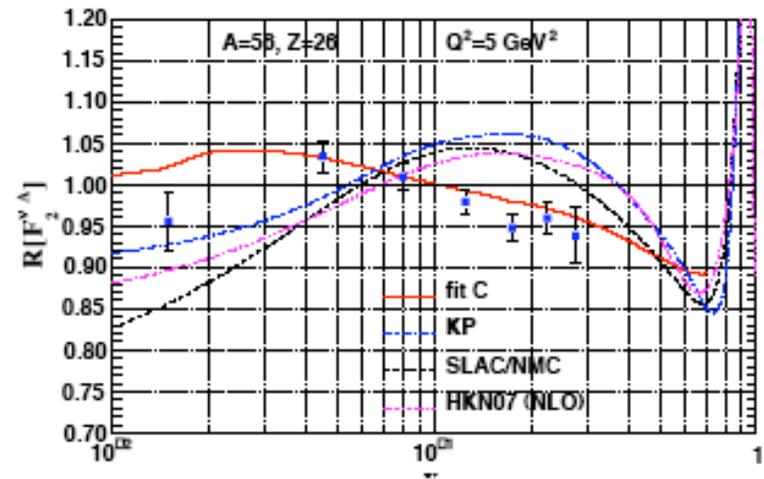
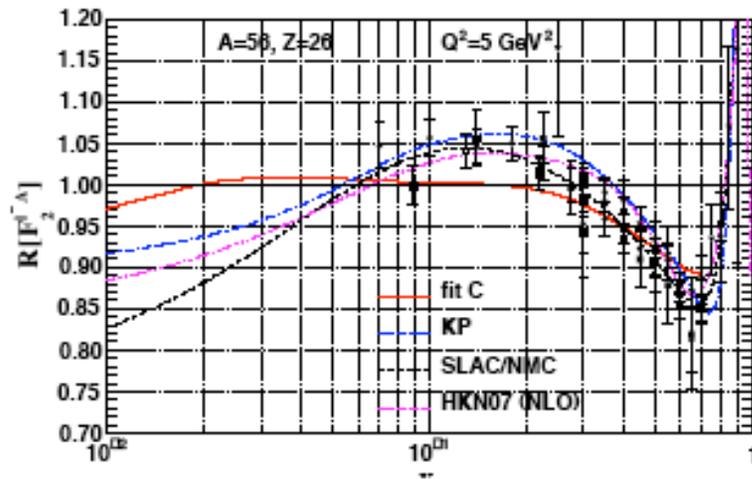


R[F<sub>2</sub>(ν Fe)]



$$W = \infty$$

$$R[F_2(\ell^\pm \text{ Fe})]$$

$$R[F_2(\nu \text{ Fe})]$$


## Fit results

Weight	Fit name	$\ell$ data	$\chi^2$ (/pt)	$\nu$ data	$\chi^2$ (/pt)	total $\chi^2$ (/pt)
$w = 0$	decut3	708	639 (0.90)	-	-	639 (0.90)
$w = 1/7$	glofac1a	708	645 (0.91)	3134	4710 (1.50)	5355 (1.39)
$w = 1/4$	glofac1c	708	654 (0.92)	3134	4501 (1.43)	5155 (1.34)
$w = 1/2$	glofac1b	708	680 (0.96)	3134	4405 (1.40)	5085 (1.32)
$w = 1$	global2b	708	736 (1.04)	3134	4277 (1.36)	5014 (1.30)
$w = \infty$	nuanua1	-	-	3134	4192 (1.33)	4192 (1.33)

