

MECCANICA STATISTICA DEI PROBLEMI DI OTTIMIZZAZIONE COMBINATORIA

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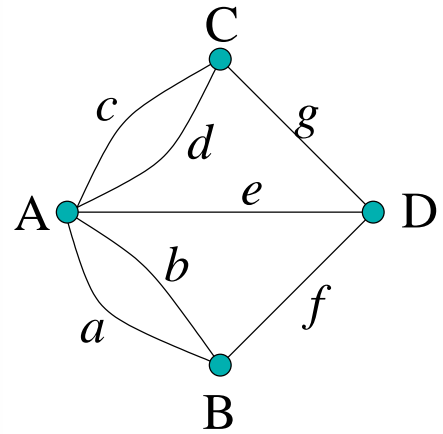
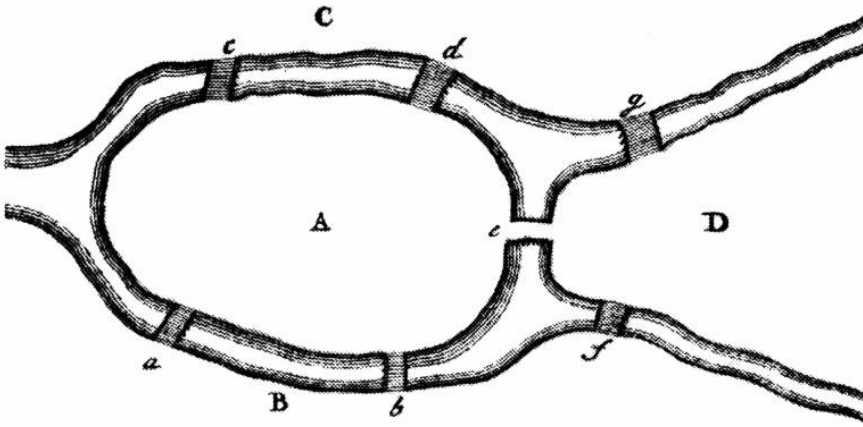
Marc Mézard (Orsay)

Martin Weigt (Göttingen)

-
- Breve introduzione alla Complessità Computazionale
 - Legame con la Meccanica Statistica
 - Semplici nozioni sugli ipergrafi random
 - Analisi di un modello semplice: 3-spin diluito
 - Conclusioni, applicazioni e prospettive future

Two old decision problems

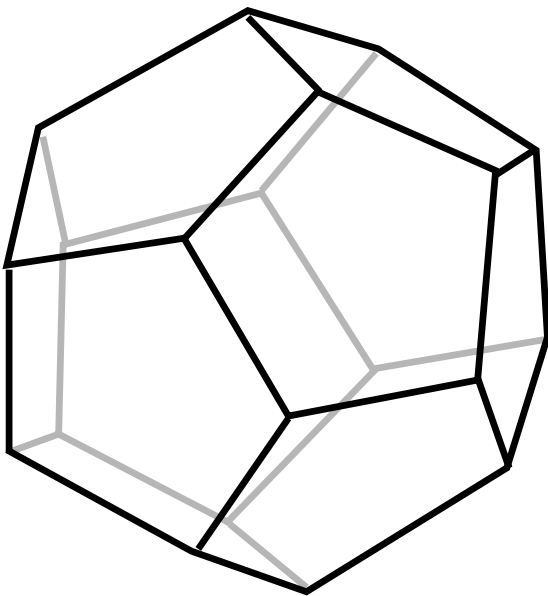
Eulerian circuits (1736)



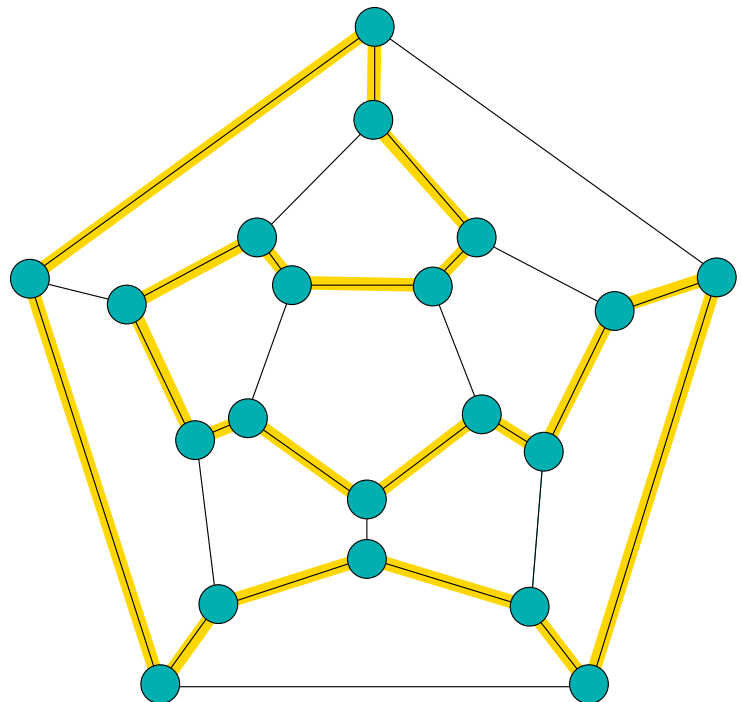
The seven bridges of Königsberg

The equivalent graph

Hamiltonian cycles (1859)



The dodecahedron



The equivalent graph

Classifying problems according to the computational resources required for their solution (e.g. CPU time and memory) in the **worst case**.

tractable (in **P**) \leftrightarrow sub-exponential: $\ln(N)$, N^α

intractable (in **NP**) \leftrightarrow exponential: 2^N , $N!$

N is the number of variables in the problem.

Main Complexity Classes

P = polynomial

NP = non-deterministic polynomial

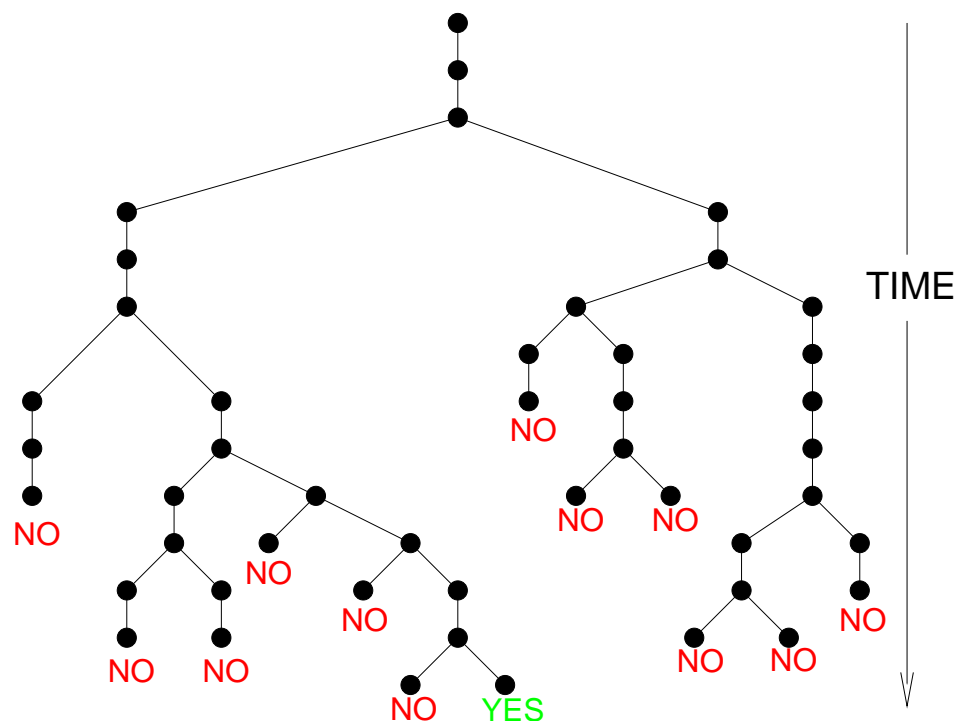


- non-determ. algorithms: goto both
- succinct certificate

P = NP ?