- ◆  $w = 0$ : **No**. Problem:  $R[F_2(\nu \text{ Fe})]$ .
- ◆  $w = 1/7$ : **No**. Problem:  $R[F_2(\nu \text{ Fe})]$ .
- ◆  $w = 1/4; 1/2$ : **No**.
  - ▼  $Q2 = 5$ : Undershoots  $R[F_2(\ell^\pm \text{ Fe})]$  for  $x < 0.2$ . Overshoots  $R[F_2(\nu \text{ Fe})]$  for  $x \in [0.1; 0.3]$ .
  - ▼  $Q2 = 20$ :  $R[F_2(\ell^\pm \text{ Fe})]$  still ok. Overshoots  $R[F_2(\nu \text{ Fe})]$ .
- ◆  $w = 1$ : **No**. Possibly there is a compromise if more strict  $Q2$  cut?
  - ▼  $Q2 = 5$ : Undershoots  $R[F_2(\ell^\pm \text{ Fe})]$  for  $x < 0.2$ .  $R[F_2(\nu \text{ Fe})]$  ok.
  - ▼  $Q2 = 20$ :  $R[F_2(\ell^\pm \text{ Fe})]$  still ok.  $R[F_2(\nu \text{ Fe})]$  ok.
- ◆  $w = \infty$ : **No**. Problem:  $R[F_2(\ell^\pm \text{ Fe})]$ .

## Quantitative $\chi^2$ Analysis

- ◆ Up to now we are giving a qualitative analysis. Consider next quantitative criterion based on  $\chi^2$
- ◆ Introduce “tolerance” (T). Condition for compatibility of two fits: The 2nd fit  $\chi^2$  should be within the 90% C.L. region of the first fit  $\chi^2$
- ◆ Charged:  $638:9 \pm 45:6$  (best fit to charged lepton and DY data)
- ◆ Neutrino:  $4192 \pm 138$  (best fit to only neutrino data)

Weight	Fit name	$\ell$ data	$\chi^2$	$\nu$ data	$\chi^2$	total $\chi^2$ (/pt)
$w = 0$	decut3	708	639 T?	-	nnnn NO	639 (0.90)
$w = 1/7$	glofac1a	708	645 YES	3134	4710 NO	5355 (1.39)
$w = 1/4$	glofac1c	708	654 YES	3134	4501 NO	5155 (1.34)
$w = 1/2$	<b>glofac1b</b>	708	680 YES	3134	4405 NO***	5085 (1.32)
$w = 1$	global2b	708	736 NO	3134	4277 YES	5014 (1.30)
$w = \infty$	nuanua1	-	nnn NO	3134	4192	4192 (1.33)