First Millennium Prize Problem

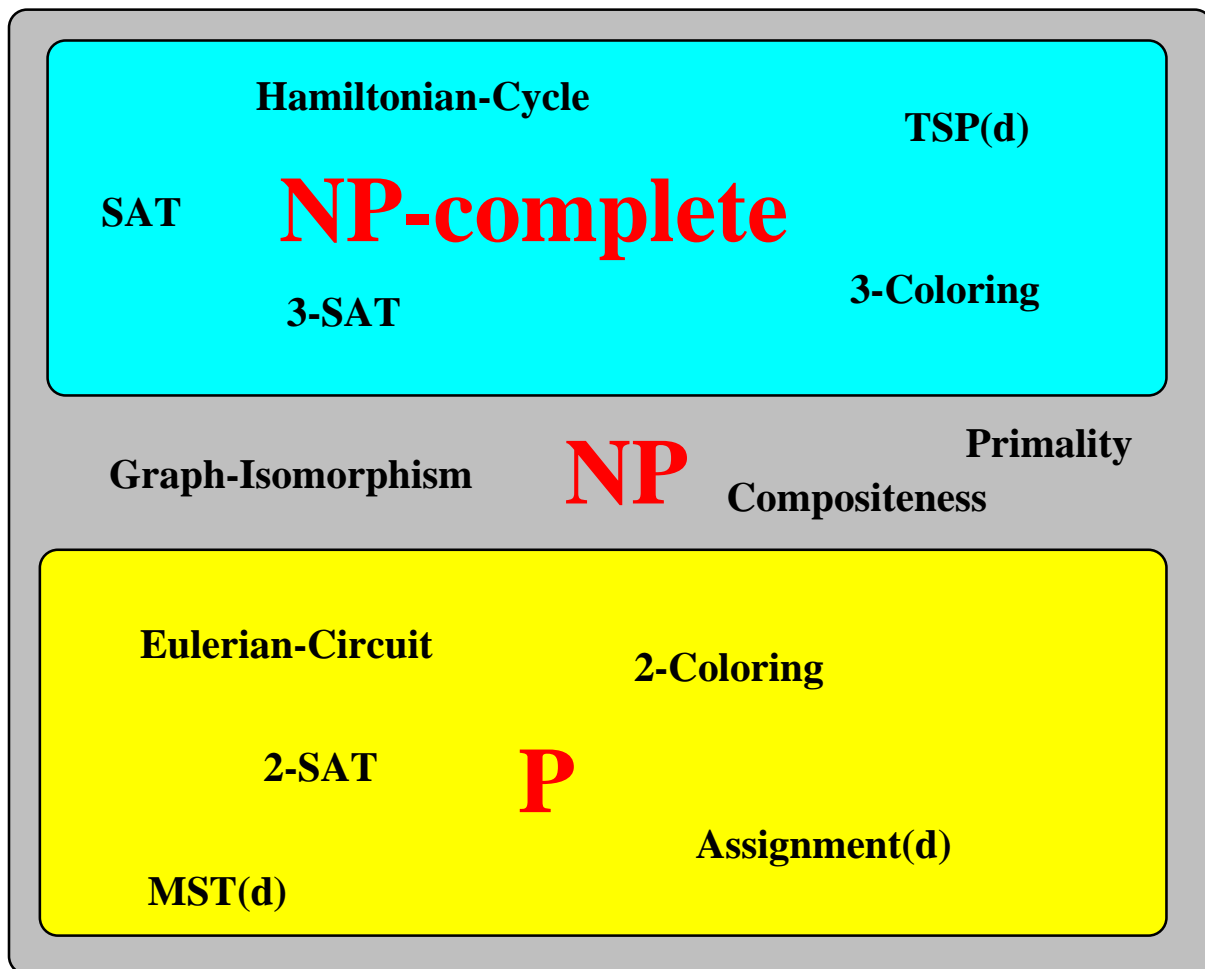
Quantum computation !



NP-completeness: combinatorial optimization

THEOREM: All problems in **NP** are polynomially reducible to **K -SAT** ($K > 2$). (Cook, 1971)

NP-complete are the **hardest NP** problems.



Beyond NP

- $\left\{ \begin{array}{l} \text{decision} \\ \text{evaluation} \\ \text{optimization} \end{array} \right\}$ problems (**NP-hardness**)
- **counting** problems (**#P, #P-complete**):
e.g. #SAT counts the entropy of SAT.

Some interesting problems

SAT

Boolean variables: $x_i \in \{0, 1\}$ $i = 1, \dots, N$

Logical operators: $\bar{\cdot} = \text{NOT}$, $\vee = \text{OR}$, $\wedge = \text{AND}$

Given a Boolean formula like

$(x_1 \vee x_{27} \vee \bar{x}_3) \wedge (\bar{x}_{11} \vee x_2) \wedge \dots \wedge (\bar{x}_{21} \vee x_9 \vee \bar{x}_8 \vee \bar{x}_{30})$

decide if a **satisfying** truth assignment exist

Random 3-SAT: 3 random literals per clause

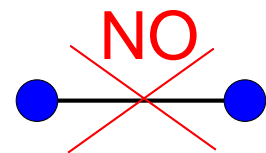
clauses = αN

Coloring

q -coloring on graphs:

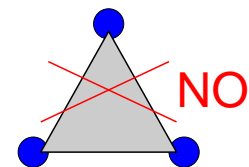
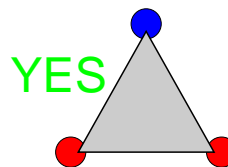
2-coloring $\in P$

3-coloring $\in \text{NP-c}$



Bicoloring on hypergraphs:

NP-complete



Computer Science

Physics

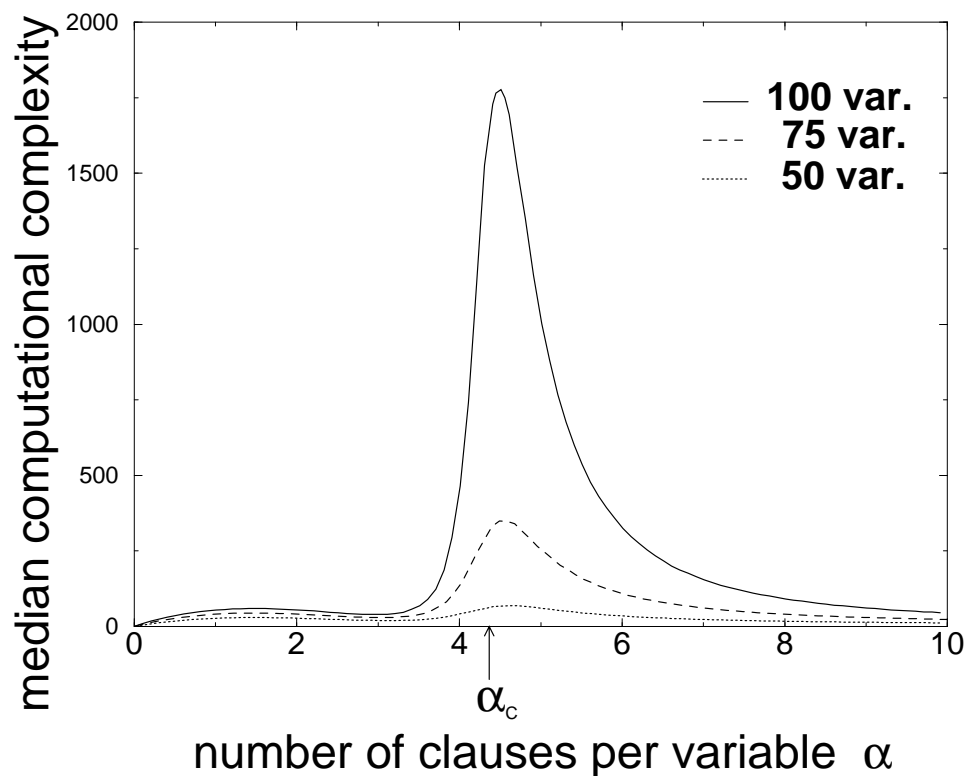
K -SAT
Graph Coloring
Vertex Cover
...