# Summary and Conclusions

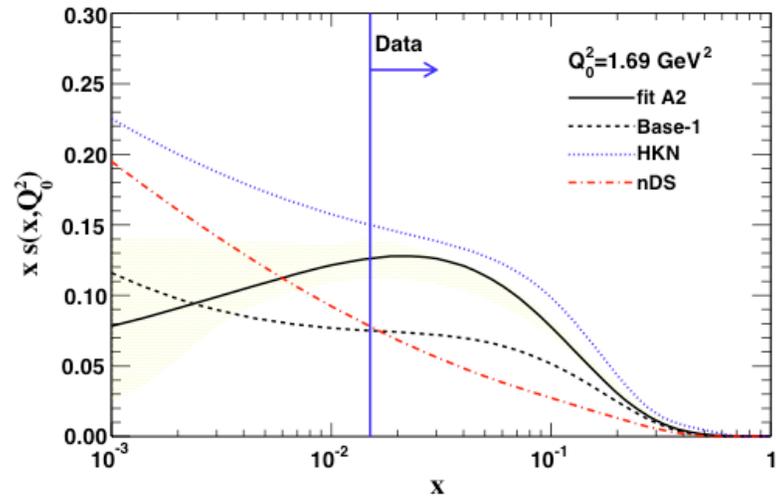
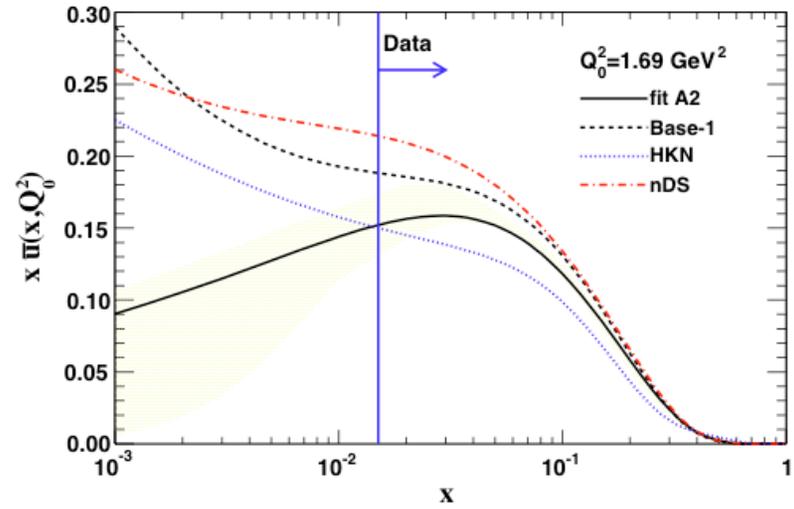
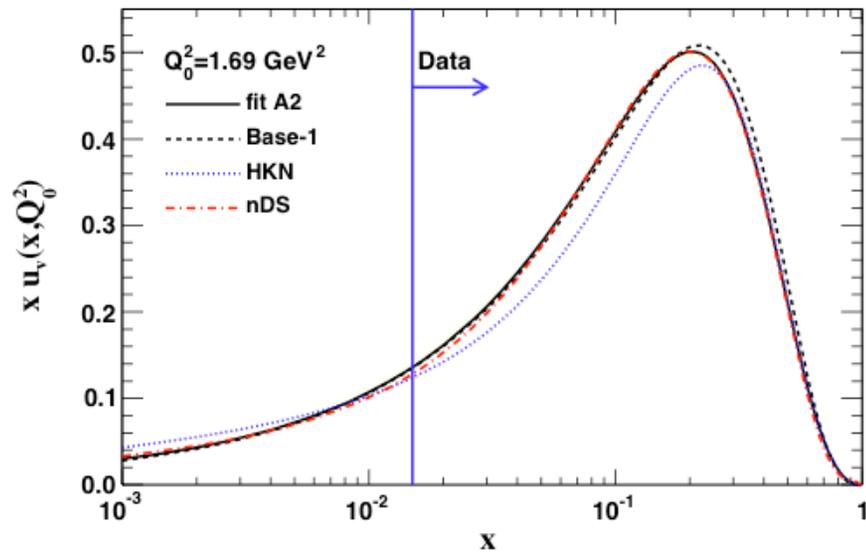
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- ◆ Neutrino scattering can provide an important look at the nucleon from a different (and complimentary) angle than electro-production.
  - ▼ The ability of neutrinos and anti-neutrinos to taste particular flavors of quarks can help isolate PDFs
- ◆ To understand the neutrino (oscillations, mixing, matter effects and  $\delta^{\text{CP}}$ ) neutrino experiments use heavy nuclear targets to obtain statistics. **Need to understand  $\nu$ -induced nuclear effects!**
  - ▼ Use the difference between  $\nu$  and  $\bar{\nu}$  to measure  $\delta^{\text{CP}}$ . **Are  $\nu$  and nuclear effects the same?**
- ◆ There are indications from **one** experiment using **one** nucleus that  **$\nu$ -induced nuclear effects are different** than  $l^{\pm}$ -nuclear effects.
  - ▼ Based on nuclear corrections factors R and the tolerance criterion, there is no good compromise fit to the  $l^{\pm}A + \text{DY} + \nu A$  data.
- ◆ Need a systematic **experimental** study of  **$\nu$ -induced nuclear effects (next talk).**
- ◆ **Need collaborative NP input to fully and correctly analyze crucial high-accuracy neutrino experiments!**