Diluted Ising Spin Glasses
Potts Models
Hard Spheres
...

Real-world NP-complete and #P-complete problems may have many easy instances

Ensemble of random NP-complete problems hard **on average** (e.g. random 3-SAT)

Random models: phase transitions, NP-hardness



Onset of complexity \leftrightarrow Phase transition

SAT/UNSAT transition: $E = \#$ violated clauses becomes greater than zero.

What makes problems hard close to α_c ?

Mapping to a Statistical Mechanics problem

Zero temperature limit of a **diluted** mean-field spin model. Much harder than usual fully connected!

For **random 3-SAT**, with $s_i = (-1)^{x_i}$,

$$\mathcal{H} = \frac{1}{8} \left(\alpha N - \sum_{i=1}^N H_i s_i + \sum_{i < j} T_{ij} s_i s_j - \sum_{i < j < k} J_{ijk} s_i s_j s_k \right)$$

$$\begin{aligned} H_i &= \sum_{\ell} \Delta_{\ell,i} \\ T_{ij} &= \sum_{\ell} \Delta_{\ell,i} \Delta_{\ell,j} \\ J_{ijk} &= \sum_{\ell} \Delta_{\ell,i} \Delta_{\ell,j} \Delta_{\ell,k} \end{aligned} \quad \Delta_{\ell,i} = \begin{cases} 1 & \text{if } \bar{x}_i \in C_{\ell} \\ -1 & \text{if } x_i \in C_{\ell} \\ 0 & \text{otherwise} \end{cases}$$

For **3-hyper-SAT** also known as 3-XOR-SAT

$$F = (x_2 \oplus x_{15} \oplus x_{33}) \wedge \dots \wedge (\bar{x}_4 \oplus \bar{x}_{21} \oplus \bar{x}_9)$$

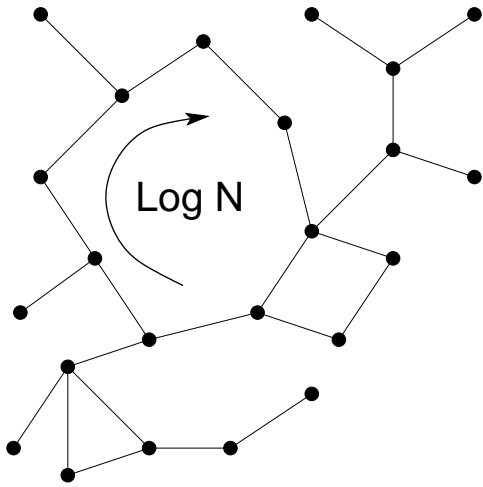
$$\mathcal{H} = \frac{1}{2} \left(\gamma N - \sum_{\{i,j,k\} \in G} J_{ijk} s_i s_j s_k \right)$$

$G = \{\text{set of } \gamma N \text{ random triples}\}$

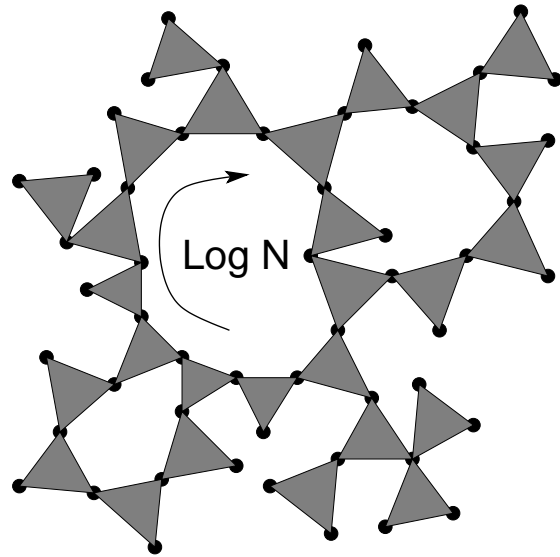
Two versions:

- **unfrustrated**, ferromagnetic: $J_{ijk} = 1$
→ 1st order ferromagnetic transition
- **frustrated**, 3-spin glass: $J_{ijk} = \pm 1$
→ SAT/UNSAT transition

Graphs and Hypergraphs



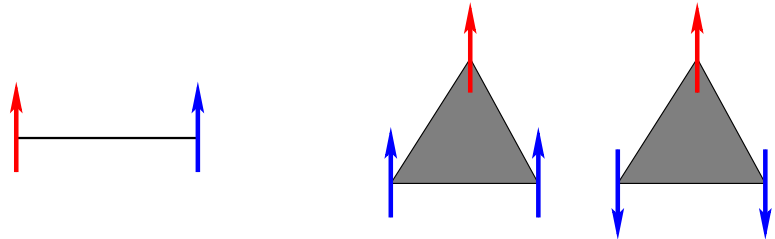
Graph (2-spin interactions)



Hypergraph (3-spin interactions)

Typically they are **non planar** and **long range**:
typical loops are of order $\ln(N)$

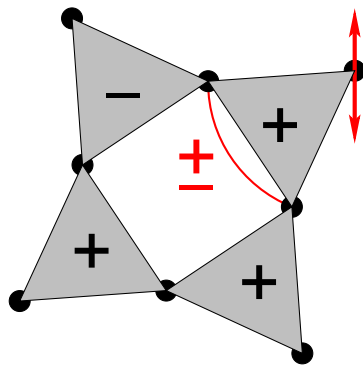
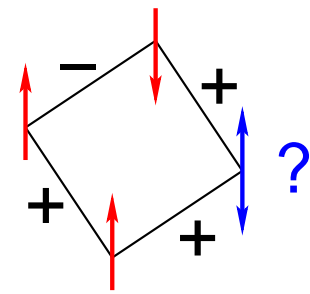
How to satisfy
a ferromagnetic
interaction



Finding the ground state of an unfrustrated spin
model is an **easy** problem on a **graph**, but it may
be **hard** on a **hypergraph**

Very different $T = 0$ dynamics (**coarsening**)

On a **graph** the **frustration** arises with **loops** at the percolation point ($\gamma_p = 1/2$). On a **hypergraph** the loops arising at the percolation point ($\gamma_p = 1/6$) give no frustration.

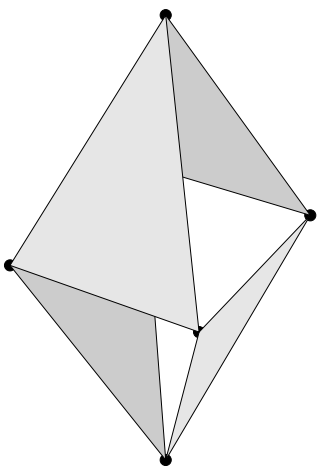


Spin on dangling ends can be freely fixed in order to change the effective interaction between the other two spins. Only **hyperloops** can generate frustration on a hypergraph.