# Additional Details

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# Iron PDFs



# Charged lepton data points

- DIS  $F_2^A/F_2^D$  data sets: 862 points (before cuts)
- DIS  $F_2^A/F_2^{A'}$  data sets: 297 points (before cuts)
- DY data sets  $\sigma_{DY}^{pA}/\sigma_{DY}^{pA'}$ : 92 points (before cuts)

Table from [Hirai et al., arXiv:0909.2329](https://arxiv.org/abs/0909.2329)

	R	Nucleus	Experiment	EPS09	HKN07	DS04	
DIS	A/D	D/p	NMC		0		
		4He	SLAC E139	0	0	0	
				NMC95	0 (5)	0	0
			Li	NMC95	0	0	
			Be	SLAC E139	0	0	0
			C	EMC-88, 90		0	
				NMC 95	0	0	0
				SLAC E139	0	0	0
				FNAL-E665		0	
			N	BCDMS 85		0	
				HERMES 03		0	
			Al	SLAC E49		0	
				SLAC E139	0	0	0
			Ca	EMC 90		0	
				NMC 95	0	0	0
				SLAC E139	0	0	0
				FNAL-E665		0	
			Fe	SLAC E87		0	
				SLAC E139	0 (15)	0	0
				SLAC E140		0	
		BCDMS 87			0		
		Cu	EMC 93	0	0		
		Kr	HERMES 03		0		
		Ag	SLAC E139	0	0	0	
		Sn	EMC 88		0		
		Au	SLAC E139	0	0	0	
			SLAC E140		0		
		Pb	FNAL-E665		0		
		A/C	Be	NMC 96	0	0	0
			Al	NMC 96	0	0	0
			Ca	NMC 95		0	
				NMC 96	0	0	0
	Fe		NMC 96	0	0	0	
	Sn		NMC 96	0 (10)	0	0	
	Pb	NMC 96	0	0	0		
	A/Li	C	NMC 95	0	0		
		Ca	NMC 95	0	0		
DY	A/D	C		0	0	0	
		Ca	FNAL-E772	0 (15)	0	0	
		Fe		0 (15)	0	0	
		W		0 (10)	0	0	
A/Be	Fe	FNAL E866	0	0			
W	0		0				
$\pi$ pro	dA/pp	Au	RHIC-PHENIX	0 (20)			

# Physics Results: Six Structure Functions for Maximal Information on PDF's

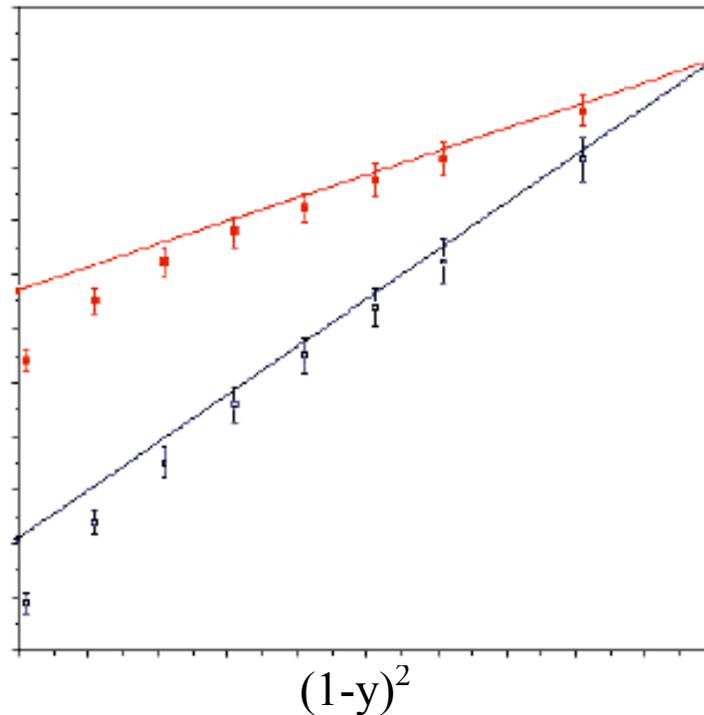
$$\frac{d\sigma^{\nu A}}{dx dQ^2} = \frac{G_F^2}{2\pi X} \left[ \frac{1}{2} (F_2^{\nu A}(x, Q^2) + xF_3^{\nu A}(x, Q^2)) + \frac{(1-y)^2}{2} (F_2^{\nu A}(x, Q^2) - xF_3^{\nu A}(x, Q^2)) \right] + y^2 F_L$$

$$\frac{d\sigma^{\bar{\nu} A}}{dx dQ^2} = \frac{G_F^2}{2\pi X} \left[ \frac{1}{2} (F_2^{\bar{\nu} A}(x, Q^2) - xF_3^{\bar{\nu} A}(x, Q^2)) + \frac{(1-y)^2}{2} (F_2^{\bar{\nu} A}(x, Q^2) + xF_3^{\bar{\nu} A}(x, Q^2)) \right]$$

$$\frac{\sigma(x, Q^2, (1-y)^2)}{G^2/2\pi X}$$

$X = 0.1 - 0.125$   
 $Q^2 = 2 - 4 \text{ GeV}^2$

**Meant to give an impression  
only!**  
**Kinematic cuts in (1-y) not  
shown.**



■ Neutrino  
■ Statistical + 5% systematic

□ Anti-Neutrino  
□ Statistical only

$R = R_{\text{whitlow}}$

# High-x PDFs

## $\nu$ - p Scattering

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$$\left. \begin{aligned} F_2^{\nu p} &= 2x (d + \bar{u} + s) \\ F_2^{\bar{\nu} p} &= 2x (\bar{d} + u + \bar{s}) \end{aligned} \right\} \xrightarrow{\text{At high } x} \boxed{\frac{F_2^{\nu p}}{F_2^{\bar{\nu} p}} = \frac{d}{u}}$$

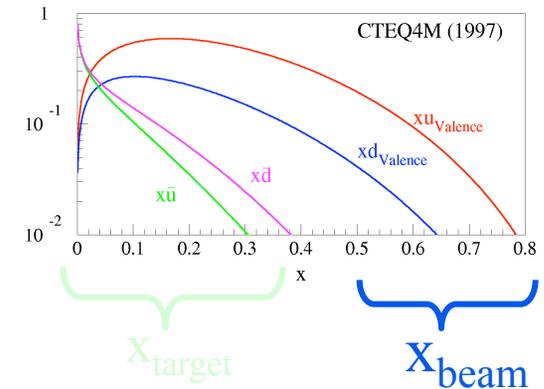
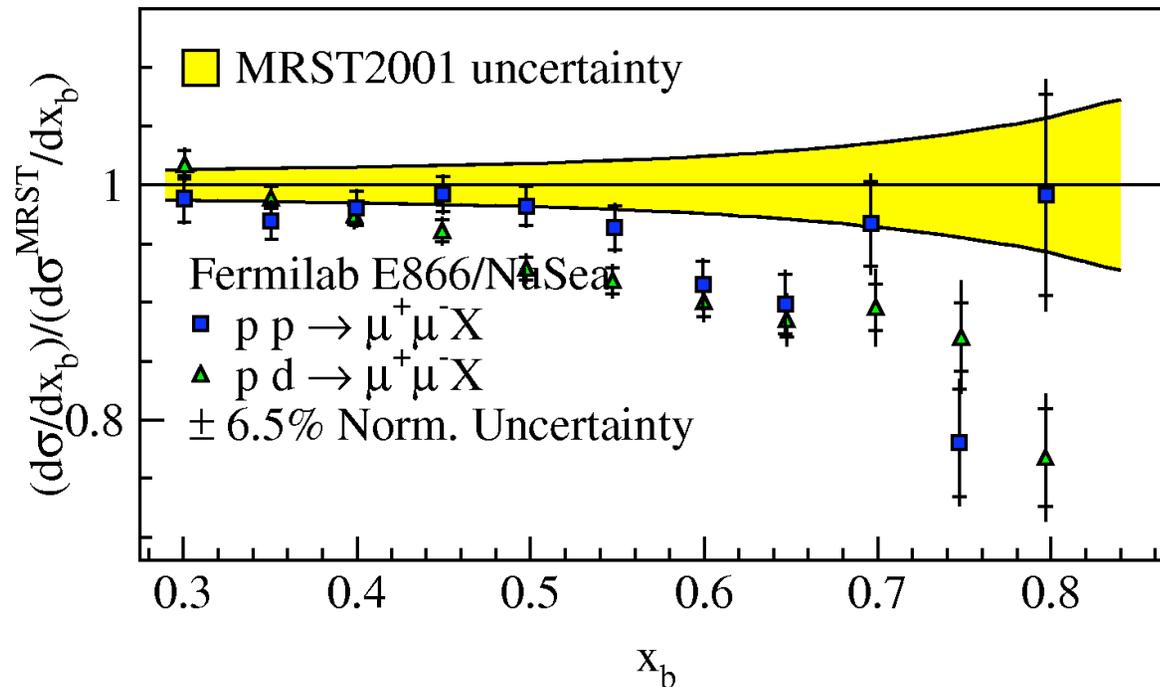
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Add in...