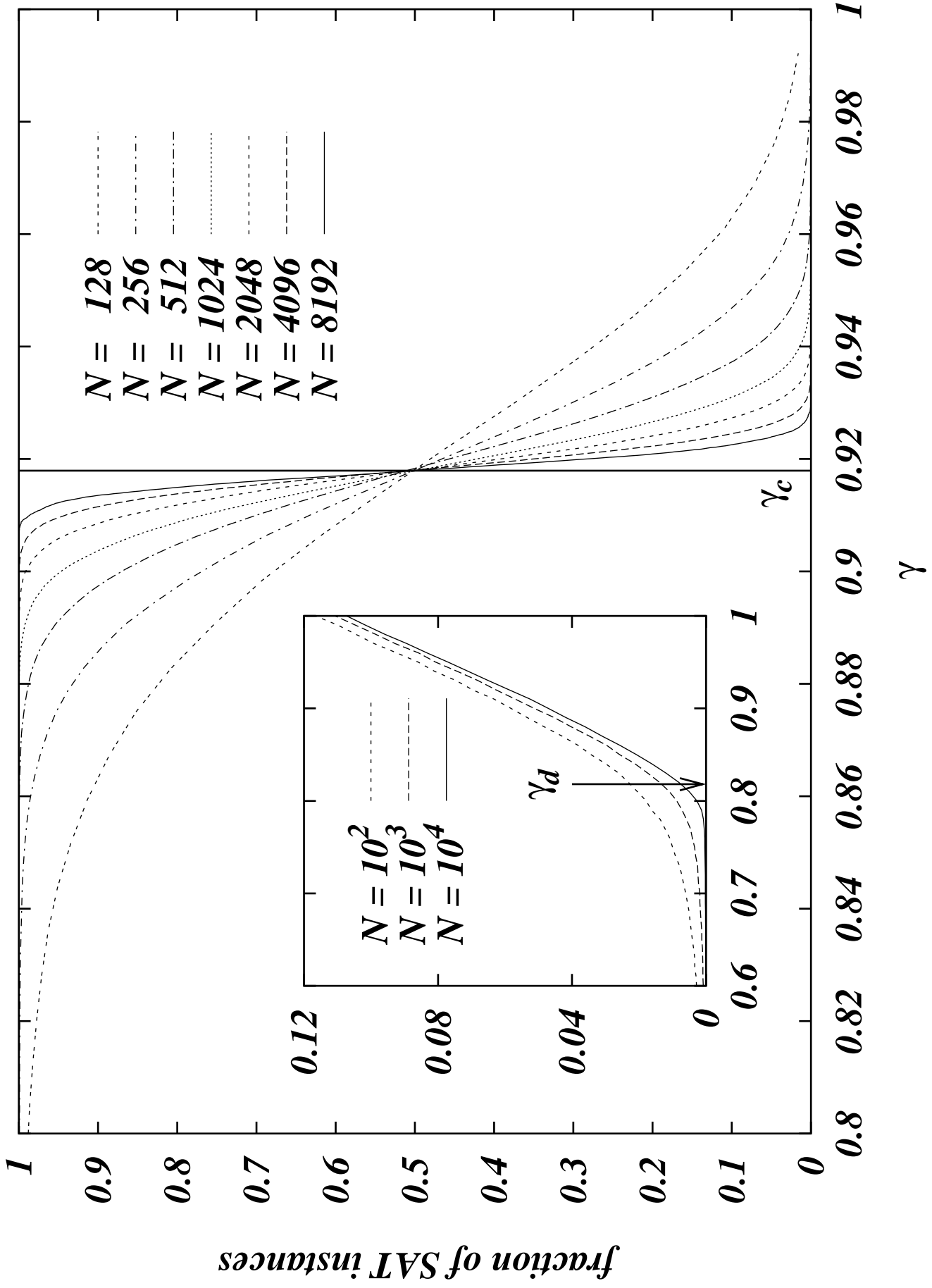
Definition of hyperloop: a non-empty set \mathcal{S} of hyperlinks $\{i, j, k\}$ such that every node (spin) appears in \mathcal{S} an **even** number of times (**zero is even**).

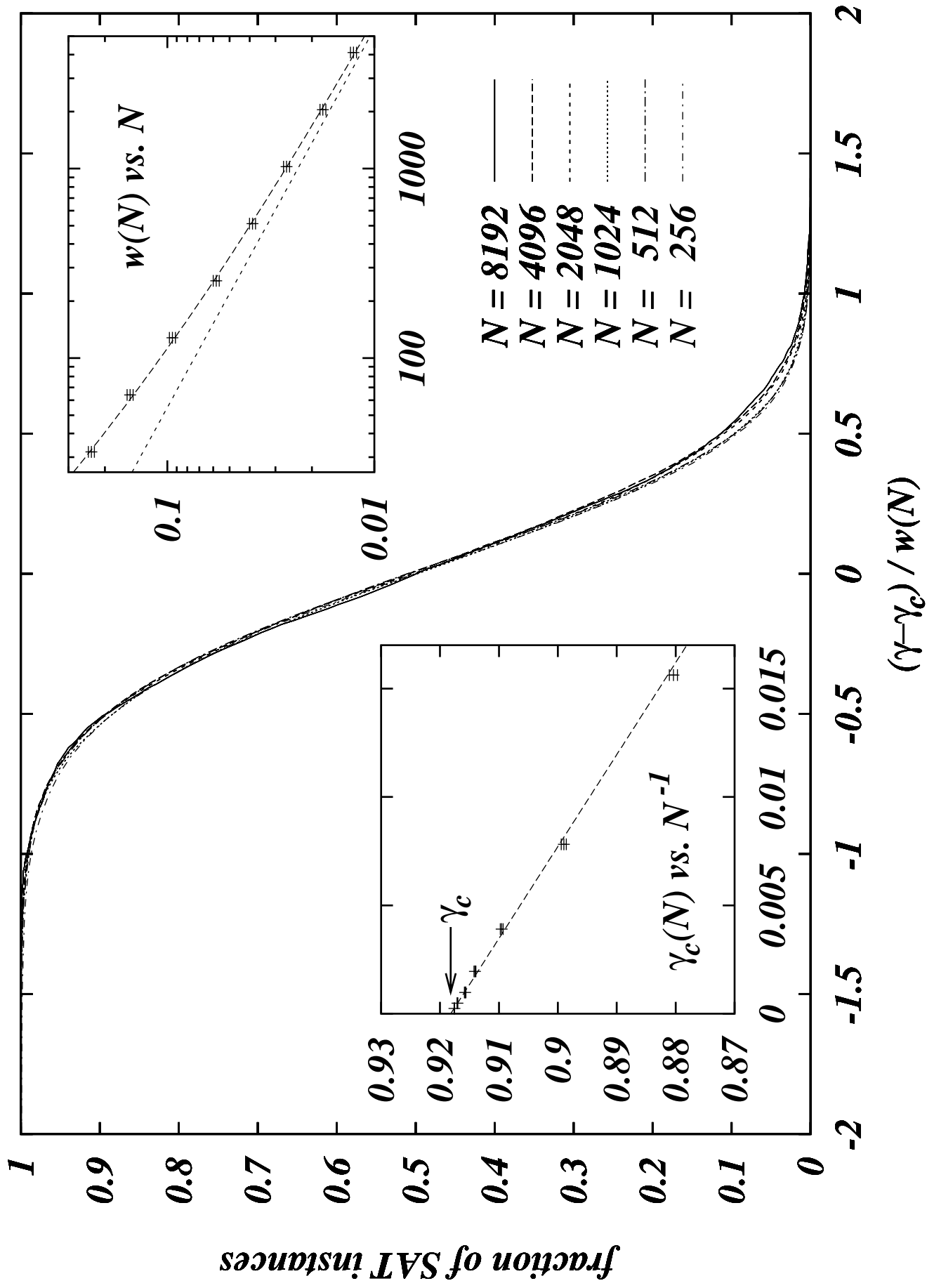
hyperloops \Rightarrow frustration

satisfied interaction: $J_{ijk} s_i s_j s_k = 1$



$$\begin{aligned}
 1 &= \prod_{\{i,j,k\} \in \mathcal{S}} J_{ijk} s_i s_j s_k = \\
 &= \prod_{\{i,j,k\} \in \mathcal{S}} J_{ijk} = \begin{cases} 1 & p = 1/2 \\ -1 & p = 1/2 \end{cases}
 \end{aligned}$$





Unfrustrated model: $J = 1$

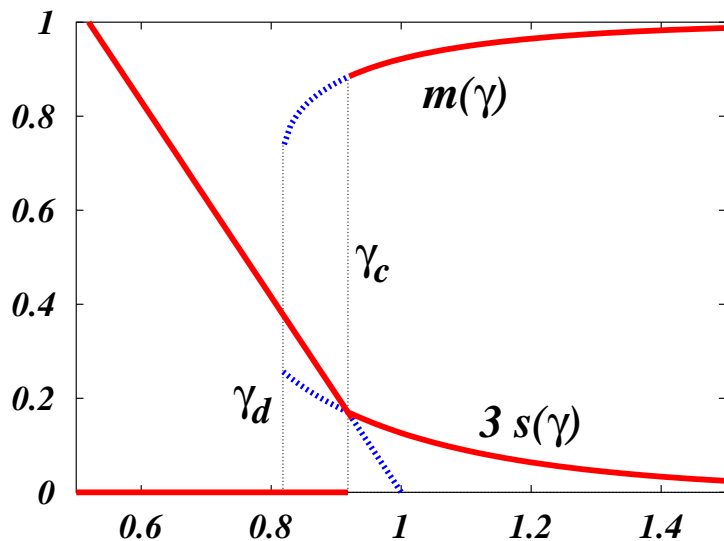
The existence of a ferromagnetic solution follows from a simple argument

$$1 - m = \sum_c e^{-3\gamma} \frac{(3\gamma)^c}{c!} \frac{c}{3\gamma} (1 - m^2)^{c-1} = e^{-3\gamma m^2}$$

Up to $\gamma_d = 0.818$ the hypergraph is like a **tree**

$m(\gamma)$ is the fraction of frozen spins, i.e. the magnetization

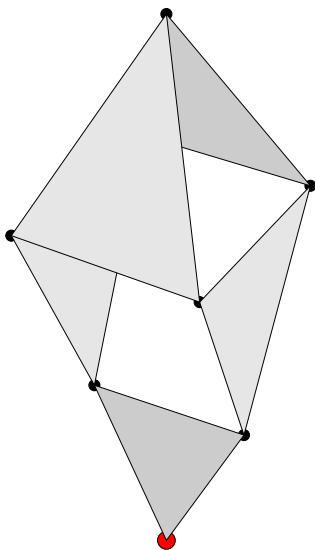
$s(\gamma)$ is the entropy
 $\gamma_c = 0.918$



The difficult task is to calculate $s(\gamma) \rightarrow$ **replicas!**

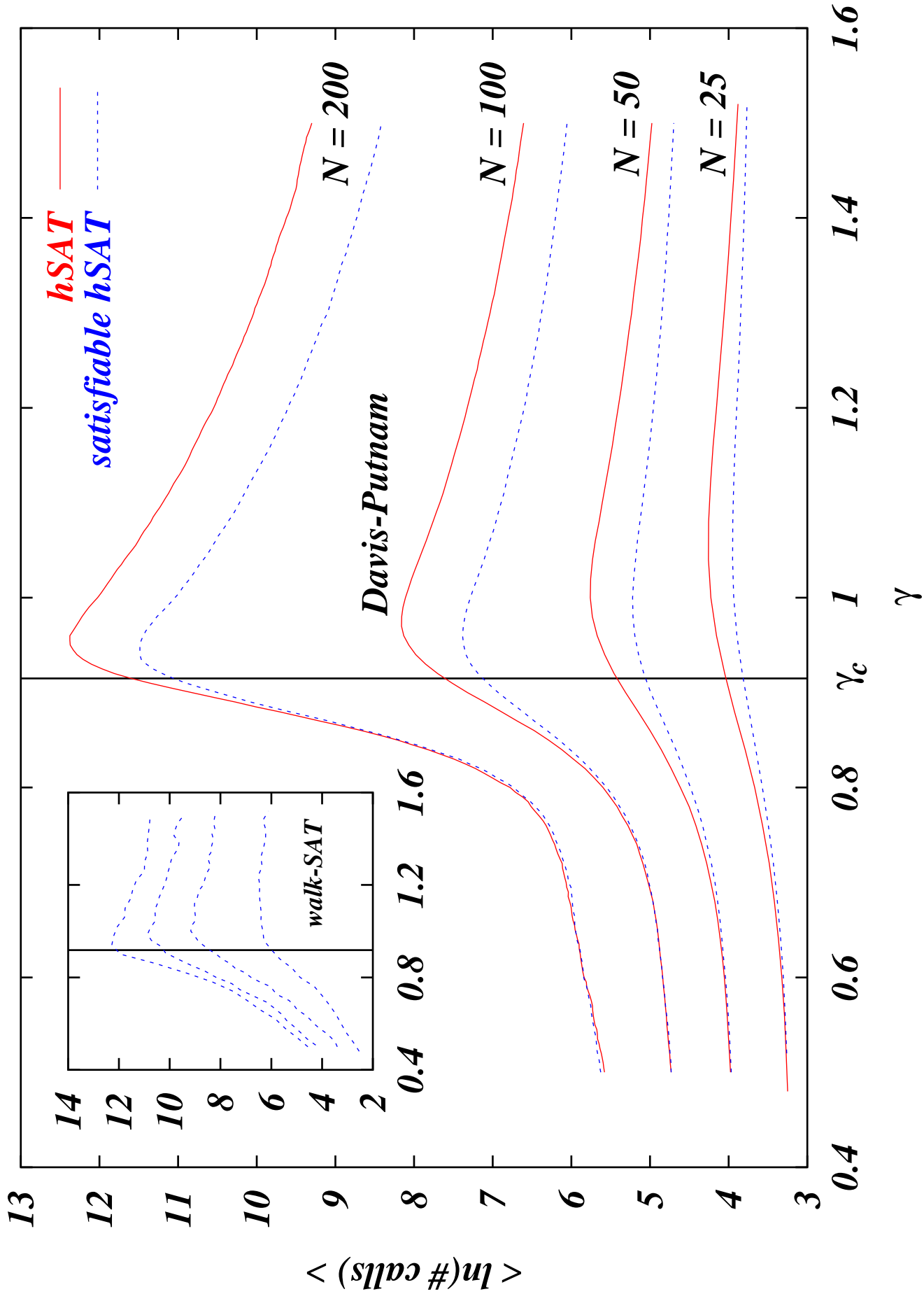
$$s(\gamma) = \ln(2)[1 - \gamma - m + 3\gamma m^2(1 - m) + \gamma m^3]$$

Challenge for mathematicians: calculate γ_c without replicas



At γ_c percolate hyperloops and **hyperfields** \Rightarrow **magnetization**

$$\begin{aligned} s_l &= \prod_{\{i,j,k\} \in \mathcal{T}} s_i s_j s_k = \\ &= \prod_{\{i,j,k\} \in \mathcal{T}} J_{ijk} = 1 \end{aligned}$$