$$\left. \begin{aligned} xF_3^{\nu p} &= 2x (d - \bar{u} + s) \\ xF_3^{\bar{\nu} p} &= 2x (-\bar{d} + u - \bar{s}) \end{aligned} \right\} \longrightarrow \begin{aligned} F_2^{\nu p} - xF_3^{\nu p} &= 4x\bar{u} \\ F_2^{\bar{\nu} p} + xF_3^{\bar{\nu} p} &= 4xu \end{aligned}$$

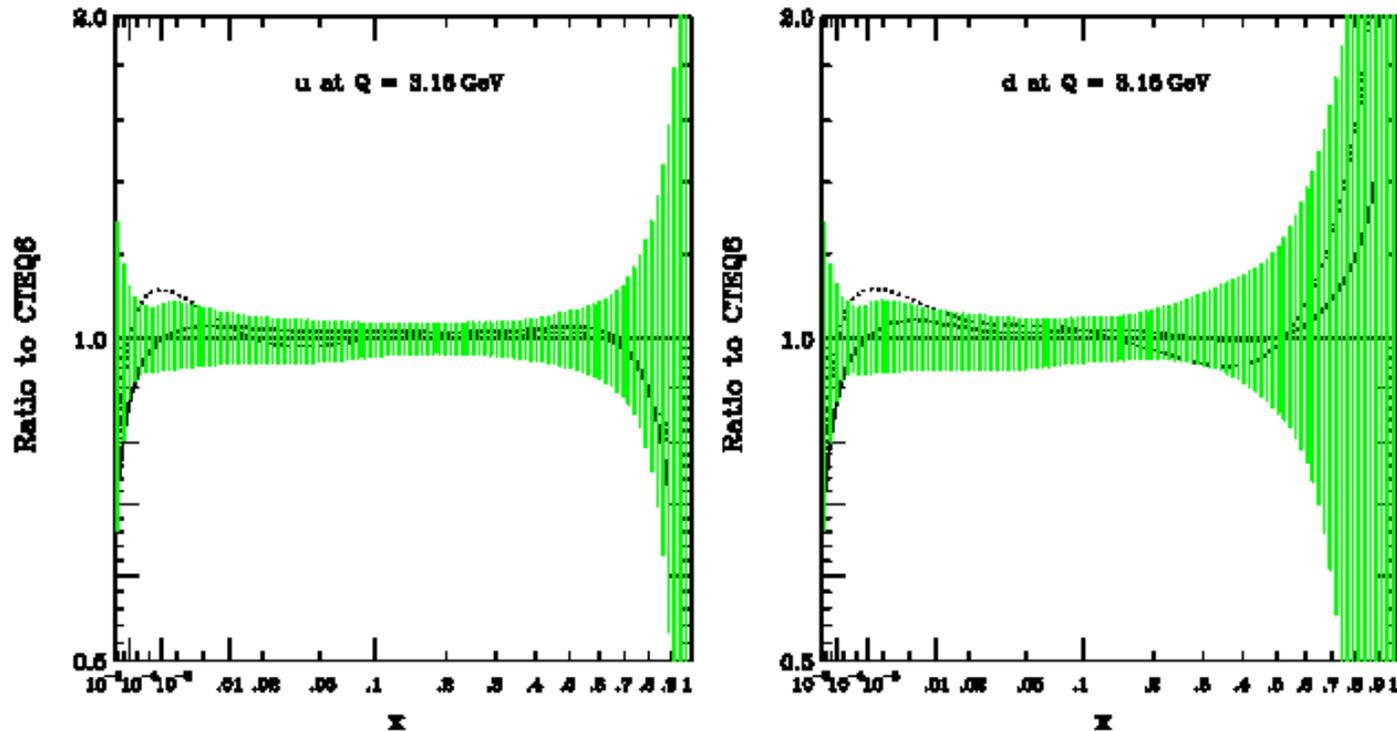
# Further indications that the valence quarks not quite right at high- $x$ ??