The structure of the configurational space

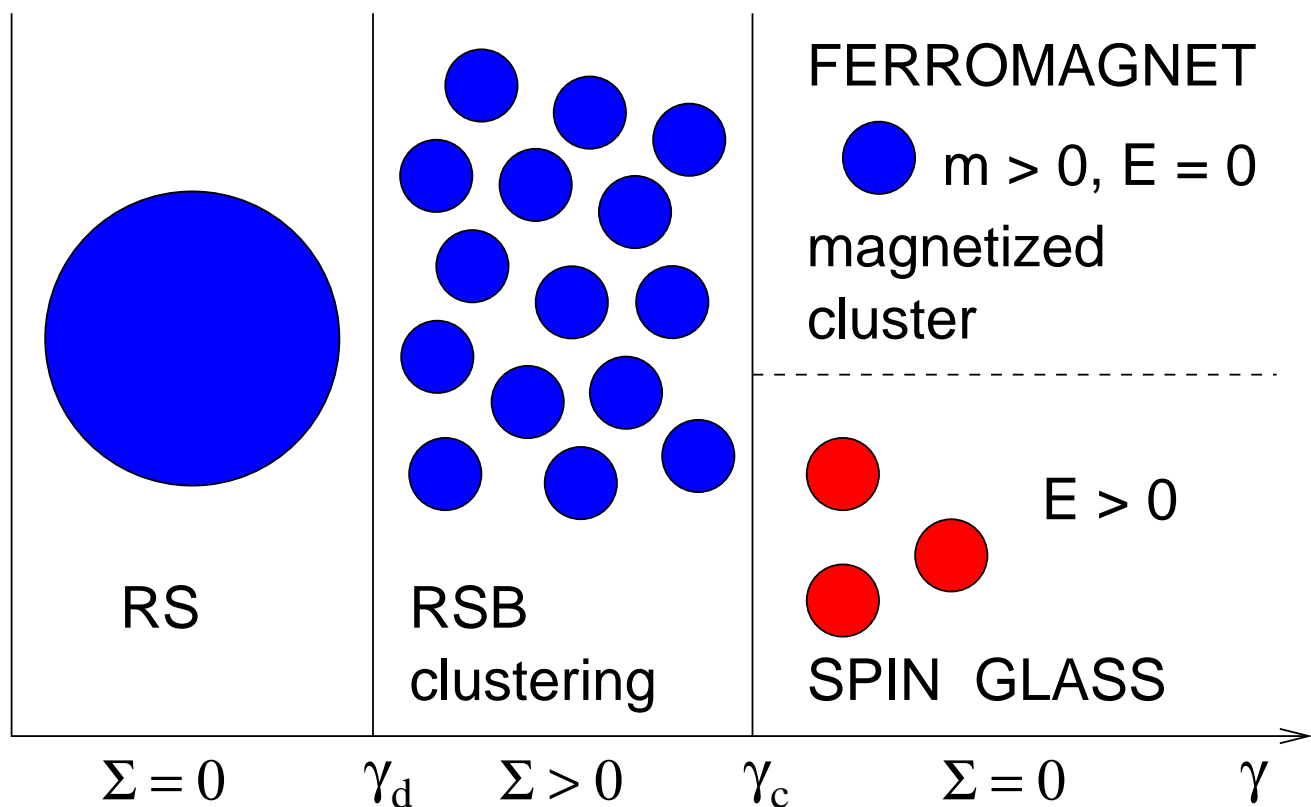
$$x_i, \eta_{ijk} \in \{0, 1\} \quad s_i = (-1)^{x_i} \quad J_{ijk} = (-1)^{\eta_{ijk}}$$

$$\{s_i\} \text{ is a solution} \iff \forall \{i, j, k\} \in I \quad s_i s_j s_k = J_{ijk} \iff$$

$$\iff (x_i \oplus x_j \oplus x_k) = \boxed{x_i + x_j + x_k \pmod{2} = \eta_{ijk}}$$

A common problem: diluted p-spin glass \equiv random p-XOR-SAT* (or low density Parity Check codes) \equiv random linear systems in finite fields (GF[2])
 *considered an open problem in theoretical computer science

If $E = 0$ (no frustration) \rightarrow Gaussian elimination
 If $E > 0$ \rightarrow exhaustive enumerations



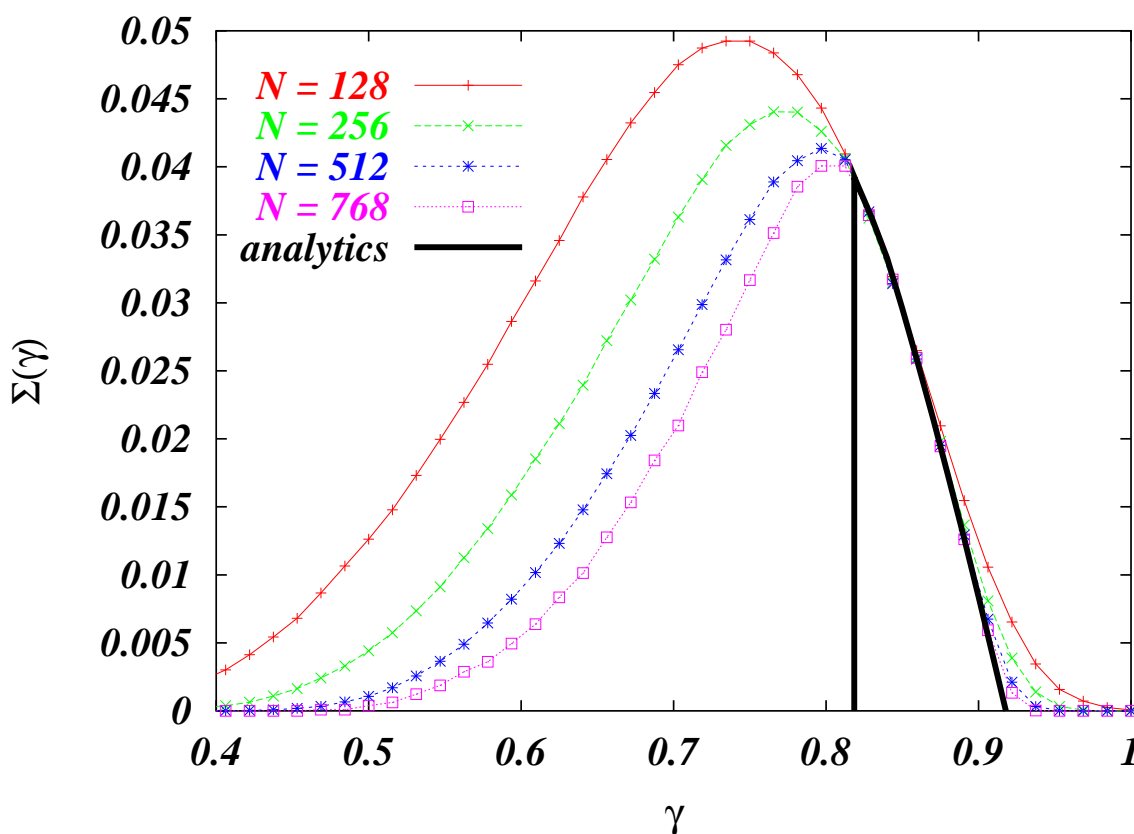
Configurational entropy: $\Sigma(\gamma) = \frac{1}{N} \ln(\# \text{ clusters})$

The structure of the configurational space

Configurational entropy: $\Sigma(\gamma) = \frac{1}{N} \ln(\# \text{ clusters})$
Since all clusters are equal

$$\begin{aligned}\Sigma(\gamma) &= \ln(2)(1 - \gamma) - s(\gamma) = \\ &= \ln(2) \left[r - 3\gamma r^2(1 - r) - \gamma r^3 \right]\end{aligned}$$

with $1 - r = \exp(-3\gamma r^2)$, is **exact** !



$\Sigma(\gamma) > 0$ for $\gamma \in [\gamma_d, \gamma_c] \Rightarrow$

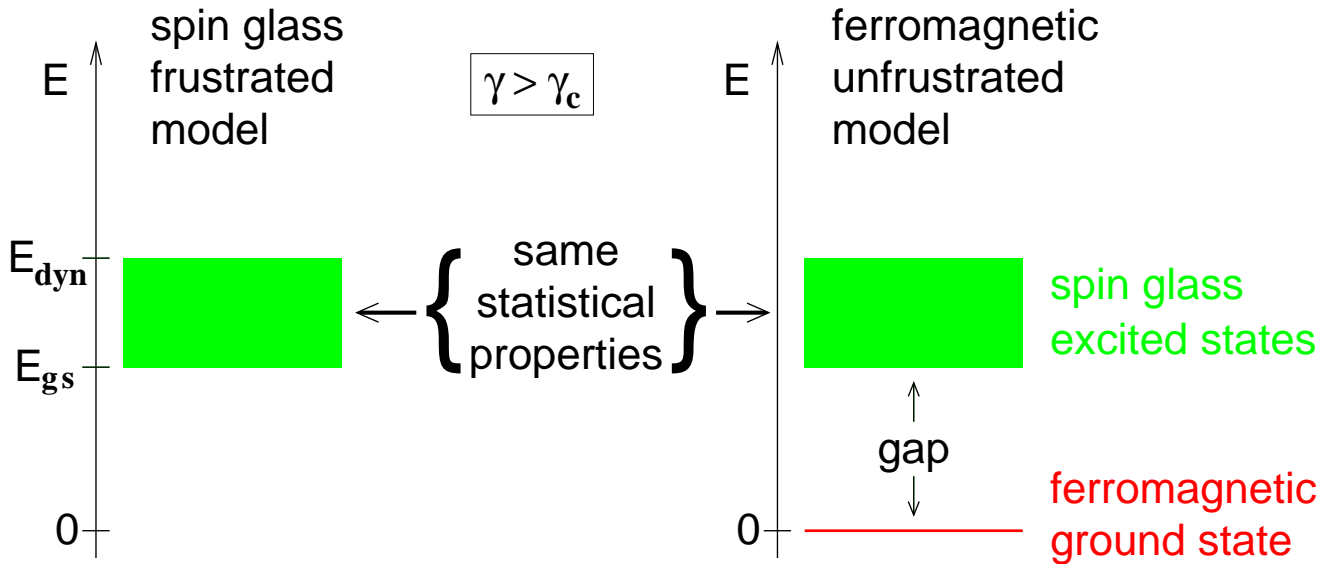
\Rightarrow clustering and search algorithms slowing down

$\Rightarrow \gamma_c = \text{SAT/UNSAT threshold for 3-XOR-SAT}$

\Rightarrow dynamical transition in memory requirements

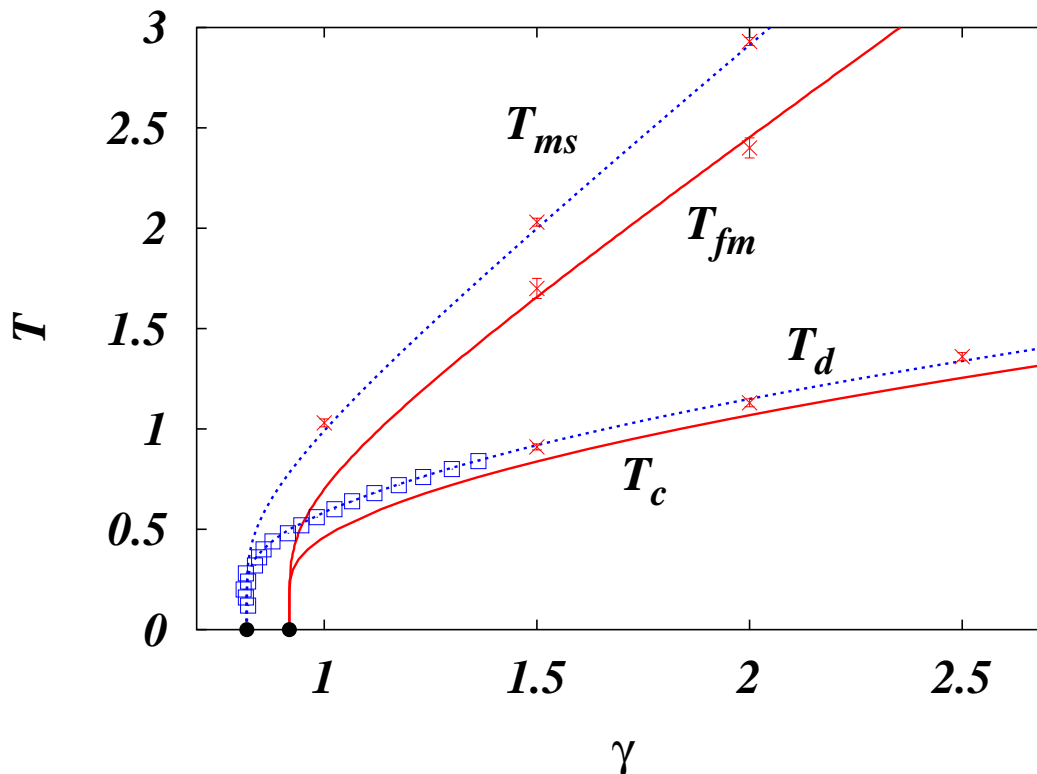
$[\mathcal{O}(N) \rightarrow \mathcal{O}(N^2)]$ solving linear systems in $\text{GF}[2]$