E866 -Drell-Yan Preliminary Results (R. Towell - Hix2004)



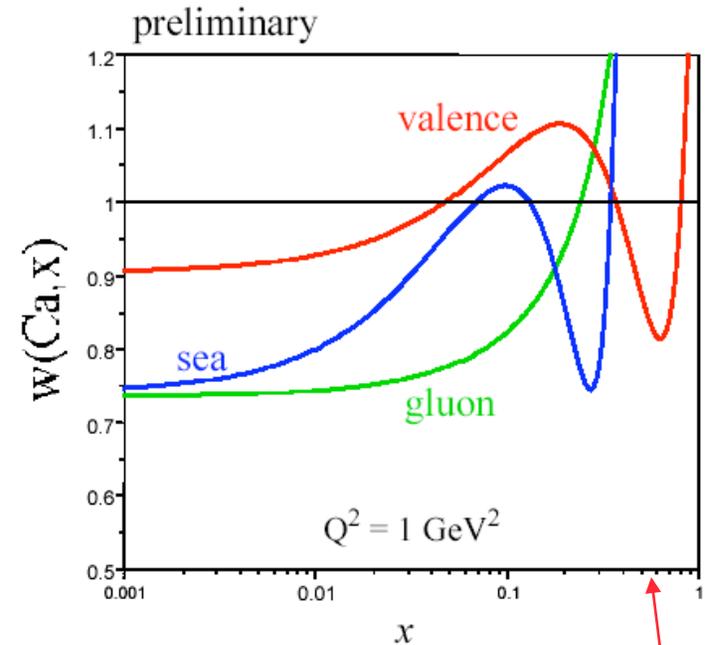
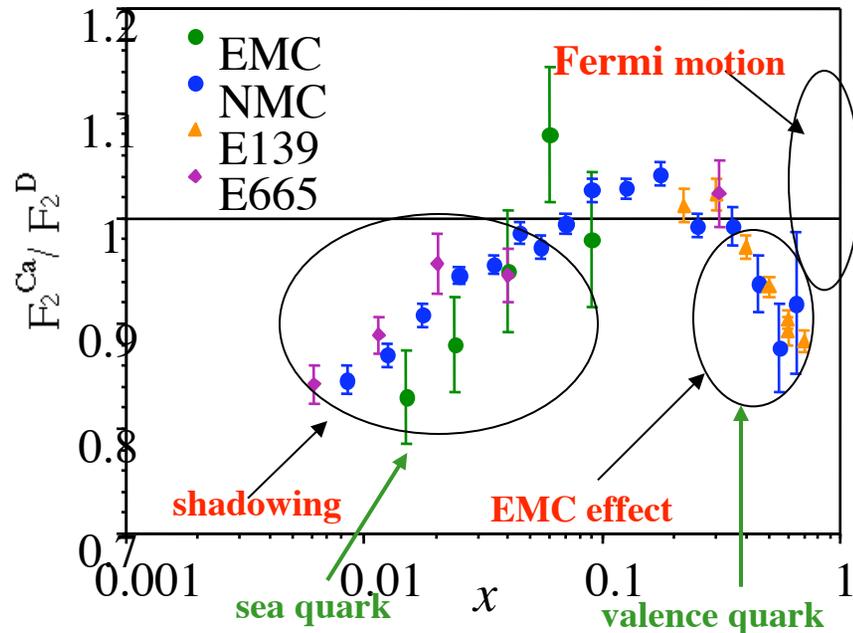
- $x_{\text{beam}}$  distribution measures  $4u + d$  as  $x \rightarrow 1$ .
- Both MRST and CTEQ overestimate valence distributions as  $x \rightarrow 1$  by 15-20%.
- Possibly related to  $d/u$  ratio as  $x \rightarrow 1$ , but requires full PDF-style fit.
- Radiative corrections have recently been calculated. (Not yet fully applied)