The structure of the configurational space

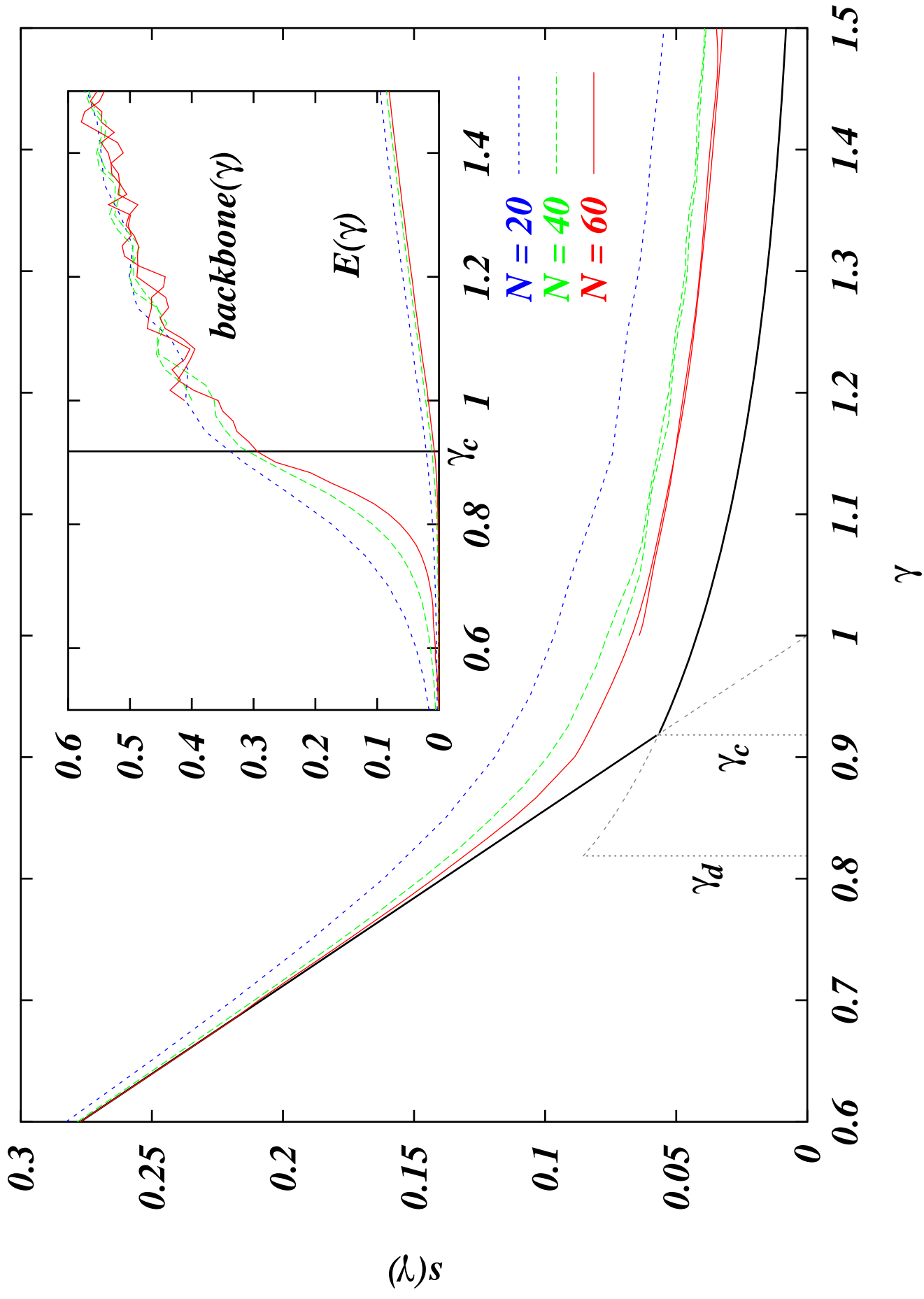


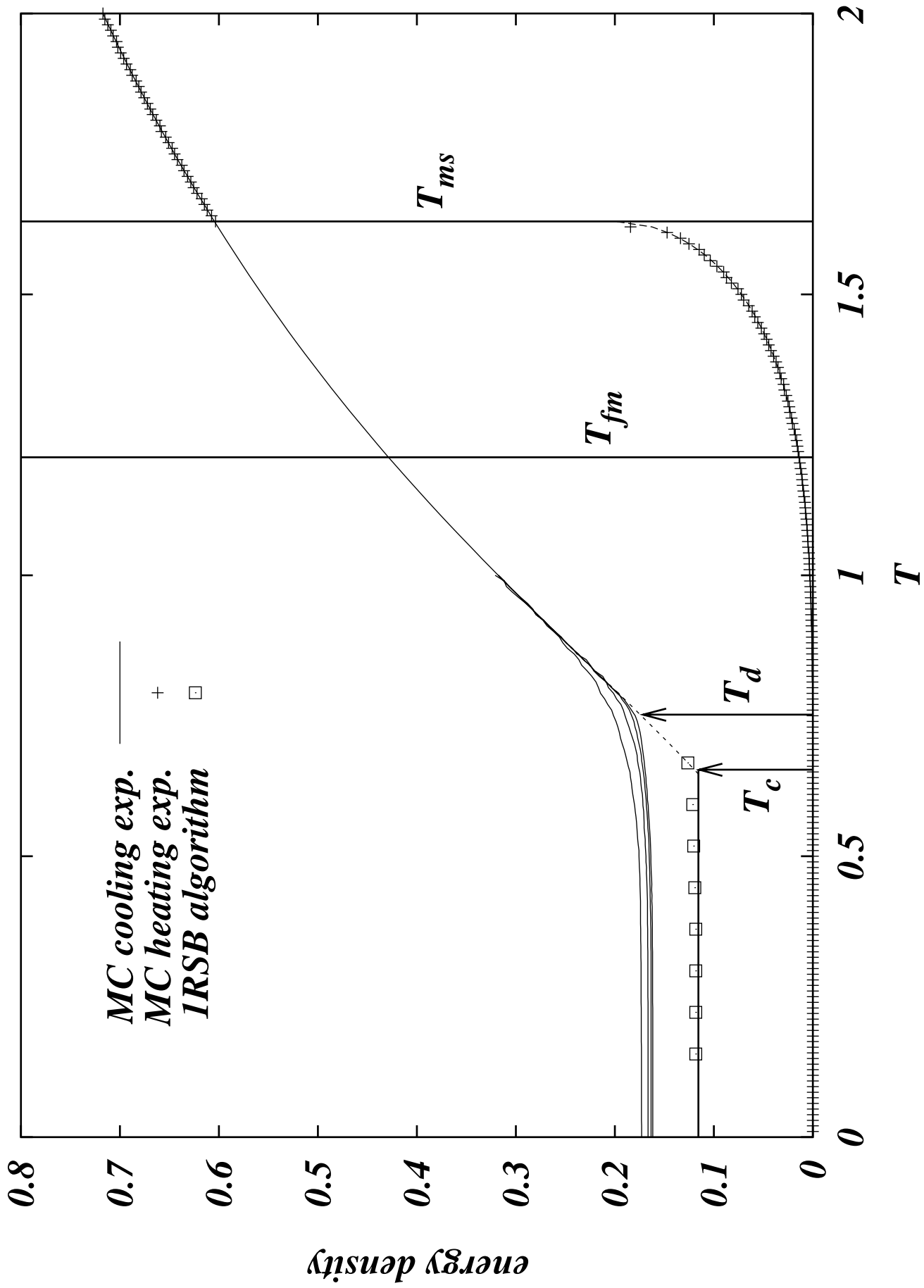
Starting from a **random** configuration, both models have the same **off-equilibrium dynamics**

Phase diagram



T is the temperature and 3γ is the average connectivity





Conclusions and perspectives

New results

- p -spin ($p > 2$) with fluctuating connectivity
 - structure of the configurational space: γ_d , γ_c and $\Sigma(\gamma)$
 - γ_c is the exact threshold for random p -XOR-SAT
 - new transitions in random hypergraphs
- p -spin ($p > 2$) and Bicoloring with fixed connectivity
 - exact 1-RSB solution: GS energy
- K -SAT and Bicoloring
 - variational bounds for α_c (at present the best!)

Some applications

- test-bed for heuristic algorithms: GS energy
- dynamical transitions in Coding and Cryptography
- solvable models for glassy systems and granular matter

Examples of open issues

- complete 1-RSB and FRSB theories (with correlations)
- out of equilibrium dynamics
- analysis of randomized algorithms
- better analysis of the configurational space in K -SAT

Some references

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S. Mertens, “*Computational Complexity for Physicists*”, cond-mat/0012185.
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S. Kirkpatrick and B. Selman, *Science* **264**, 1297 (1994).
- **Statistical Mechanics of K -SAT:**
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R. Monasson, R. Zecchina, S. Kirkpatrick, B. Selman and L. Troyansky, *Nature* **400**, 133 (1999).
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S. Franz, M. Mézard, F. Ricci-Tersenghi, M. Weigt and R. Zecchina, preprint cond-mat/0103026.
- **Exact solution for 3-spin at $T = 0$:**
S. Franz, M. Leone, F. Ricci-Tersenghi and R. Zecchina, preprint cond-mat/0103328.