# Present Status: $\nu$ -scattering High $x_{Bj}$ parton distributions



- ◆ Ratio of CTEQ5M (solid) and MRST2001 (dotted) to CTEQ6 for the u and d quarks at  $Q^2 = 10 \text{ GeV}^2$ . The shaded green envelopes demonstrate the range of possible distributions from the CTEQ6 error analysis.
- ◆ CTEQ / MINERvA working group to investigate high- $x_{Bj}$  region.

# Knowledge of Nuclear Effects with Neutrinos: essentially NON-EXISTENT



- ◆  $F_2$  / nucleon changes as a function of  $A$ . Measured in  $\mu/e - A$  not in  $\nu - A$
- ◆ Good reason to consider nuclear effects are DIFFERENT in  $\nu - A$ .
  - ▼ Presence of axial-vector current.
  - ▼ SPECULATION: Much stronger shadowing for  $\nu - A$  but somewhat weaker “EMC” effect.
  - ▼ Different nuclear effects for valence and sea --> different shadowing for  $xF_3$  compared to  $F_2$ .
  - ▼ Different nuclear effects for d and u quarks.

# Formalism

- ◆ PDF Parameterized at  $Q_0 = 1.3 \text{ GeV}$  as

$$xf_i(x, Q_0) = \begin{cases} A_0 x^{A_1} (1-x)^{A_2} e^{A_3 x} (1+e^{A_4 x})^{A_5} & : i = u_v, d_v, g, \bar{u} + \bar{d}, s, \bar{s}, \\ A_0 x^{A_1} (1-x)^{A_2} + (1+A_3 x)(1-x)^{A_4} & : i = \bar{d}/\bar{u}, \end{cases}$$

- ◆ PDFs for a nucleus are constructed as:

$$f_i^A(x, Q) = \frac{Z}{A} f_i^{p/A}(x, Q) + \frac{(A-Z)}{A} f_i^{n/A}(x, Q)$$

- ◆ Resulting in nuclear structure functions:

$$F_i^A(x, Q) = \frac{Z}{A} F_i^{p/A}(x, Q) + \frac{(A-Z)}{A} F_i^{n/A}(x, Q)$$

- ◆ The differential cross sections for CC scattering off a nucleus::

$$\begin{aligned} \frac{d^2\sigma}{dx dy} {}^{(\bar{\nu})A} &= \frac{G^2 ME}{\pi} \left[ \left(1 - y - \frac{Mxy}{2E}\right) F_2 {}^{(\bar{\nu})A} \right. \\ &\quad \left. + \frac{y^2}{2} 2xF_1 {}^{(\bar{\nu})A} \pm y\left(1 - \frac{y}{2}\right) xF_3 {}^{(\bar{\nu})A} \right] \end{aligned}$$